Dynamical Models of Disc Galaxies and Fueling of the Central SMBH

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Special thanks to Nuria Piñol-Ferrer, Javier Blasco-Herrea, Michael Gatchell, Roland Karlsson

Guillermo Haro 2015 Workshop, INAOE, Mexico

- What processes could generate flows in discs
- Where can we find signatures of such processes
- Can we unveil flows in discs
- How do we observe the observables
- Applying the models to data to close in to the SMBH

Gas transport to galaxy centers means losing momentum

On galactic scales, this can be initiated and driven by

- Rotational magnetic instabilities (Sellwood & Balbus 1999)
- Minor Mergers (Hernquist & Mihos 1995)
- Turbulence (Lynden-Bell & Pringle 1974)
- Shocks (Roberts 1969)
- Sheared gravitational instabilities (Goldreich & Lynden-Bell 1965)
- Density waves (Lindblad 1964; Lin & Shu 1964)

Nuclear spirals correlate with nuclear activity

• Flocculent or Grand-design

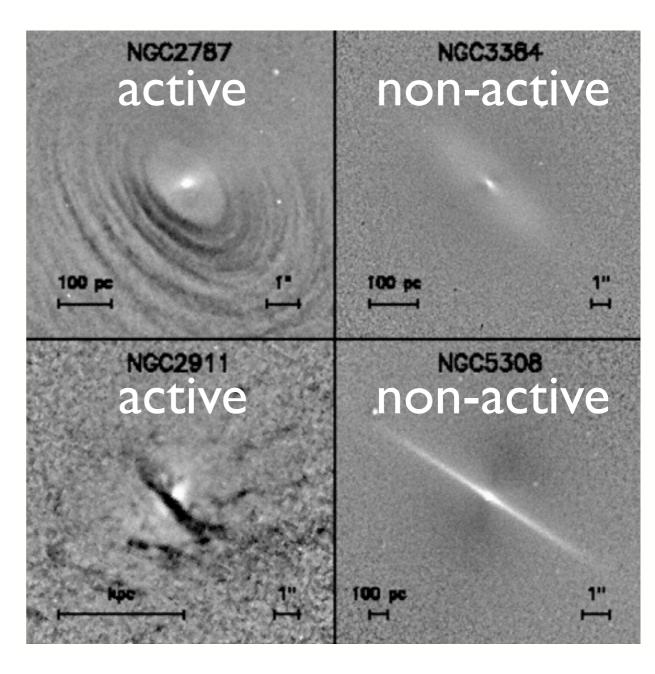
e.g., Elmegreen et al. (1998) Montenegro (1998) Englmaier & Shlosman (2000) Maciejewski (2004)

• <u>Deviations in density ~ 5-10 %</u>

e.g., Lou et al. (2001) Prieto et al. (2005) Fathi et al. (2011)

• Non-circular motions are significant

e.g., Fathi (2004) Storchi-Bergmann et al. (2007)

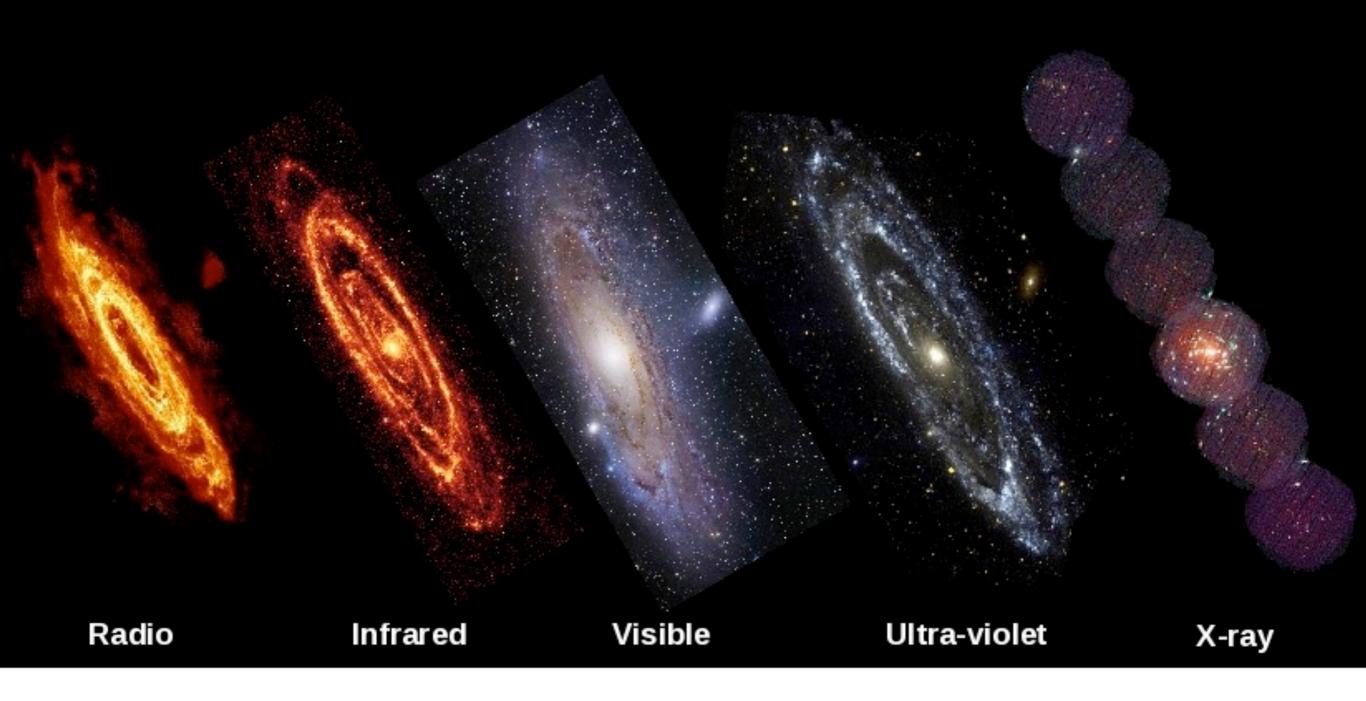


Martini & Pogge (1999) Regan & Mulchaey (1999) **Density wave theory**

DEVELOPMENT OF A "BAR-TYPE" DENSITY WAVE IN THE CENTRAL LAYER OF A GALAXY

Per Olof Lindblad Stockholm Observatory April 1961 ttt.astro.su.se/~po

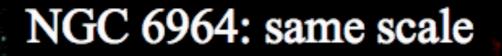
The importance of observing different tracers



Radio:WSRT/R. Braun; Infrared:NASA/Spitzer/K. Gordon; Visible: R. Gendler; Ultraviolet: NASA/GALEX; X-ray: ESA/XMM/W. Pietsch

The importance of observing different tracers

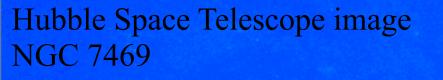
Optical (stars)

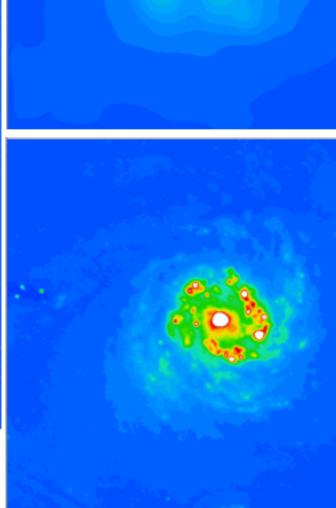


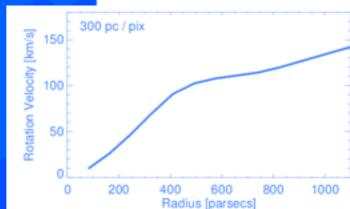
radio 21cm (hydrogen gas)

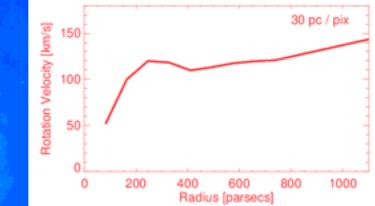
Boomsma (2008)

The importance of high-resolution data

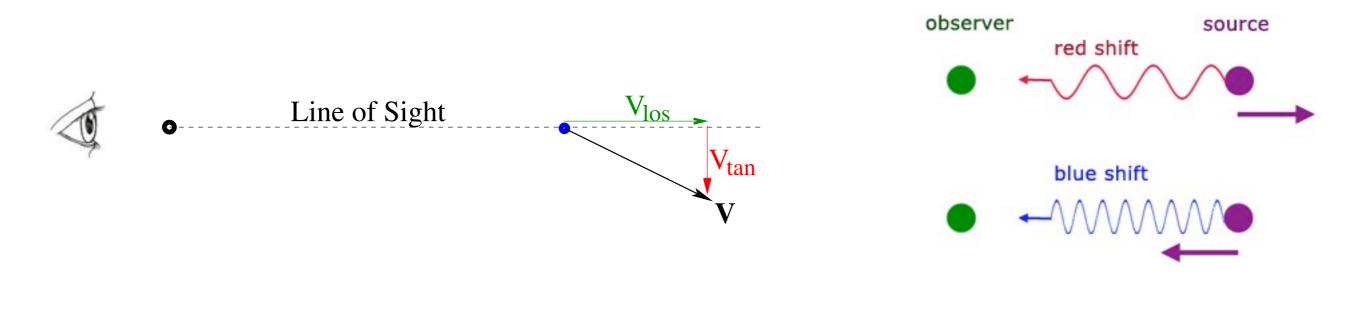


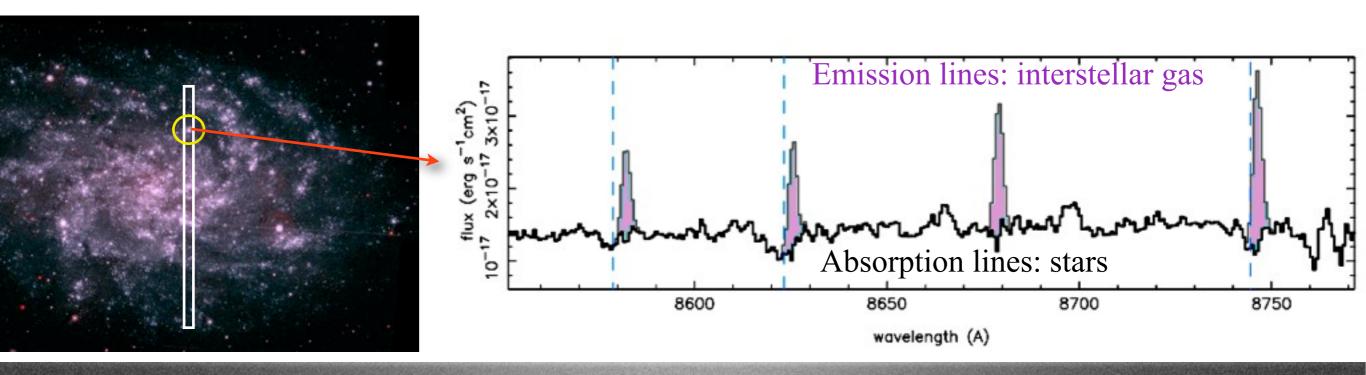




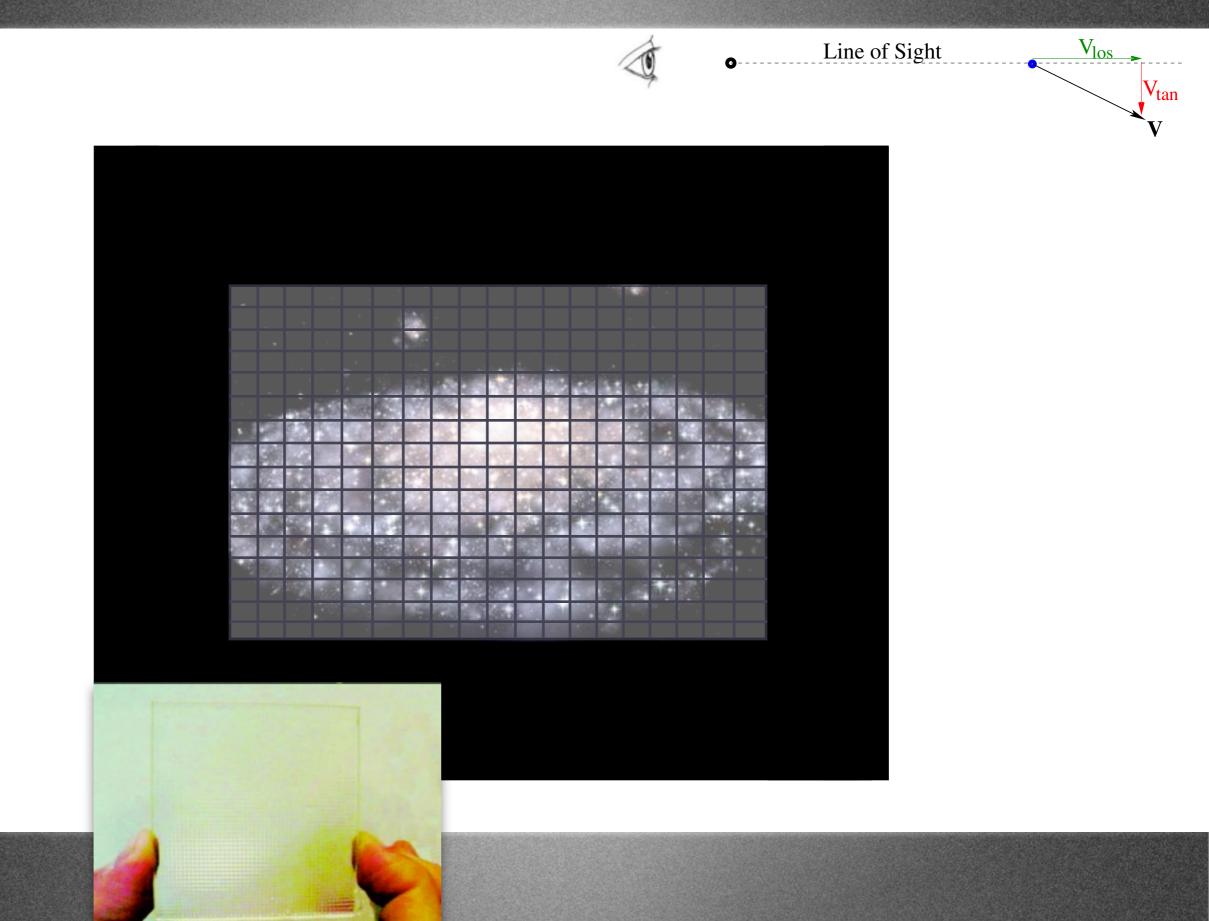


Principles of observing kinematic features

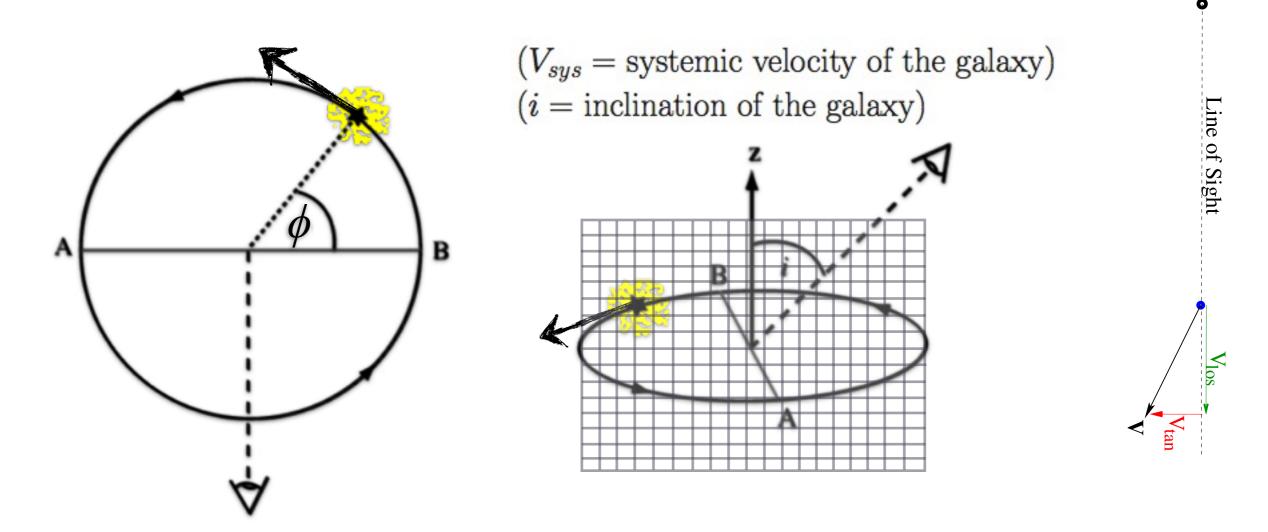




Two-dimensional spectroscopy



Interpreting kinematic observations



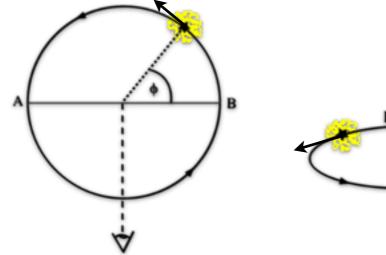
Viewed edge-on:

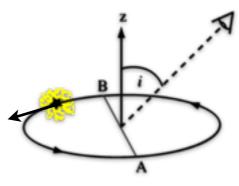
$$V_{los}(R, i = 90^{\circ}) = V_{sys} + V(R)\cos\phi$$

Viewed at an angle:

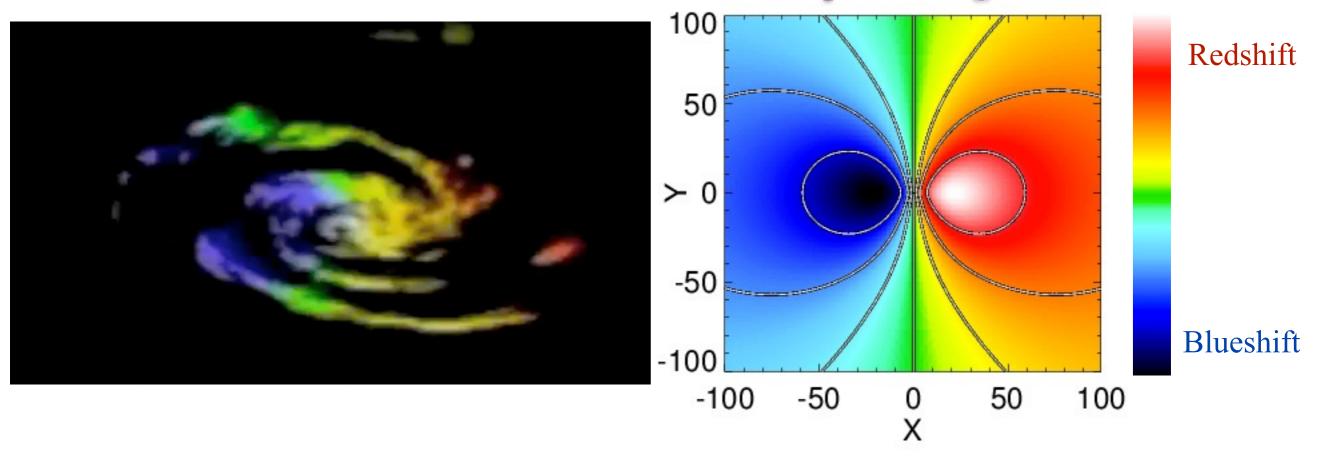
 $V_{los}(R,i) = V_{sys} + V(R) \cos \phi \sin i$

Two dimensional velocity fields

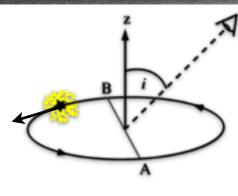




A Spider Diagram

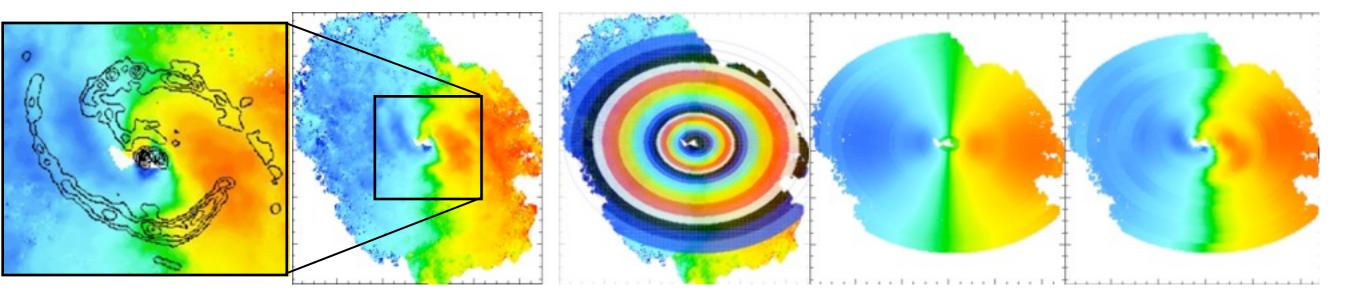


Analysing high-resolution data



Assuming differential rotation, we section a velocity field into concentric rings and fit the disk centre, inclination, PA, systemic and rotation velocity

 $V_{los}(R,i) = V_{sys} + V(R) \cos \phi \sin i$



$$V_{\text{los}} = V_{\text{sys}} + \sum_{n=1}^{k} \left(c_n \cos(n\phi) + s_n \sin(n\phi) \right) \sin i$$

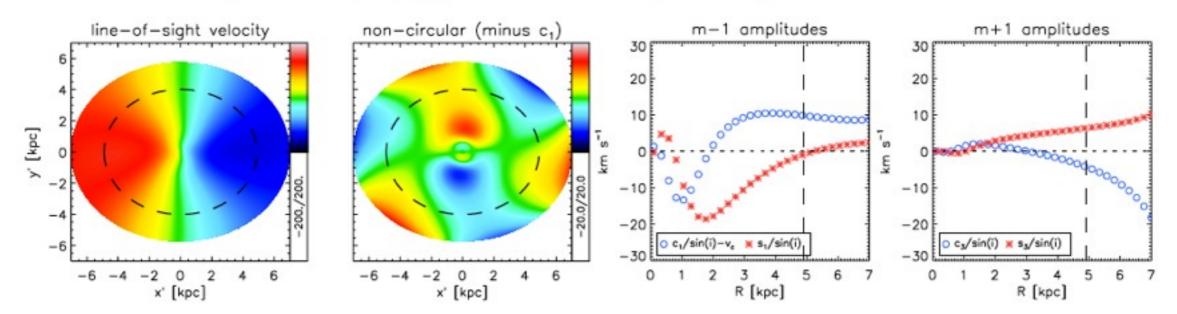
Fathi et al. (2005; 2007; 2008); van de Ven & Fathi (2010)

Modelling a weakly perturbed gravitational potential

 $\Phi(R,\theta) = \Phi_0(R) + \Phi_m(R) \cos[m(\theta - \Omega_p t)]$

where Φ_0 is the axisymmetric potential, Ω_p is the angular speed of the rotating potential, and *m* is the harmonic number of the distortion. For the case of m = 2 the potential is barred. The equations of motion are written as

$$\begin{aligned} &-\frac{\partial \Phi}{\partial R} &= \ddot{R} + 2\lambda \dot{R} + (\dot{\phi} + \Omega_p)^2 R \\ &-\frac{\partial \Phi}{R\partial \phi} &= R \ddot{\phi} + 2(\dot{\phi} + \Omega_p) \dot{R} \end{aligned}$$



$$\ddot{R}_1 + 2\lambda\kappa_0\dot{R}_1 + \kappa_0^2R_1 = -R_0(A\cos\eta + B\sin\eta)$$

Shlosman & Noguchi (1993); Wada (1994); Fathi (2004); van de Ven & Fathi (2010)

More elaborate analytic dynamical models

• Apply the formalism of Lindblad et al. (1996)

see also Wada (1994) and Sakamoto et al. (1999)

$$\Phi(r,\theta) = \Phi_0(r) + \Phi_1(r,\theta) = \Phi_0(r) - \sum_{m=1}^n \Psi_m(r) \cos m(\theta - \vartheta_m(r))$$

• We advance the method by treating the bar and spiral perturbation simultaneously

• Solve the linearised equations of motion in polar coordinates

$$\ddot{\xi} + 2\lambda\dot{\xi} - 2\Omega\dot{\eta} - 4\Omega A\xi = -\frac{\partial\Phi_1}{\partial r} = \sum_{m=1}^n [C_m \cos m(\theta - \vartheta_m) + E_m \sin m(\theta - \vartheta_m)]$$
$$\ddot{\eta} + 2\Omega\dot{\xi} + 2\lambda\dot{\eta} + 4\lambda A\xi = -\frac{1}{r}\frac{\partial\Phi_1}{\partial\theta} = -\sum_{m=1}^n D_m \sin m(\theta - \vartheta_m),$$

• Calculate the orbits in an arbitrary gravitational potential and for given viewing angle, damping efficiency, bar speed and strength and spiral pitch angle)

Deriving flows in the disc

$$V_{\rm los} = \sin i \left[v_{\phi}(R, \psi) \cos \psi + v_R(R, \psi) \sin \psi \right]$$

 $v_R(R, \psi) = v_c(R)[c_R \cos m\psi + s_R \sin m\psi]$

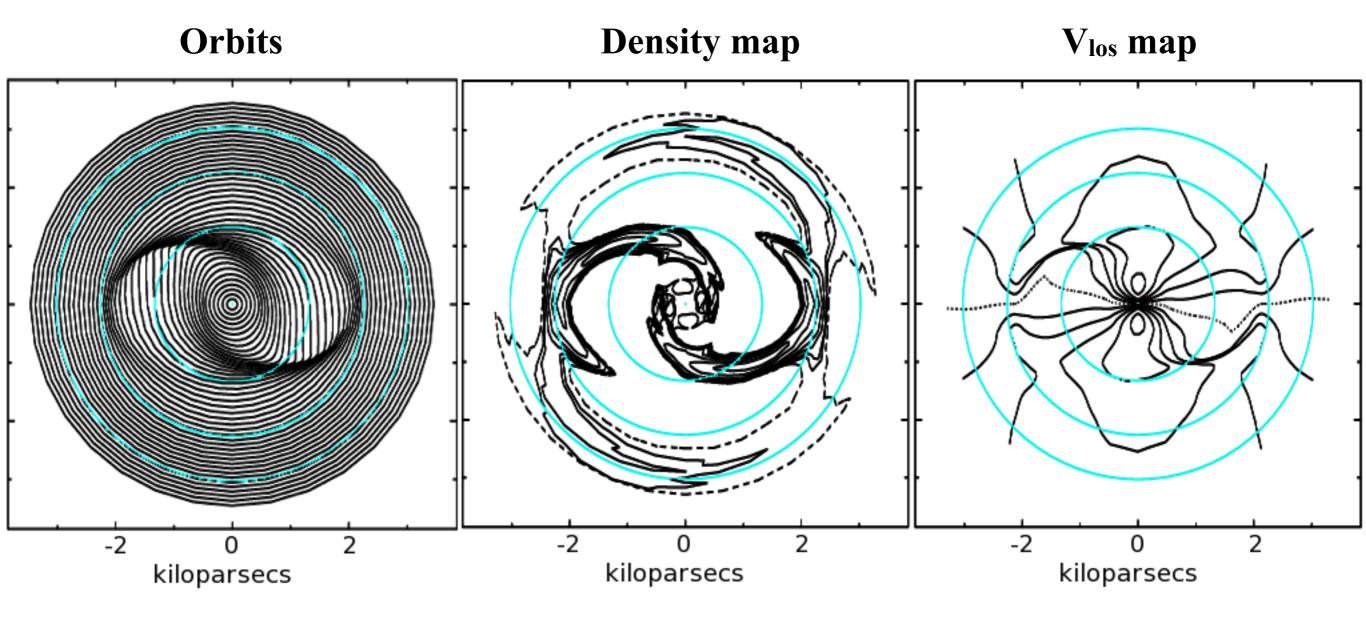
 $v_{\phi}(R, \psi) = v_c(R)[1 + c_{\phi} \cos m\psi + s_{\phi} \sin m\psi]$

$$v_{\text{flow}} = \left(\frac{\Lambda^2}{\Delta^2 + \Lambda^2}\right)^{1/2} v_R$$

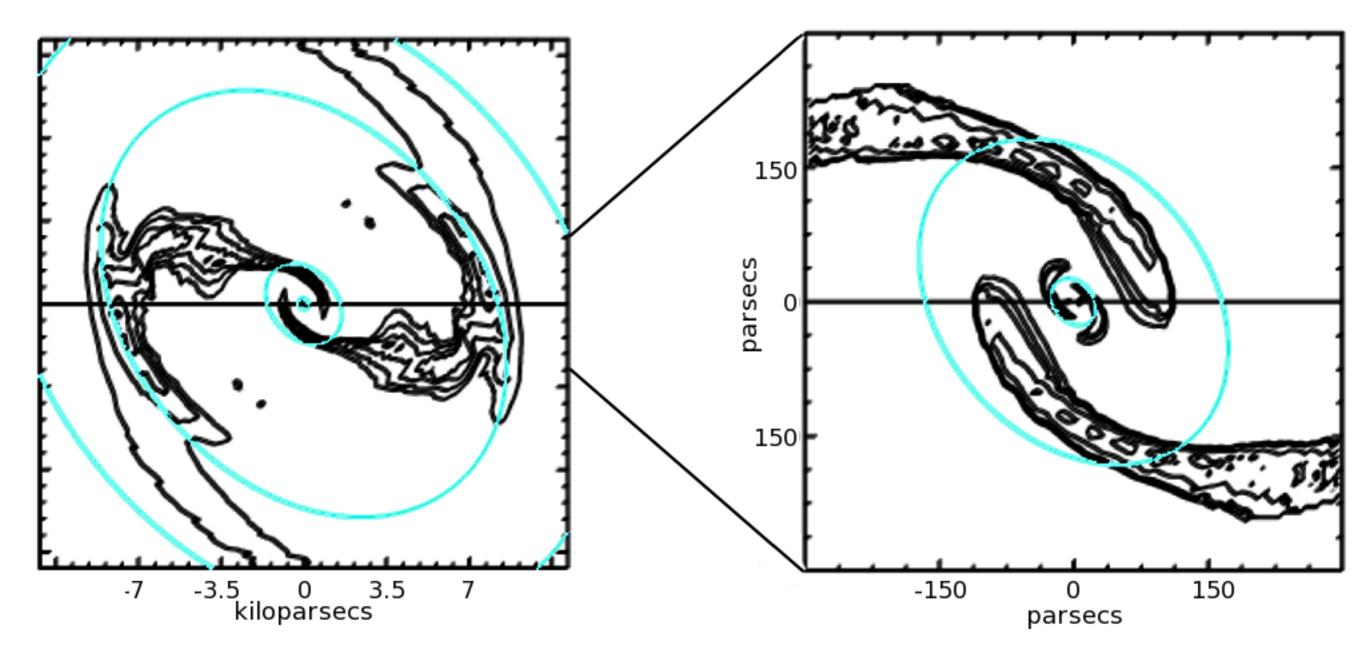
 $\Delta = \kappa^2 - m^2 (\Omega - \Omega_p)^2, \quad \Lambda = 2\lambda \kappa m (\Omega - \Omega_p)$

functions of $\Omega(R)$ and $\kappa(R)$

Orbits, densities and line of sight velocities

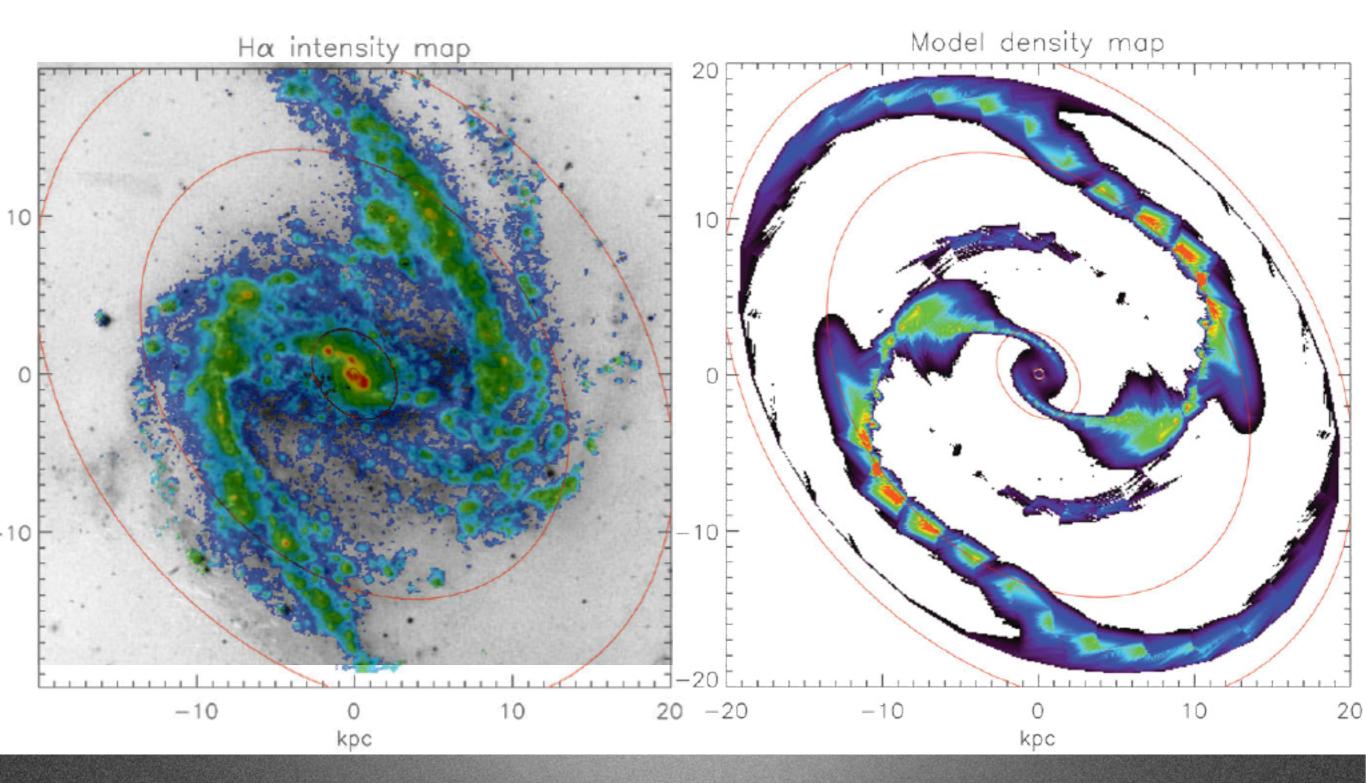


Zooming into the centre



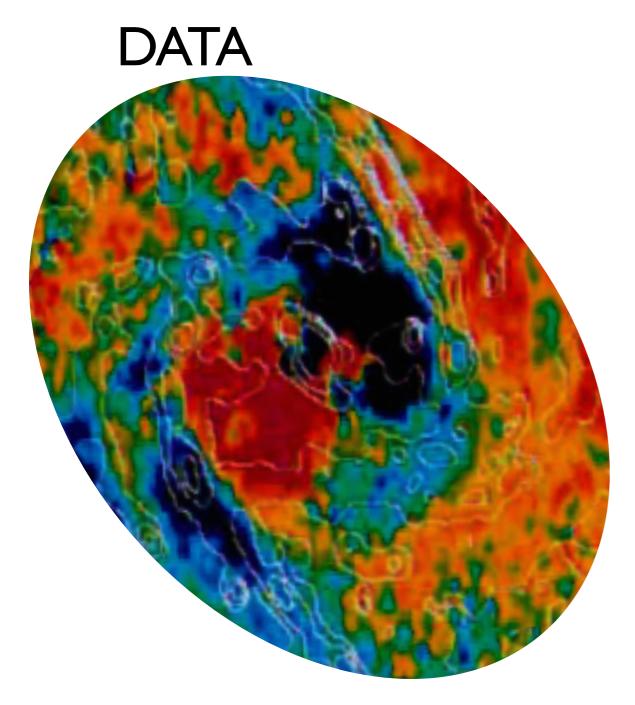
Modelling observational data: ionized gas

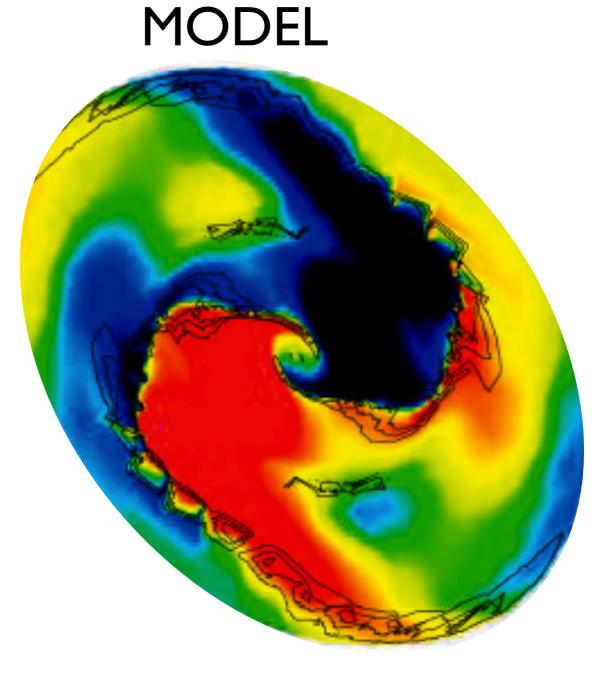
NGC 1365 @ 17 Mpc distance



Piñol-Ferrer, Lindblad, Fathi (2012)

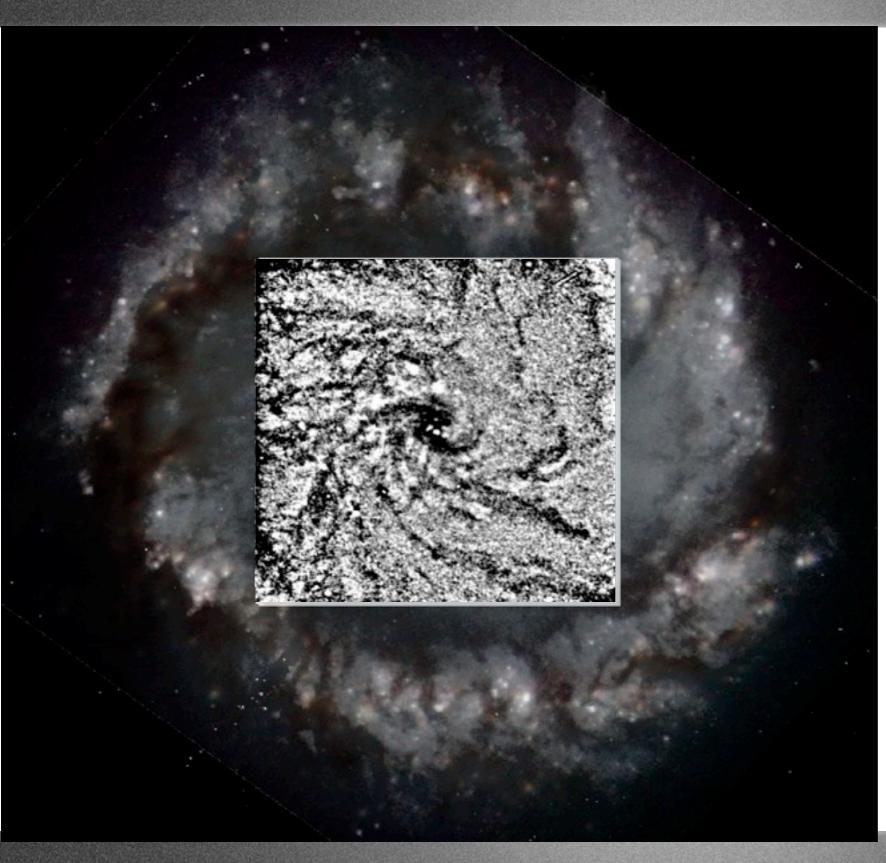
Modelling observational data: neutral gas





Jörsäter & van Moorsel (1996)

Piñol-Ferrer, Lindblad, Fathi (2012)



NGC 1097 @ 14.5 Mpc distance

Large-scale dynamics:

Fabry-Perot data VLA data

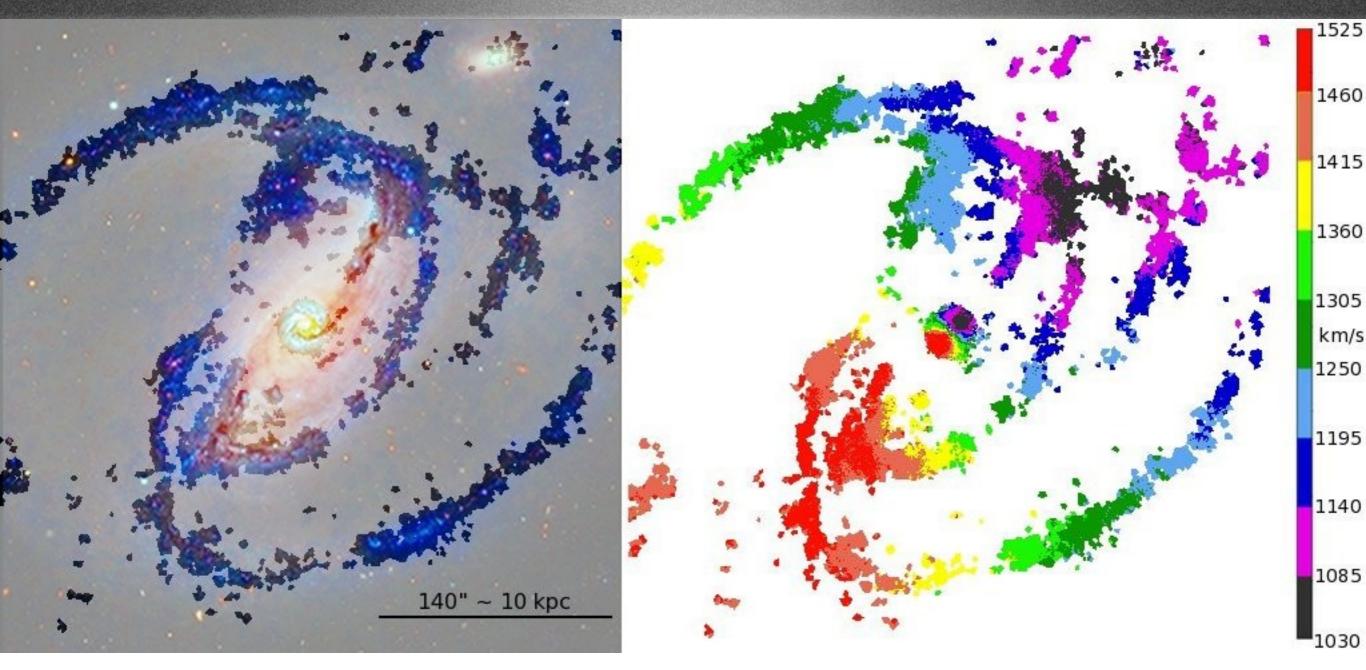
Small-scale dynamics:

8-m Gemini 2D spectra SMA data (archival) ALMA (Cycle 0 and 1)

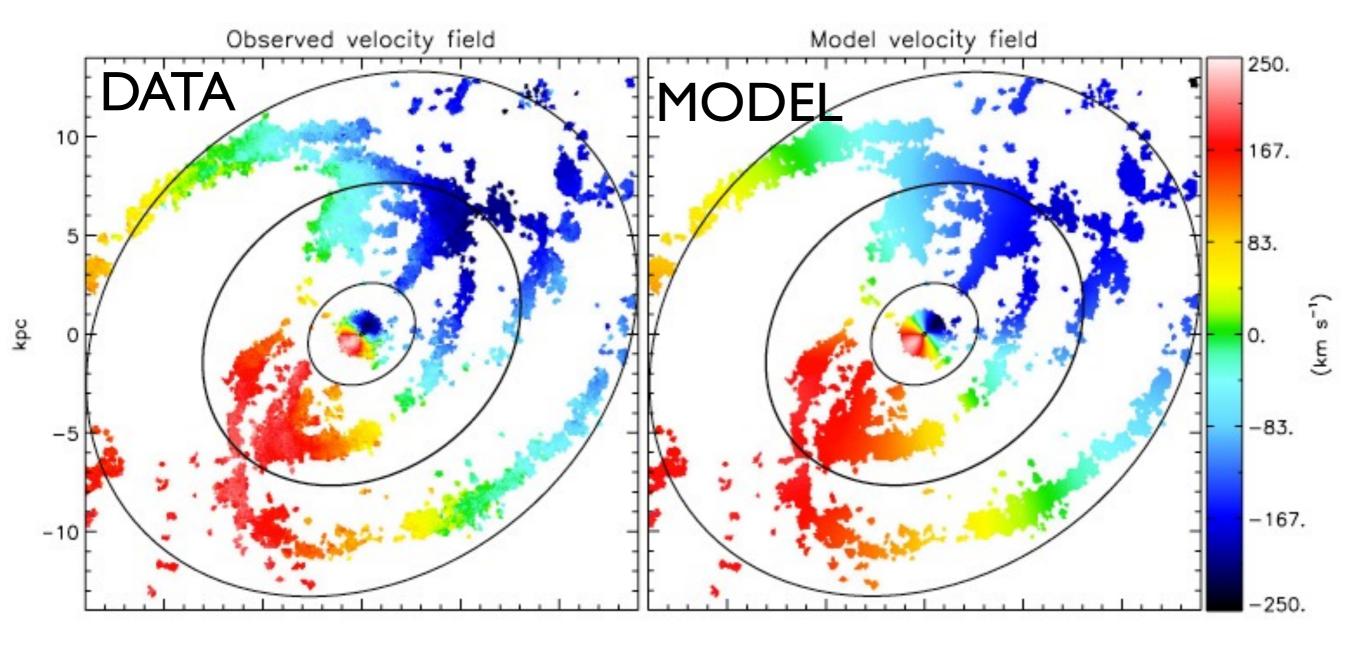
Single-dish: APEX, JCMT, IRAM

+ our own Hubble Space Telescope images and spectra (GO 9872)

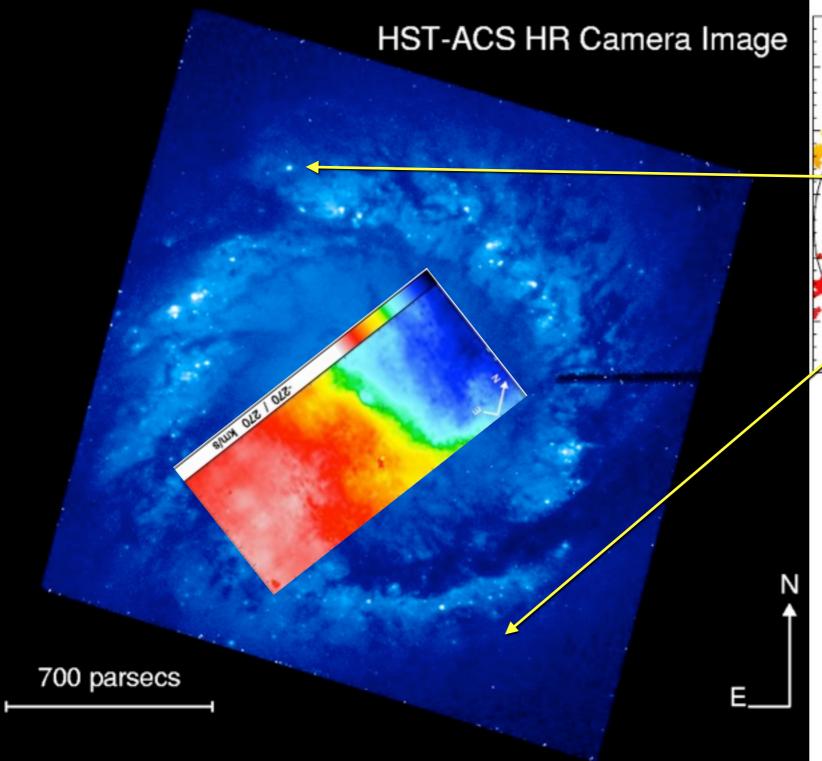
Fathi et al. (2006)

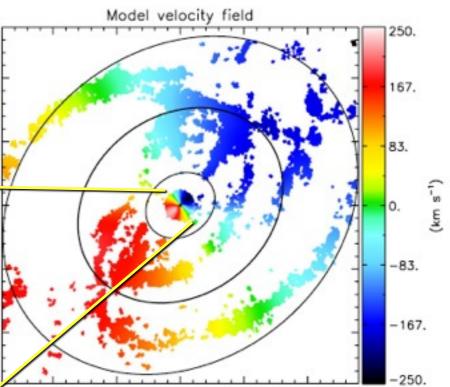


H α line observed with our Fabry-Perot Interferometer covering 8'x8' field at 0.8"/pix with a velocity sampling = 15 km/s VLA DnC configuration data at 21 km/s and beam FWHM of 56"



Piñol-Ferrer, Fathi et al. (2014)

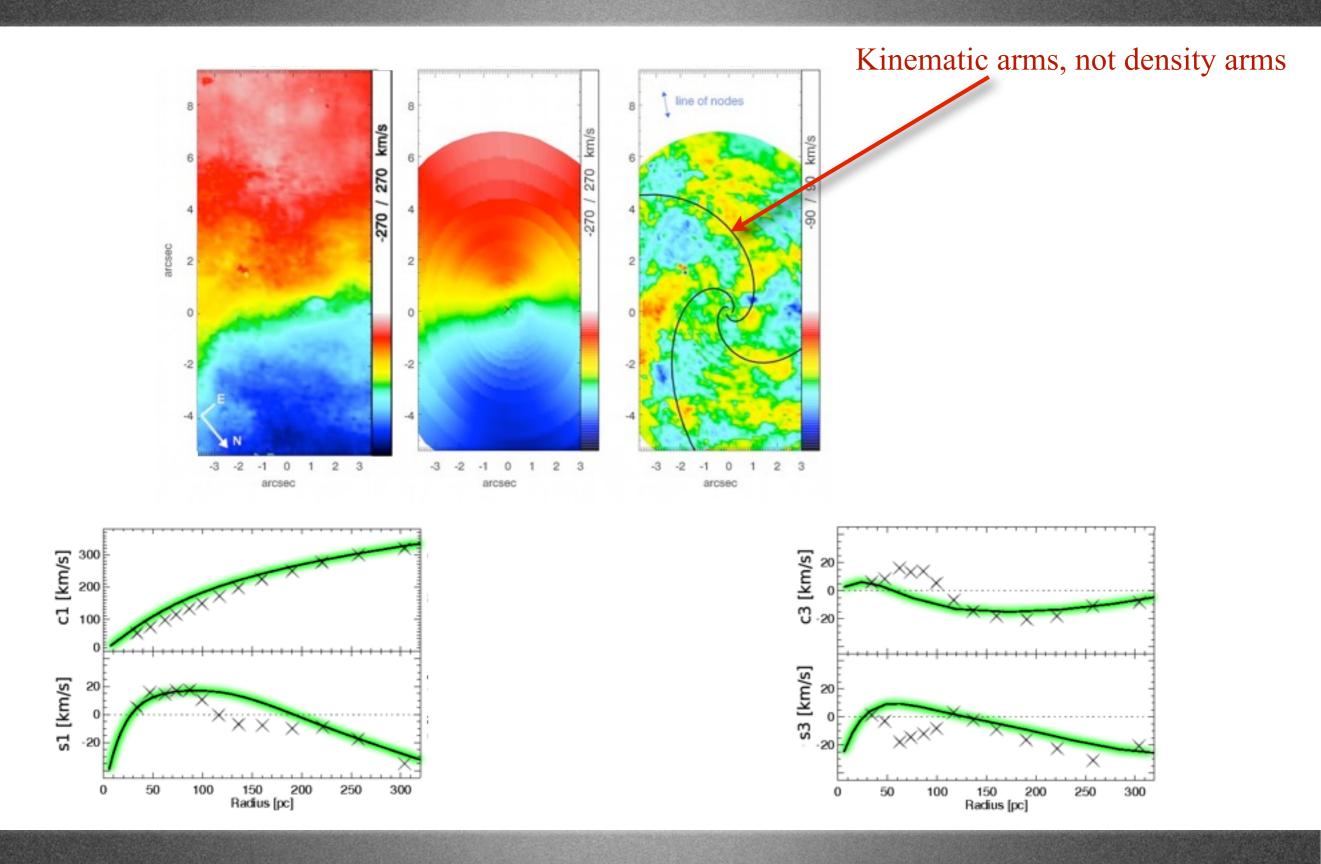




GMOS-2D spectra cover Hα, [NII], [SII] lines at 0.1"/pix (0.8" / pix in big map) at 85 km/s (FWHM)

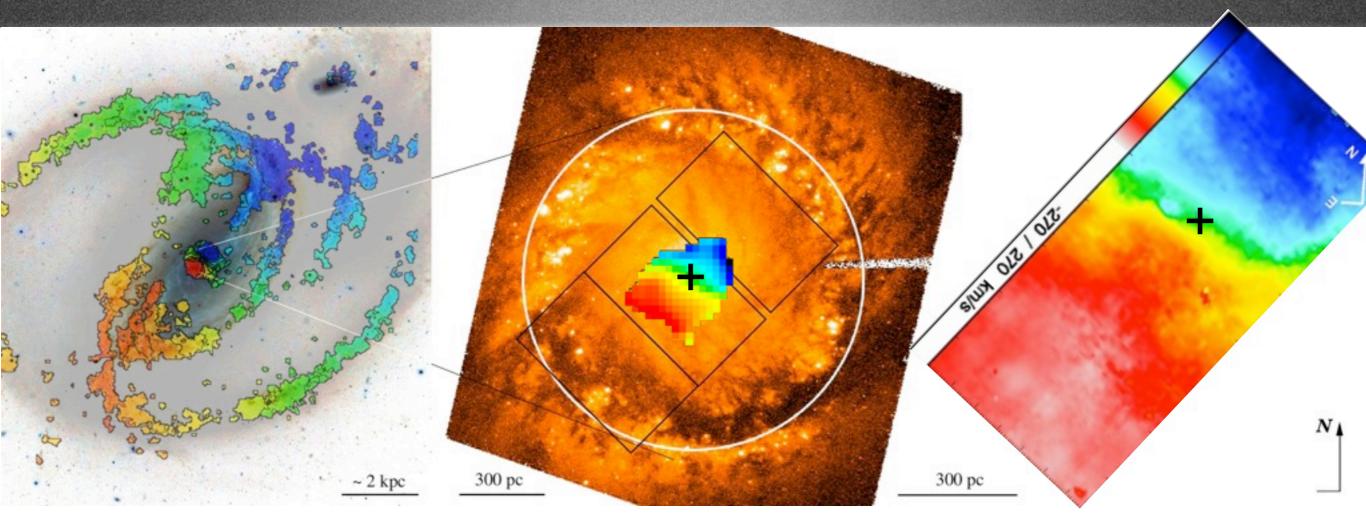
Fathi et al. (2006)

Direct evidence for flows along nuclear spiral arms



Fathi et al. (2006); van de Ven & Fathi (2010)

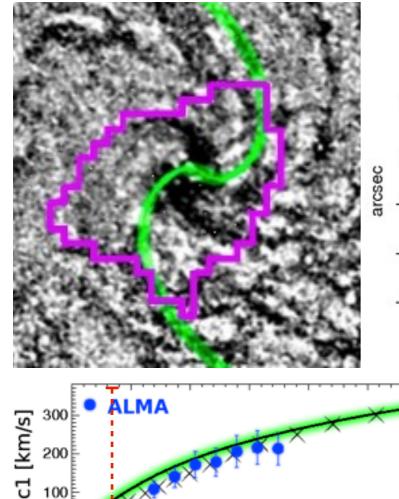
Synergy between optical, mm and radio observations



- ALMA Band 7 observations with 15 antennas (Cycle-0 data: PI Kotaro Kohno)
- HCN(4–3) line at rest frequency of 354.505 GHz (488.28125 kHz)
- Primary beam of 18.1"
- Synthesized beam of 1.5" x 1.2" (105 x 84 pc) sampled at 0.3"/pix (56"/pix with VLA)

Fuelling the Supermassive Black Hole in NGC 1097

with model density arms

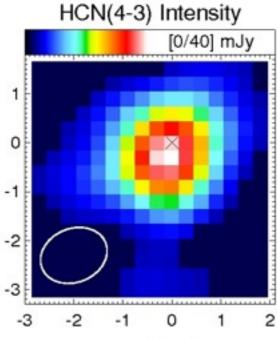


s1 [km/s]

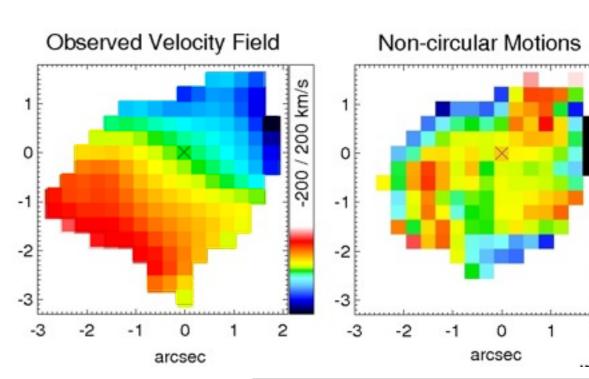
20 [-

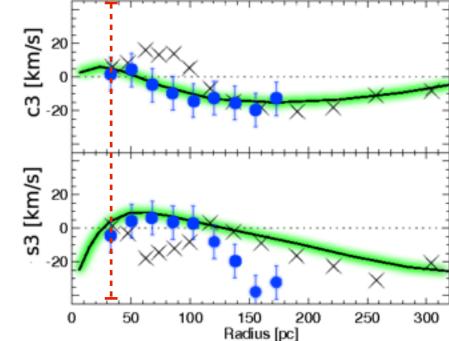
-20

Radius [pc]



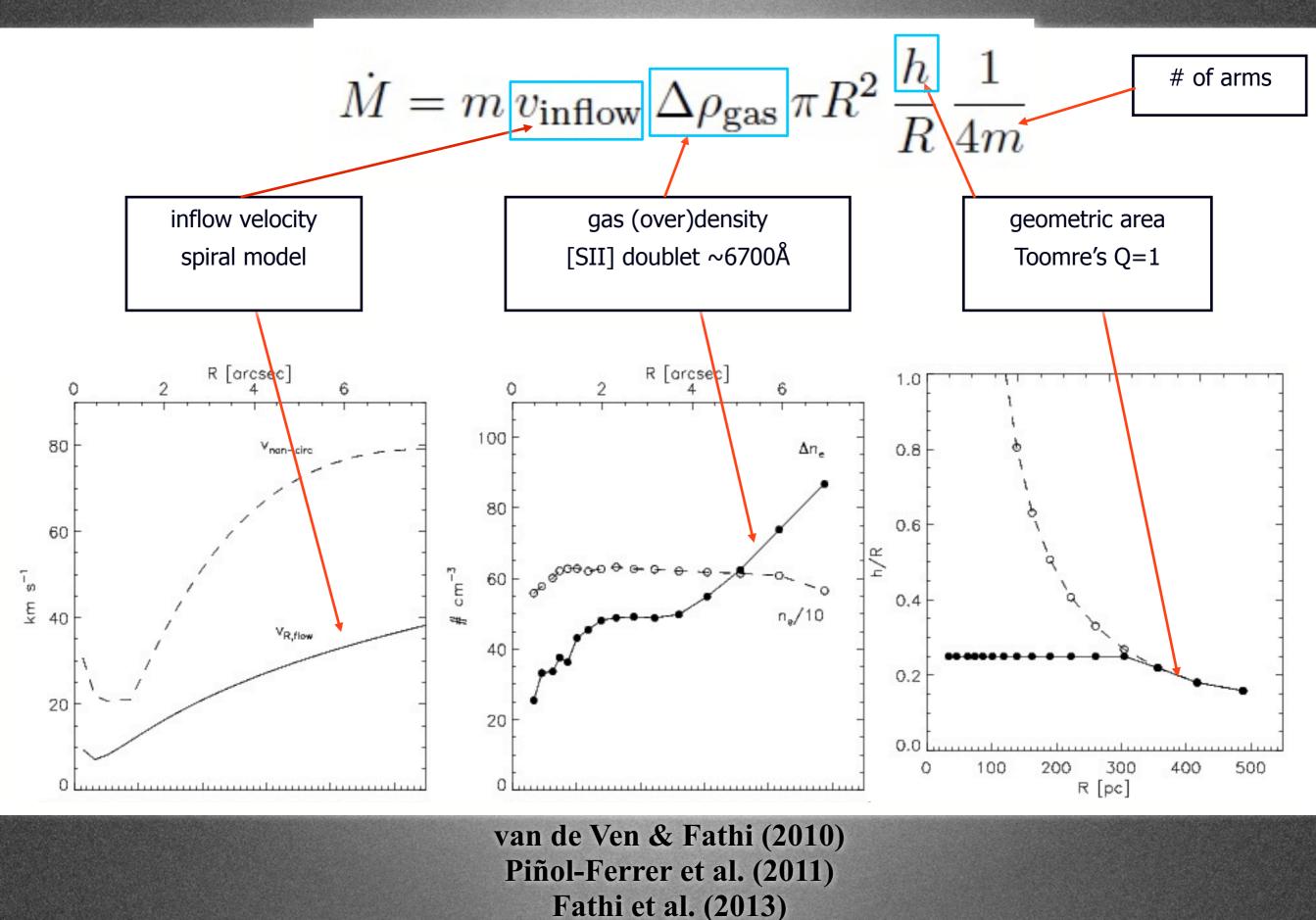






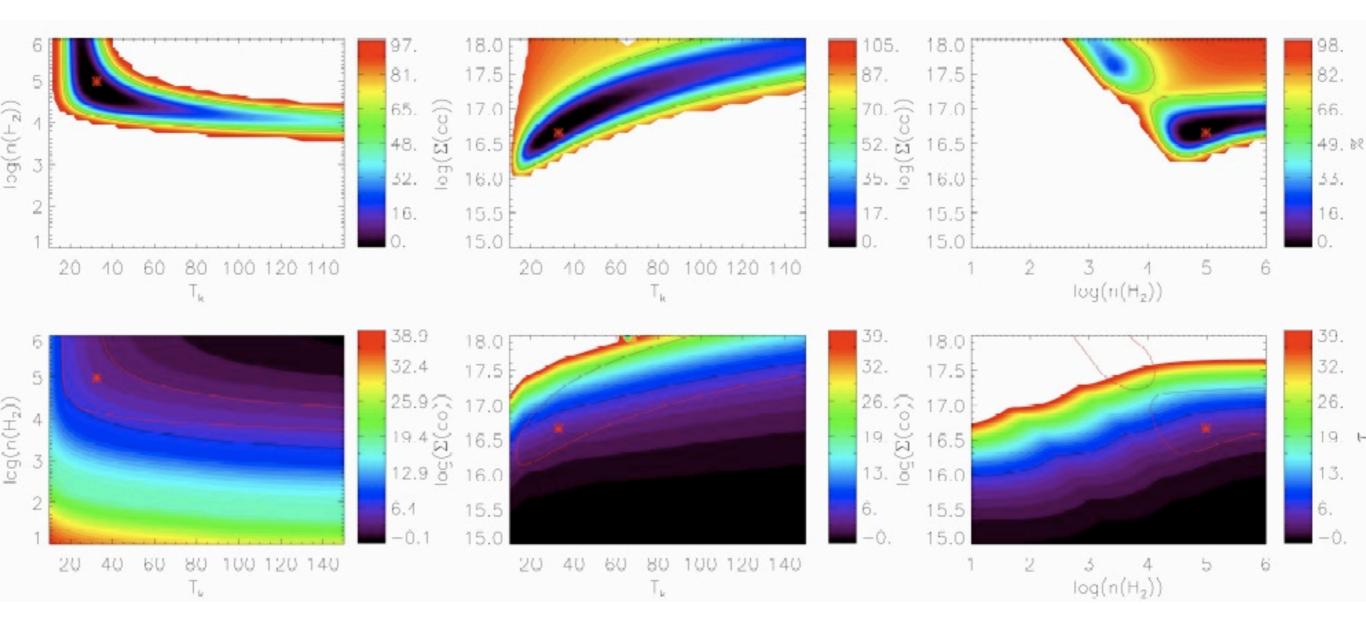
-50 / 50 km/s

Ionised and molecular gas inflow rates



Ionised and molecular gas inflow rates

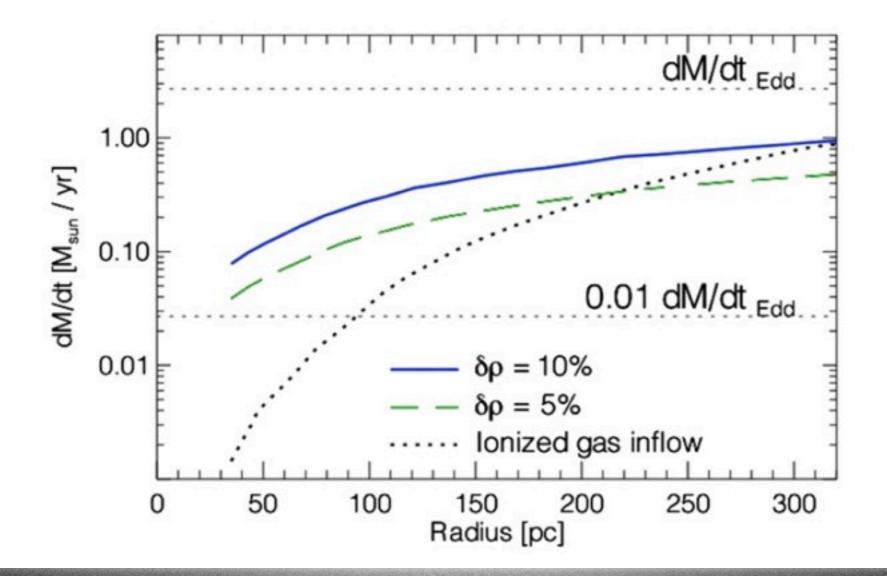
- [SII] lines to measure the ionised gas
- CO isotopologues to measure the molecular gas (Local thermal equilibrium method)
- HCN to measure the densest cores



Piñol-Ferrer et al. (2011); Fathi et al. (2013)

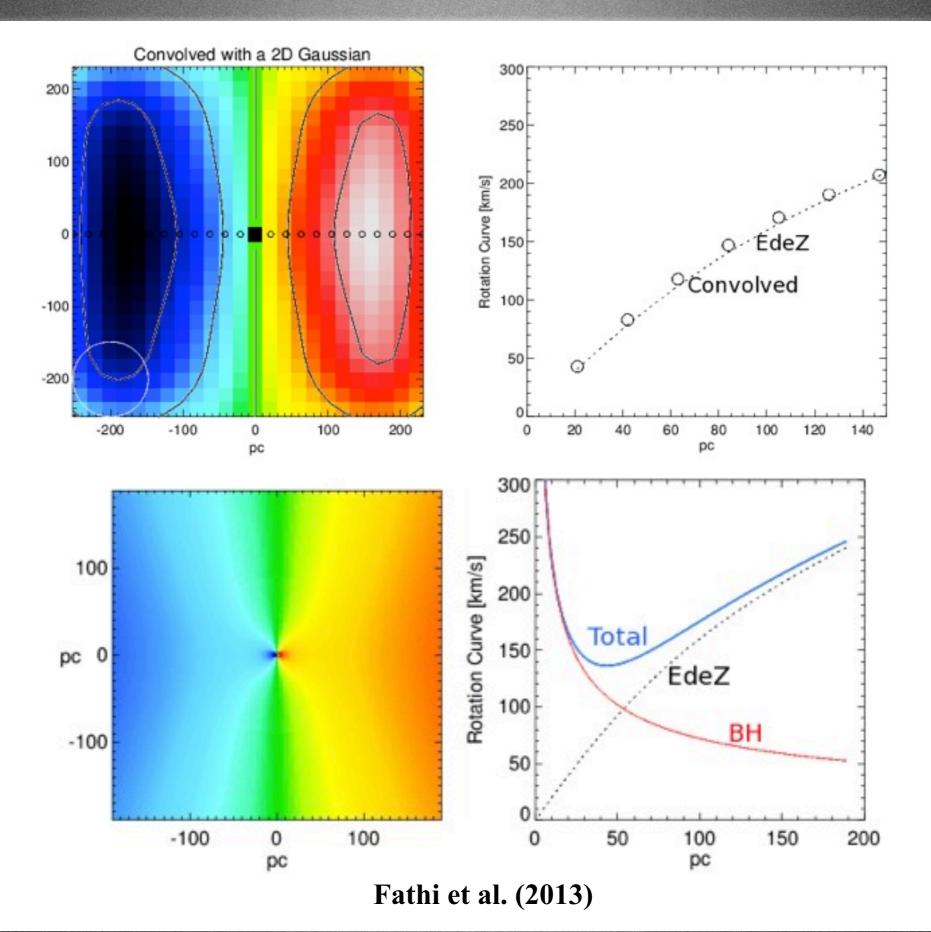
Ionised and molecular gas inflow rates

- NGC 1097 hosts a 1.2×10⁸ M☉ Supermassive Black Hole
- This implies an Eddington accretion rate ($\epsilon=0.1$) = (dM/dt)_{Edd} = 2.7 M \odot /yr



Piñol-Ferrer et al. (2011); Fathi et al. (2013)

How close are we to the Supermassive Black Hole



- Interstellar gas flows and galactic scale dynamics are best studied using a multiwavelength approach (and using 2D spectroscopy)
- The tools that we have built and developed over the past 15 years have proven successful in a number of observed galaxies
- Our new analytic approach reproduces all main observed features seen in various wavelengths: outer arms, main bar and the nuclear spirals (for the first time)
- ALMA Cy-0 observations kinematically confirm the nuclear spirals down to 40 pc from the 10⁸ M⊙ Supermassive Black Hole in NGC 1097
- Gas inflow rate upper limit = $0.2 \text{ M} \odot/\text{yr} (0.07 \text{ M}_{\text{Edd}})$ at 40 pc distance

Thank you