



Formation of Massive Black Hole Seeds in First Galaxies

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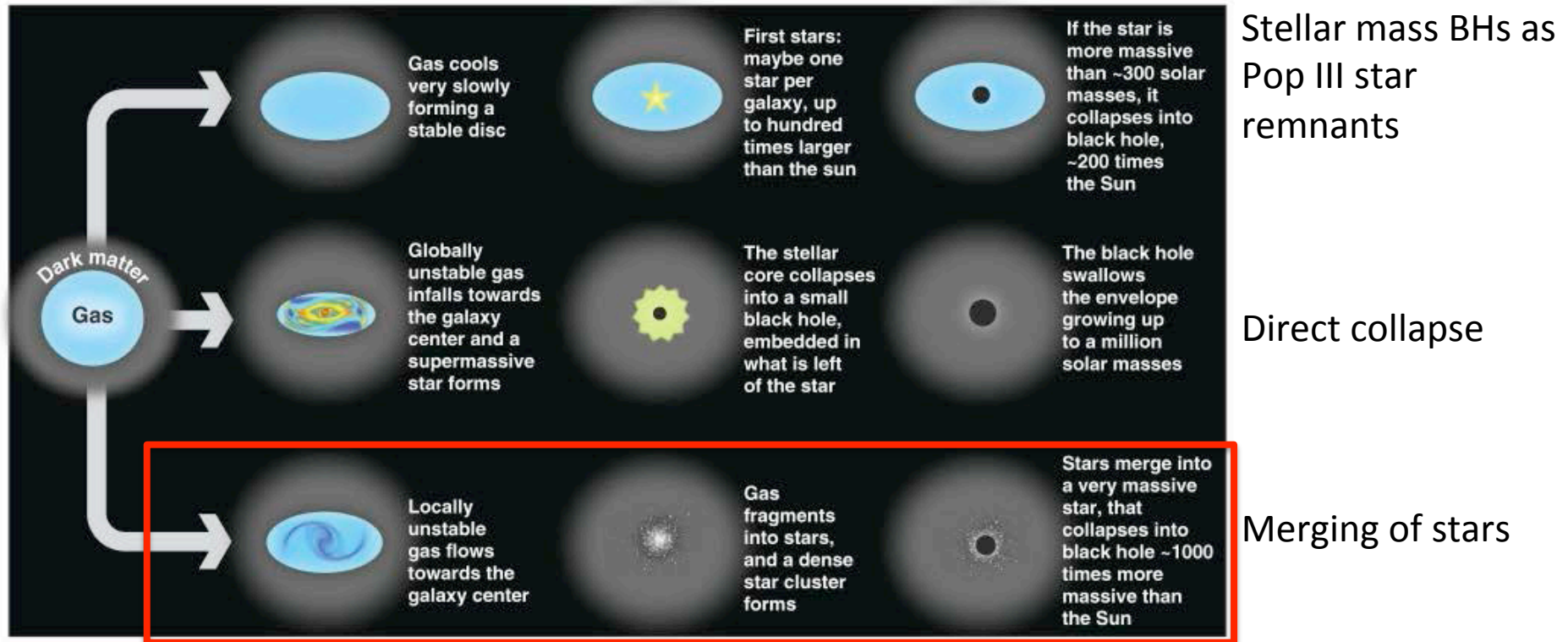
Sadegh Khochfar (Edinburgh)

“Forming and Fueling Supermassive Black Hole Seeds”

Guillermo Haro 2015 Workshop: July 6-24, 2015, Tonantzintla Puebla, Mexico

Models of formation of MBHs

Volonteri (2012)



See also Rees (1984)

Star clusters

Galactic Bulge



$M_{\text{star}} \sim 10^{10-12} M_{\text{sun}}$
 $R \sim 5 \text{ kpc}$



Open clusters

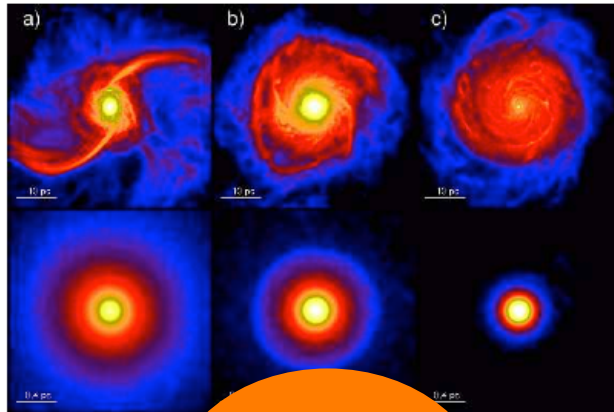
$M_{\text{star}} \sim 10^{2-3} M_{\text{sun}}$
 $R \sim 10 \text{ pc}$

Globular clusters

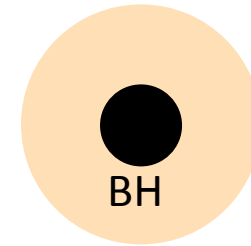


$M_{\text{star}} \sim 10^{4-7} M_{\text{sun}}$
 $R \sim 5 \text{ pc}$

Formation of dens star clusters



(Mayer+10, Nature)



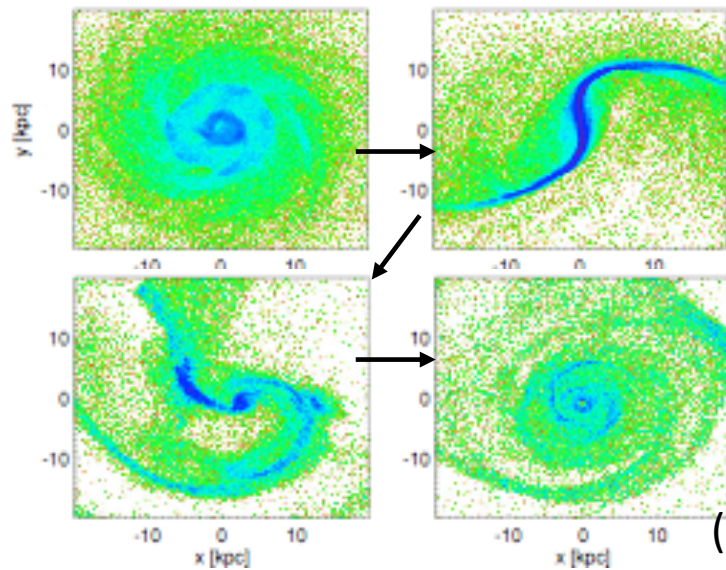
Super-massive star

-> Direct collapse

(e.g., Begelman06; Lodato&Natarajan06)

Compact gas cloud

Or



(Hopkins+09)

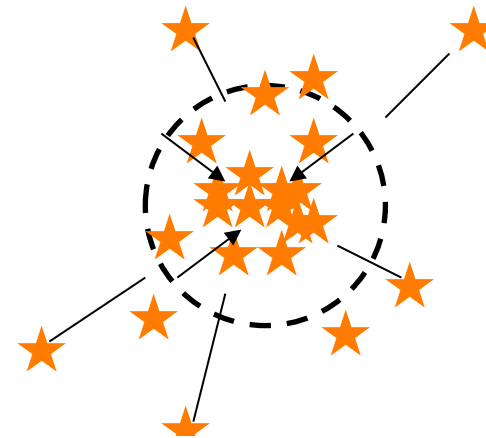
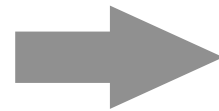
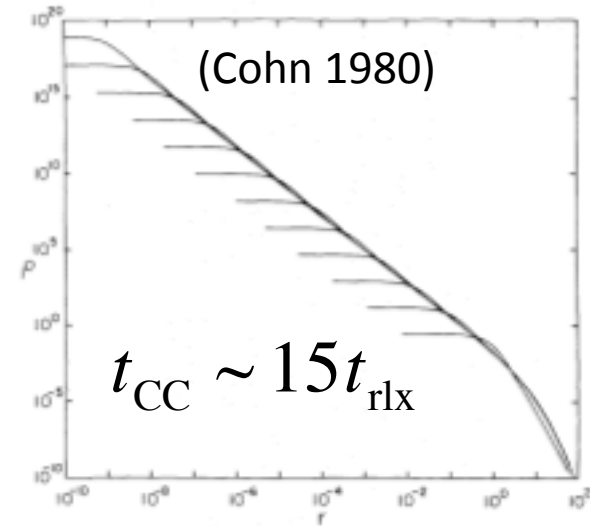
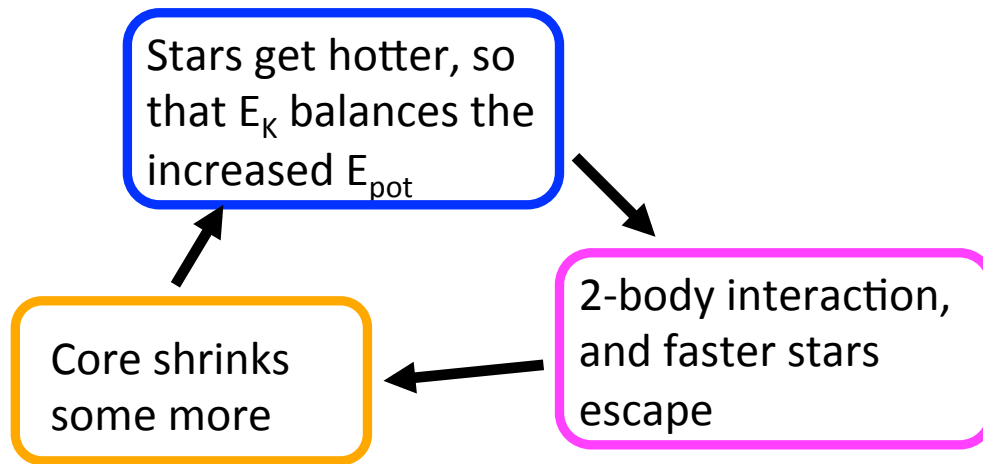


Compact star clusters

(e.g., Ferrara+13, Harley+15)

Core collapse

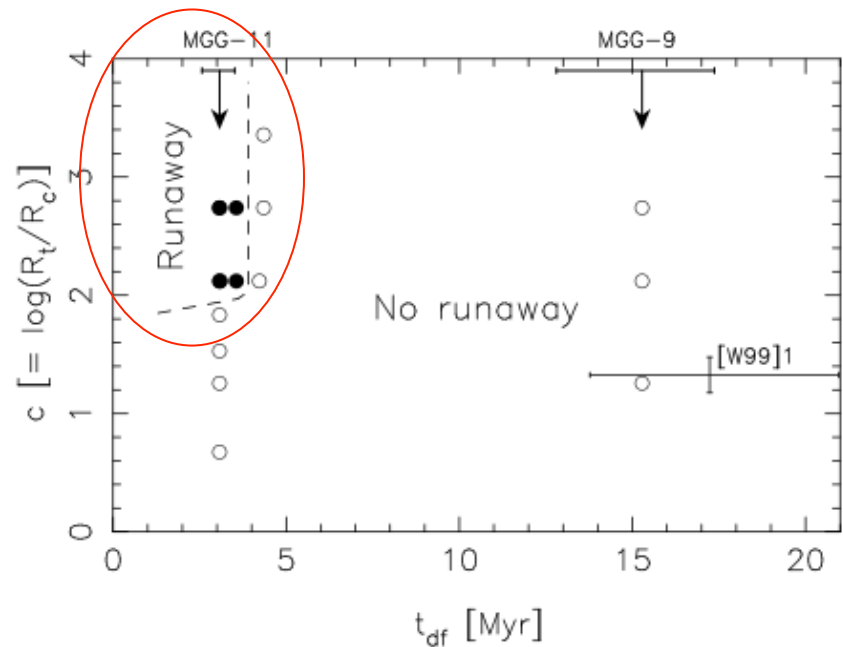
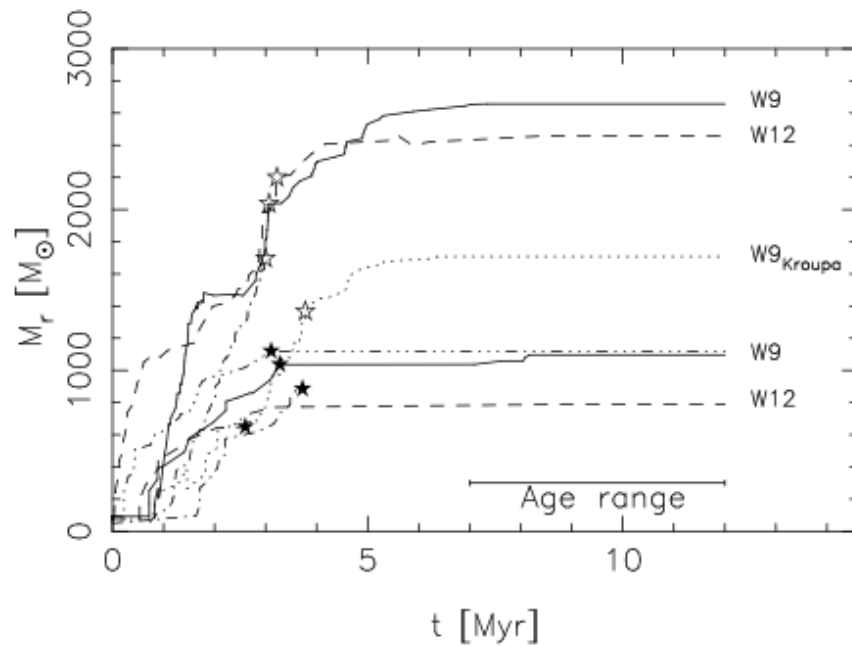
Star clusters cause core collapse by 2-body interaction



Very Massive Stars in compact star clusters

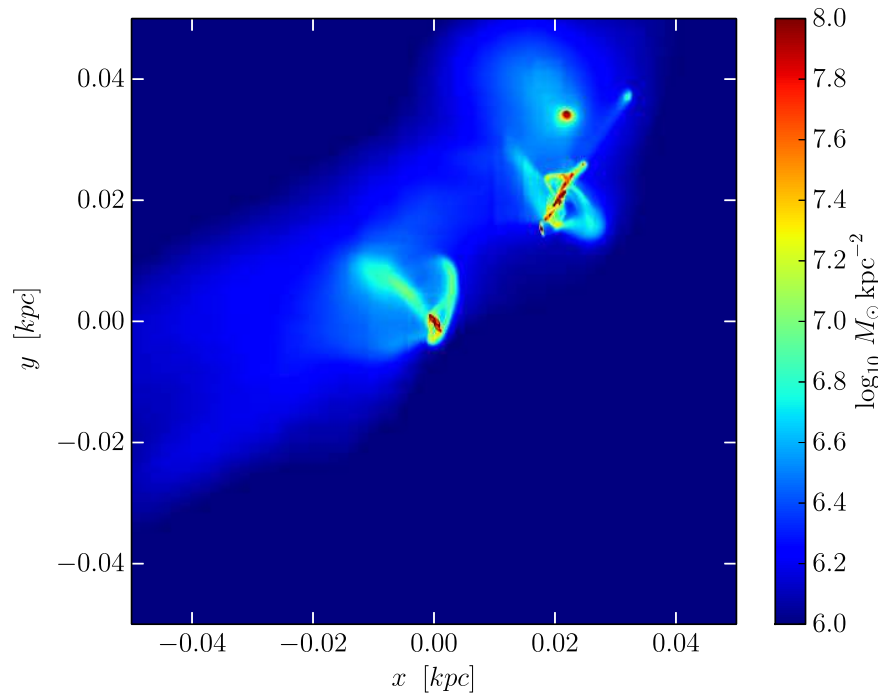
Portegies Zwart et al. (2004, Nature)

N-body simulations
 Compact star clusters \rightarrow core-collapse \rightarrow
 merging of stars \rightarrow very massive stars



Recent cosmological simulation

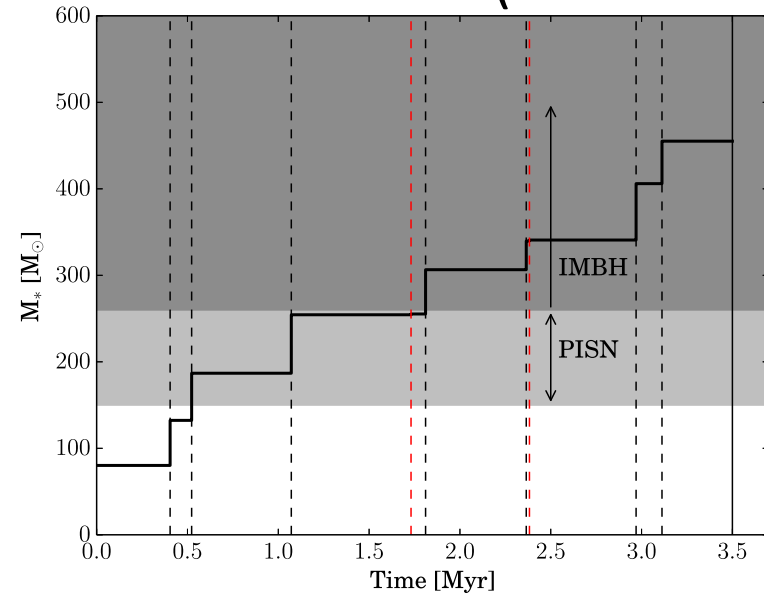
Katz+(2015)



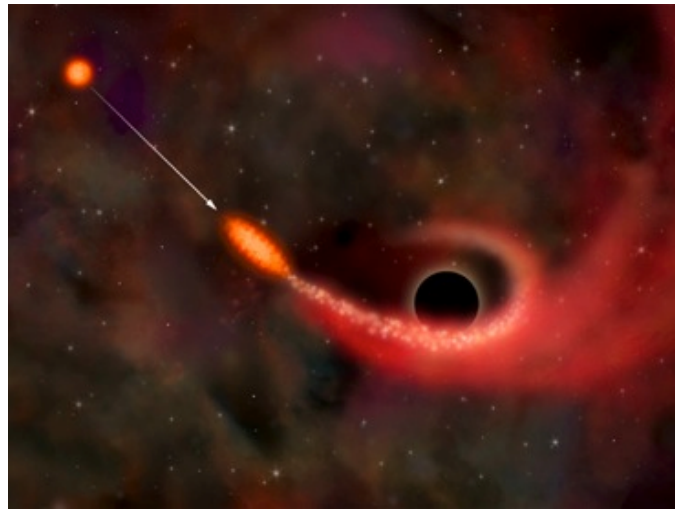
Mini-halos merger at $z \gtrsim 20$

Cosmological AMR simulation (RAMSES)
+
N-body simulation

Formation of VMS ($\sim 450 M_{\odot}$)



Tidal disruption & Loss cone depletion

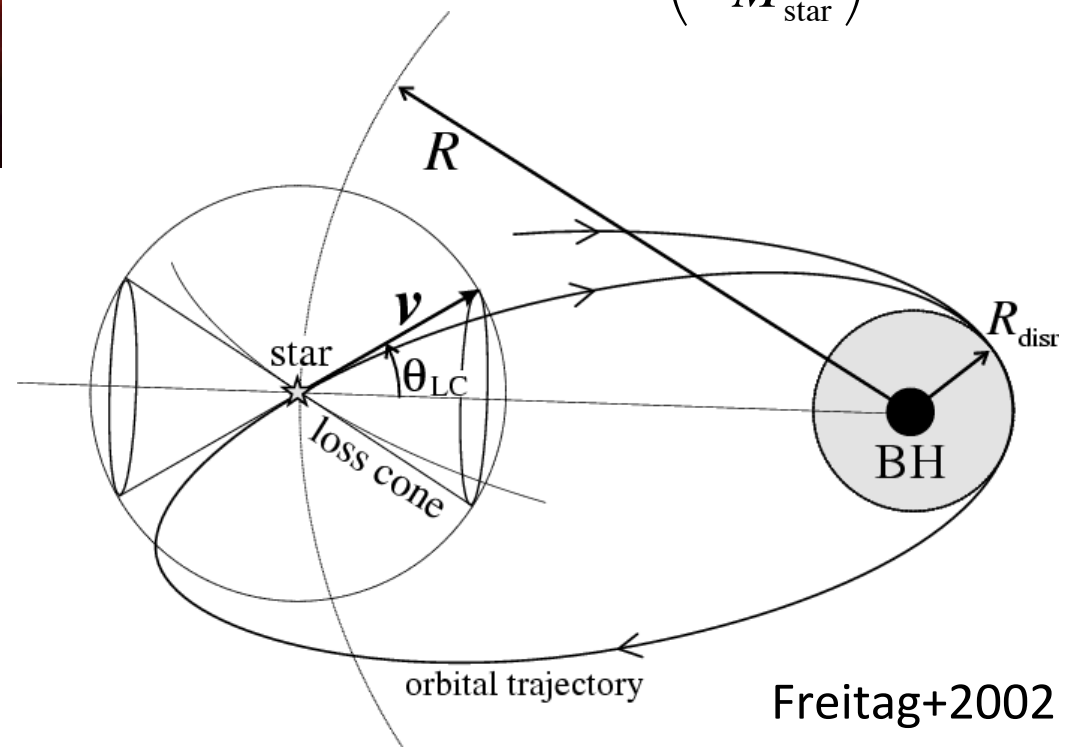


Stars passing within tidal radius of BHs are destroyed and accreted onto the BHs

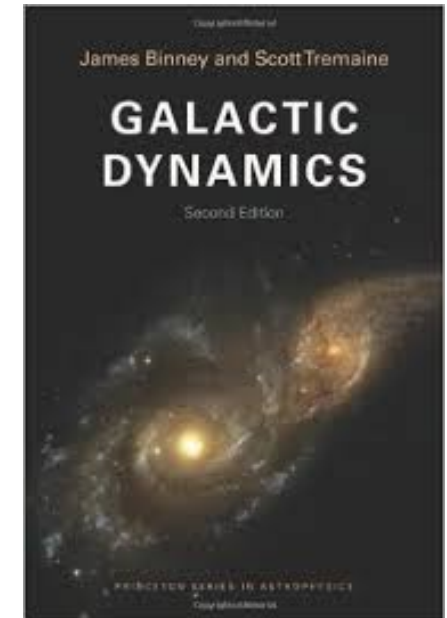
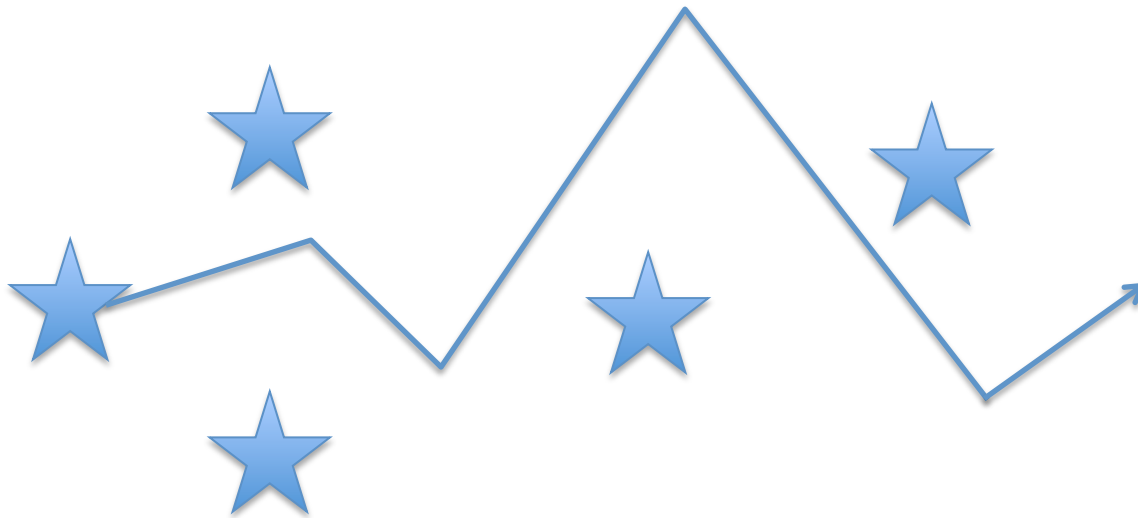
$$R_{\text{disr}} \cong \left(2 \frac{M_{\text{BH}}}{M_{\text{star}}} \right)^{1/3} R_{\text{star}}$$

Consequently, loss cone is made

$$\theta_{\text{LC}}^2 \cong 2 \frac{GM_{\text{BH}} R_{\text{disr}}}{v^2 R^2}$$



2-body relaxation

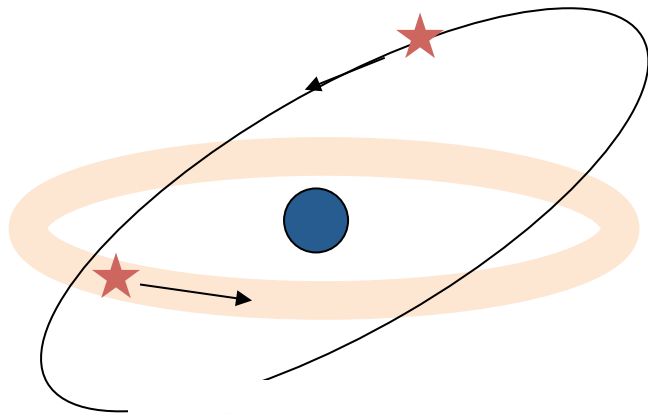


Time scale until stellar orbit is totally 90 degree changed by star-star collisions

$$\tau_{\text{NR}}(r) = 0.34 \frac{\sigma^3}{G^2 \rho m_2 \ln(\Lambda)}$$

Resonant Relaxation (RR)

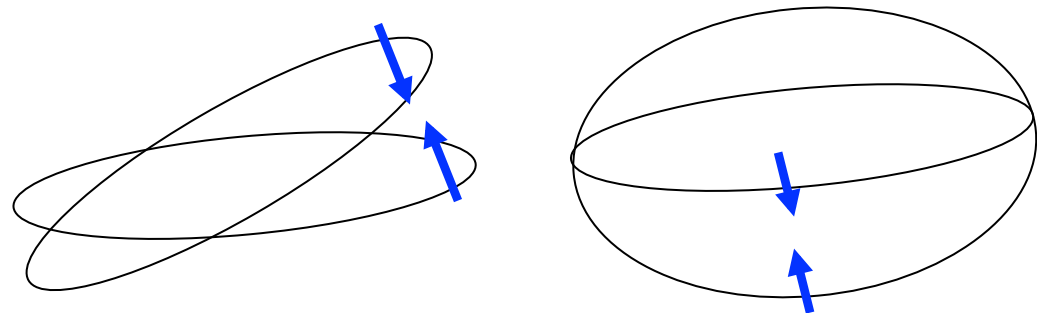
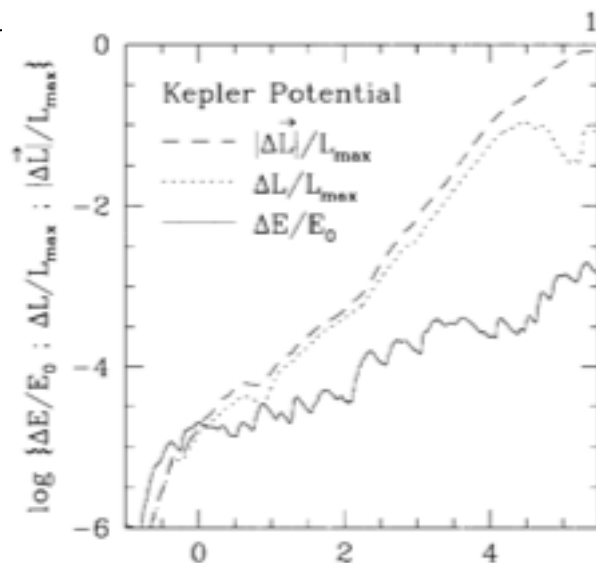
Rauch & Tremaine (1996)



Star in Kepler motion can give a torque to other stars for a time $t < t_{\text{prec}}$

$$\dot{J} \sim \frac{\sqrt{N} G m}{R^2} \times R$$

$$\frac{\Delta J}{J_c} = \frac{\dot{J} t}{J_c} \sim \sqrt{N} \left(\frac{M_{\text{star}}}{M_{\bullet}} \right) \frac{t}{2\pi \sqrt{R^3 / GM_{\bullet}}}$$



Model: Disk model & star cluster

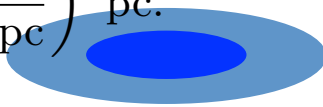
1: Disk surface density

Isothermal Mestel disk

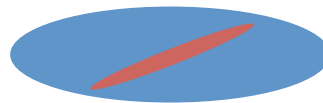
$$\Sigma(r) = \Sigma_0 \frac{R_d}{r} \quad (\text{Mestel 1963})$$

$$\Sigma_0 = 70 \left(\frac{V_h}{15 \text{ km s}^{-1}} \right) M_\odot \text{ pc}^{-2},$$

$$R_d = 100 \left(\frac{\lambda}{0.05} \right) \left(\frac{R_{\text{vir}}}{700 \text{ pc}} \right) \text{ pc.}$$



2: Gas inflow by bar instability after Major merging



$$M_{\text{Inf}} = 2\pi\Sigma_0 \frac{1}{v_\phi} \frac{GM_{\text{bar}}}{R_0} \Delta\tau$$

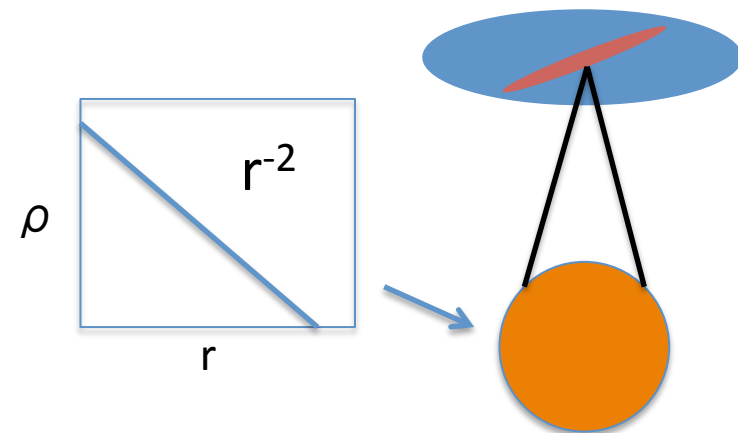
$$\Delta\tau \sim 3\text{Myr}$$

(e.g., Hopkins et al. 2009)

3: Spherical density structure at galactic centers with a singular isothermal solution

$$\rho_{\text{gas}} = \frac{c_s^2}{2\pi G r^2}$$

$$r_{\text{cl}} = 0.22 \text{ pc} \left(\frac{\sigma_s}{10 \text{ km s}^{-1}} \right)^{-2} \left(\frac{M_{\text{inf}}}{10^5 M_\odot} \right)$$

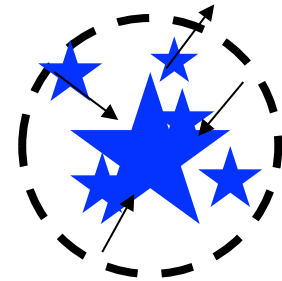


Model: Growth of BH

4: Core-collapse timescale

$$t_{\text{CC}} \sim t_{\text{df}} = \frac{1.91}{\ln \Lambda} \frac{r^2 \sigma_{3\text{D}}}{G m_{\text{max}}}$$

(Fujii & Portegies Zwart 2014)

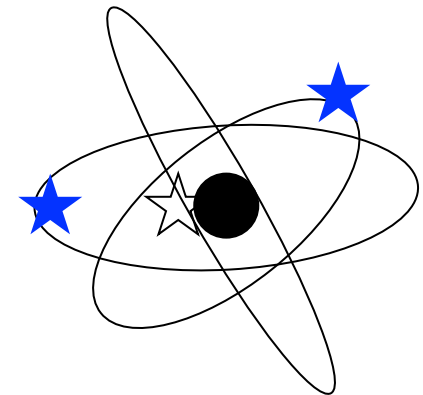


5: Growth of massive mass

$$M_{\text{VMS}} = m_{\text{init}} + \quad (\text{Portegies Zwart \& McMillan 2002})$$

$$4 \times 10^{-3} M_{\text{cluster}} f_c \ln \lambda_c \ln \left(\frac{3 \text{ Myr}}{t_{\text{CC}}} \right)$$

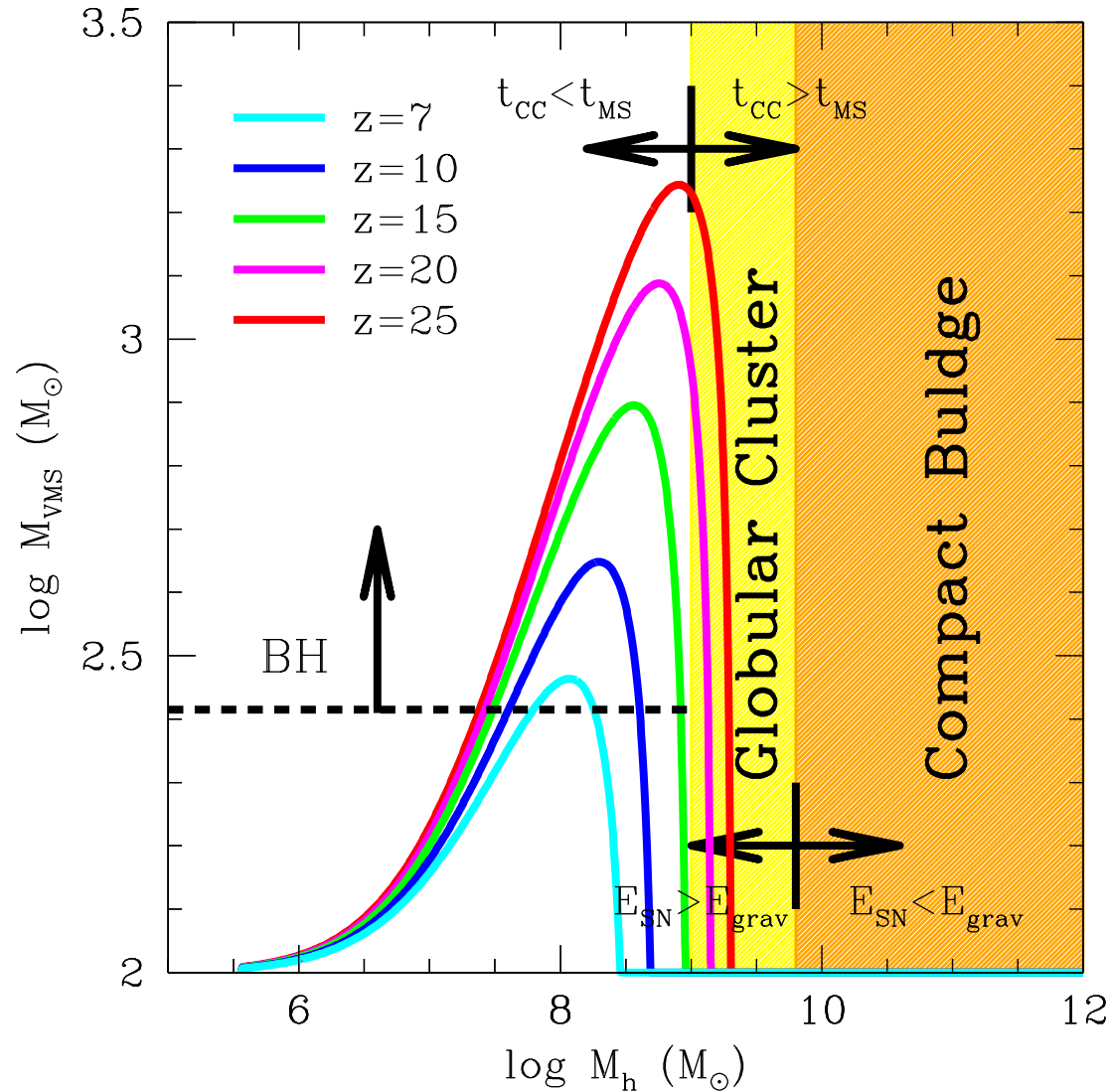
6: Initial Mass Function: Salpeter IMF with $0.1-100 M_{\text{sun}}$



7: Growth of BH

$$M_{\text{BH}}(t) = M_{\text{BH}}^{\text{init}} + \int_0^t dt \int_0^{R_{\text{cl}}} \frac{4\pi r^2 \rho_{\text{star}}(r, t)}{t_{\text{relax}}(r, t)} dr.$$

Very Massive Stars

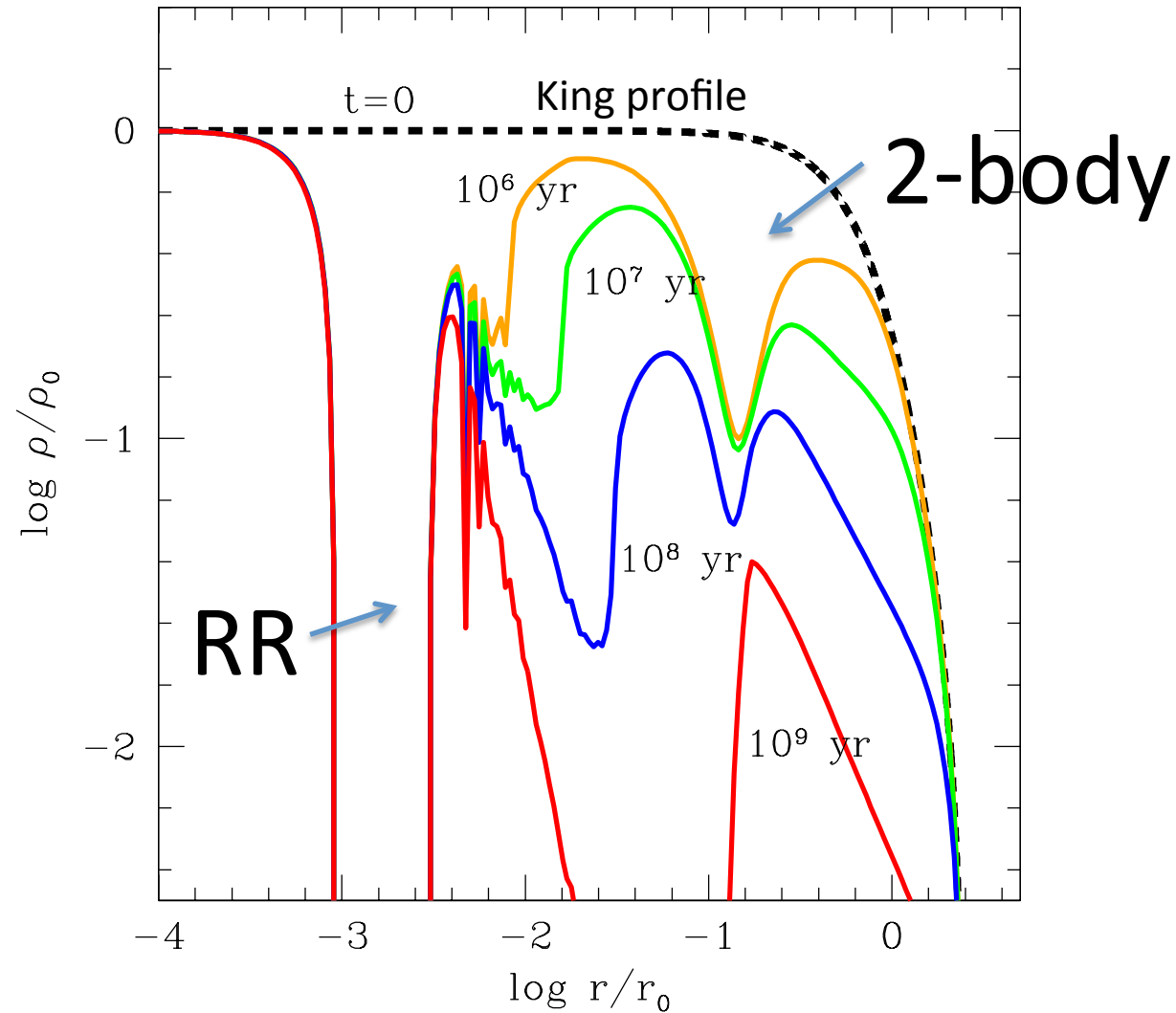


$$E_{\text{SN}} \sim 0.007 \times M_{\text{cl}} \times 10^{51} \text{erg}$$

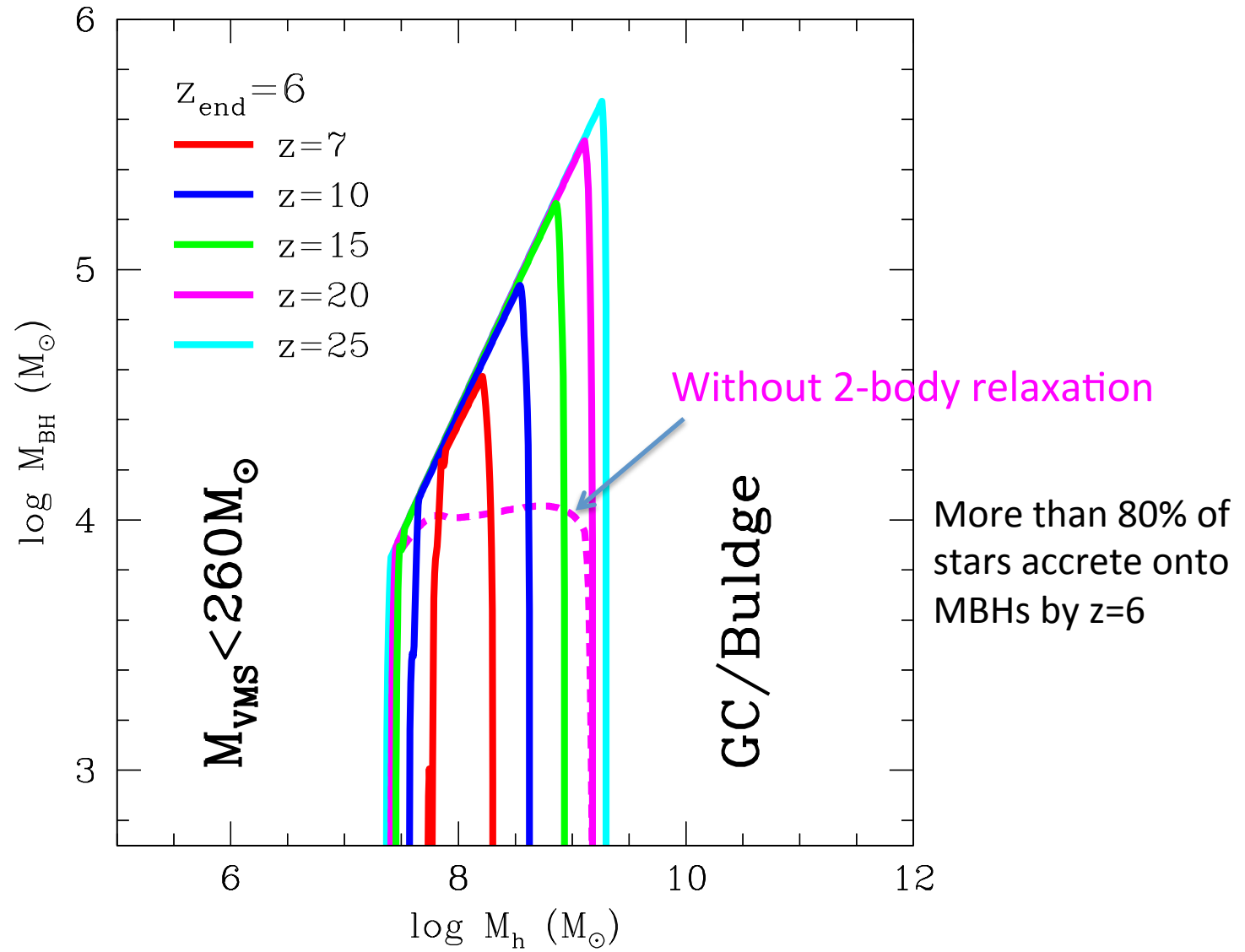
$$E_{\text{grav}} \sim \lambda \frac{GM_h^2}{R_{\text{vir}}}$$

Mass of VMSs decreases with redshift due to the bigger size of clusters

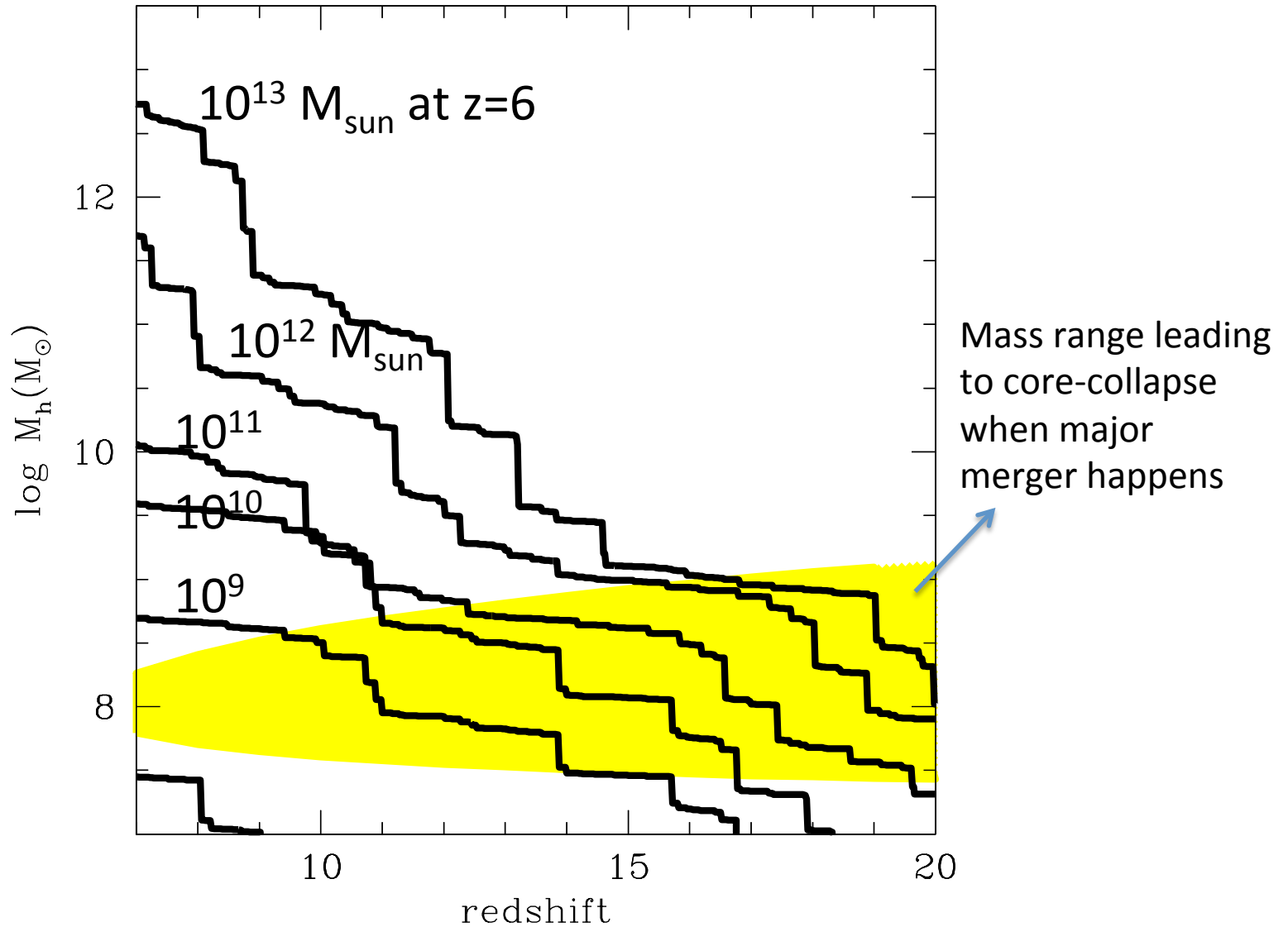
Stellar density distribution



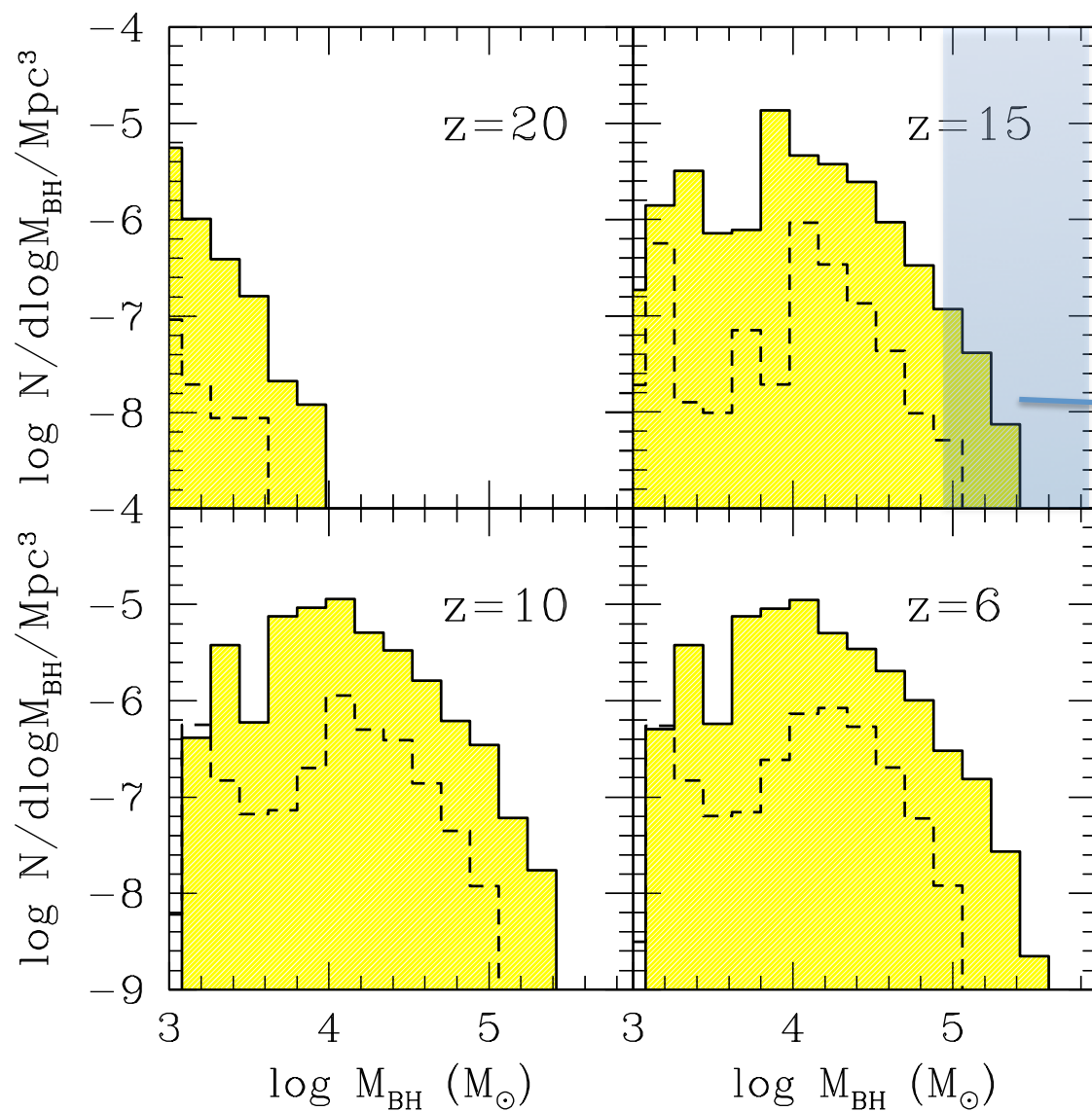
Black hole mass



Halo Merger Tree



BH mass function

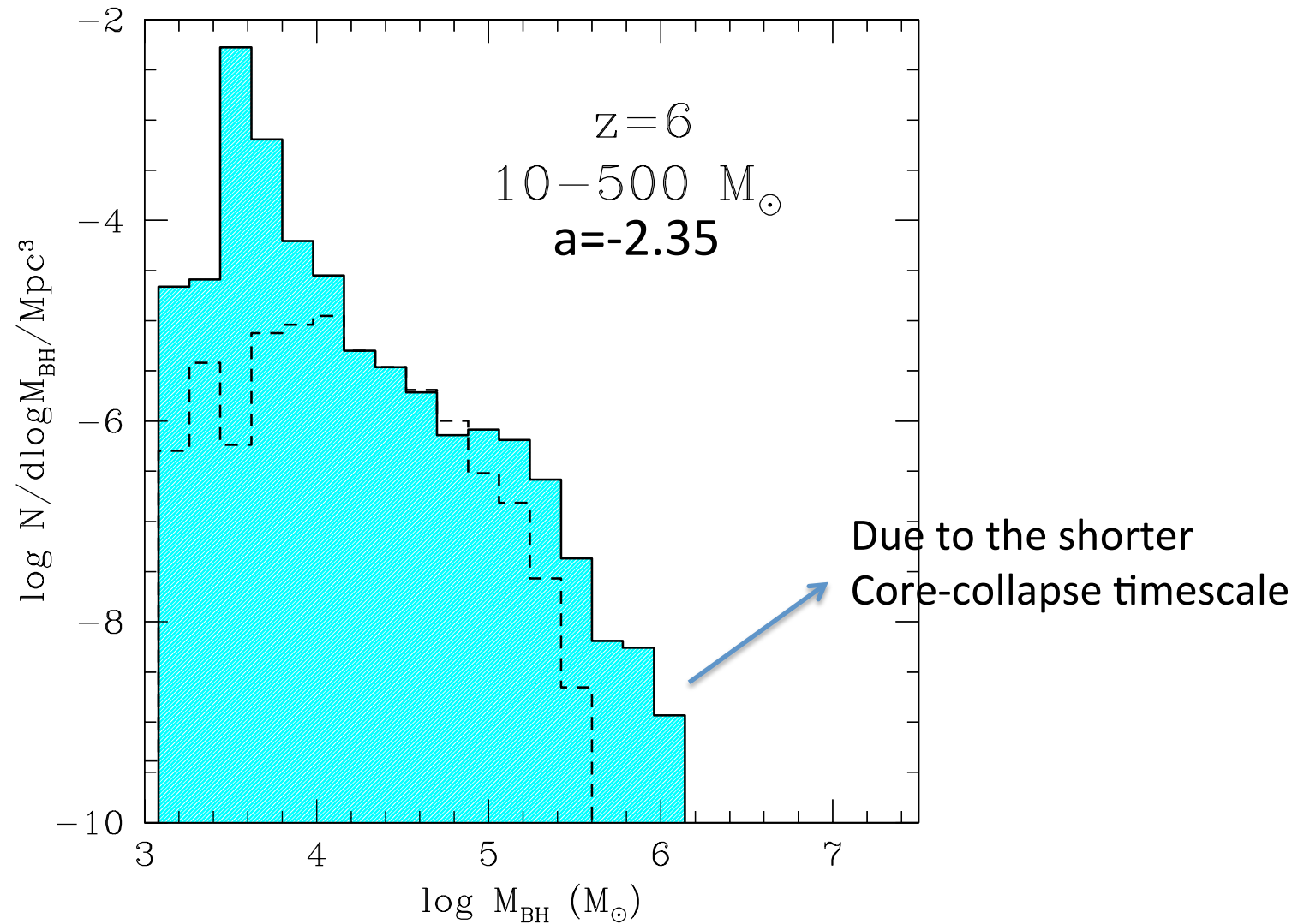


SMBH at $z\sim 6-7$

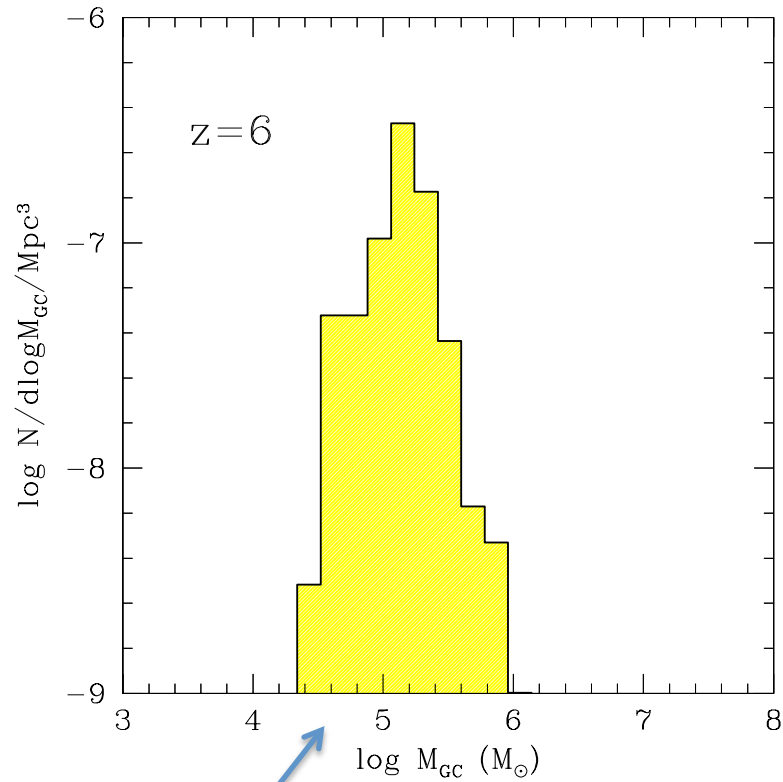
$$t = \frac{\epsilon}{1 - \epsilon} \tau_{\text{Sal}} \ln 10^4 \sim 0.46 \text{ Gyr}$$

Merger Tree
+
Seth-Tormen Halo
mass function

If Pop III star clusters

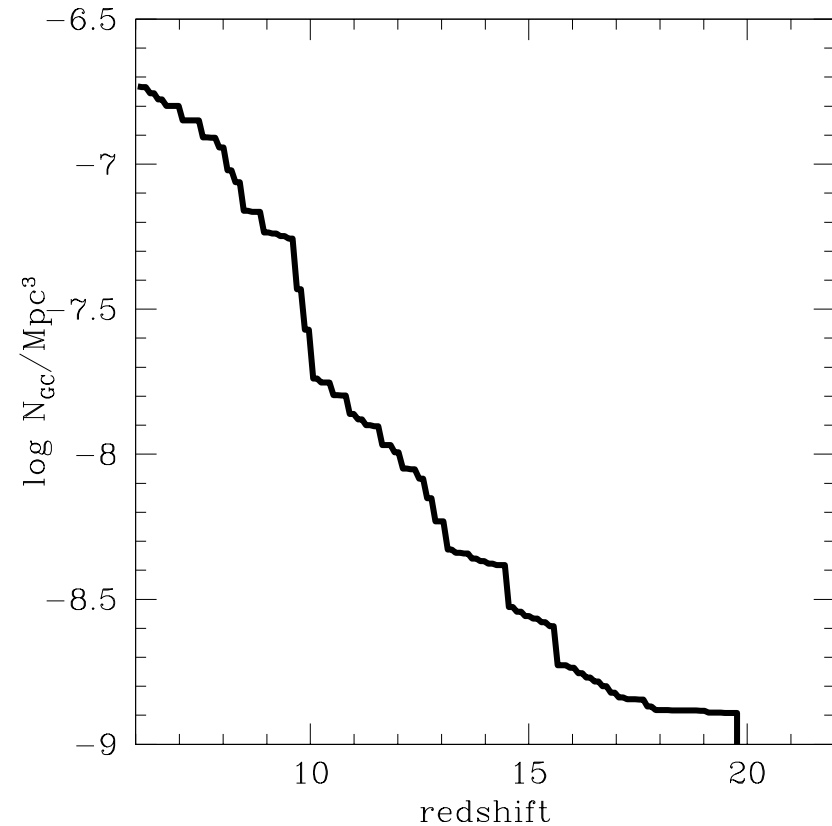


Globular Clusters



Mach the mass range of the observed local GCs

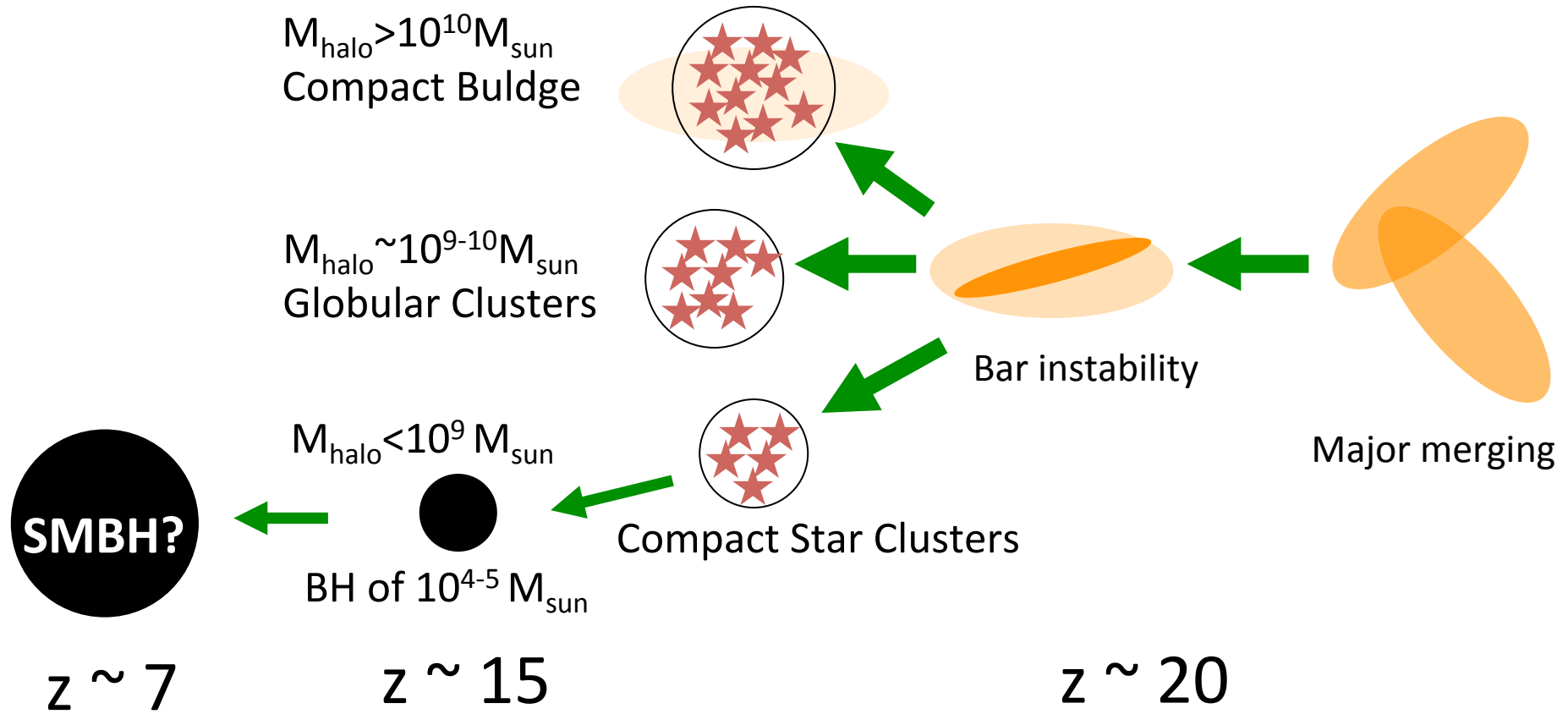
Number density of our modeled GCs
Is much smaller than that of local
observed GCs.



Summary

- We model the formation of SMBH seeds in merging first galaxies.
- Merging of galaxies with $M_h \sim 10^{8-9} M_{\text{sun}}$ leads to compact star clusters, resulting in very massive stars of $\sim 1000 M_{\text{sun}}$ after the core-collapse
- Massive BHs can grow up to $\sim 10^5 M_{\text{sun}}$ via stellar relaxation processes (2-body and resonant relaxation of stars)

Schematic view of our model



END