

How Big Can Supermassive BH Grow?

Omar López-Cruz

omarlx@inaoep.mx

Instituto Nacional de Astrofísica, Óptica y Electrónica (INAOE)
Sta. Ma. Tonantzintla, Puebla, México.

GH 2015: "Formation and Fueling of Supermassive Black Hole Seeds"
July 18, 2015

Collaborators

- Christopher Añorve (FACITE, Universidad de Sinaloa, México) First Licenciatura en Astronomía in México
- Mark Birkinshaw (Astrophysics, U Bristol, UK)
- Diana M. Worrall (Astrophysics, U Bristol, UK)
- Héctor J. Ibarra-Medel (INAOE, México, now at IA UNAM)
- Wayne A. Barkhouse (DPA, University of North Dakota, USA)
- Juan Pablo Torres-Papaqui (DAUG, México)
- Verónica Motta (DFA, U Valparaíso, Chile)

Outline

- I will try to persuade you that the known BH scaling law breaks down for cD galaxies (a special kind of BCG).
- And that the final BH masses might have been set by initial conditions.

$$1 \text{ kpc} = 3.08521 \times 10^{21} \text{ cm}$$

DGCG: Driver for GALFIT on Cluster Galaxies

GALFIT fits two-dimensional surface brightness to galaxies on digital images.

GALFIT functions: Sérsic, exponential, Gaussian, King Profile, Nuker, PSF profiles.

DGCG adapts GALFIT (Peng et al. 2002) to cluster galaxies.

GALFIT →



Añorve 2012, Ph.D.
Thesis, INAOE.

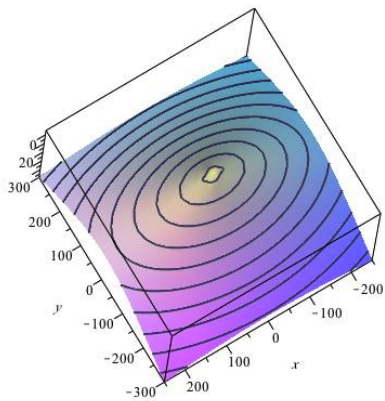
Surface Brightness Models

E galaxies and bulges of S and S0 galaxies are well described by the Sérsic function **Sérsic (1963)**:

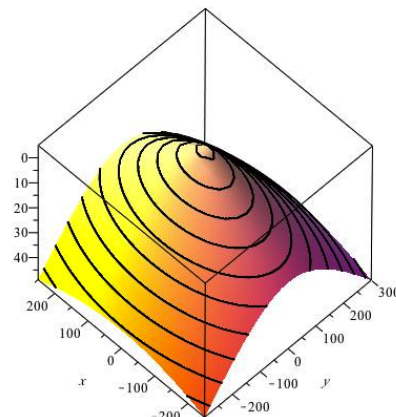
$$I(R) = I_e \exp \left(-k \left[\left(\frac{R}{R_e} \right)^{1/n} - 1 \right] \right) . \quad (1)$$

Disks of S and S0 galaxies are well fitted by the exponential function **de Vaucouleurs (1953)**

$$I(R) = I_0 \exp \left(-\frac{R}{R_s} \right) \quad (2)$$

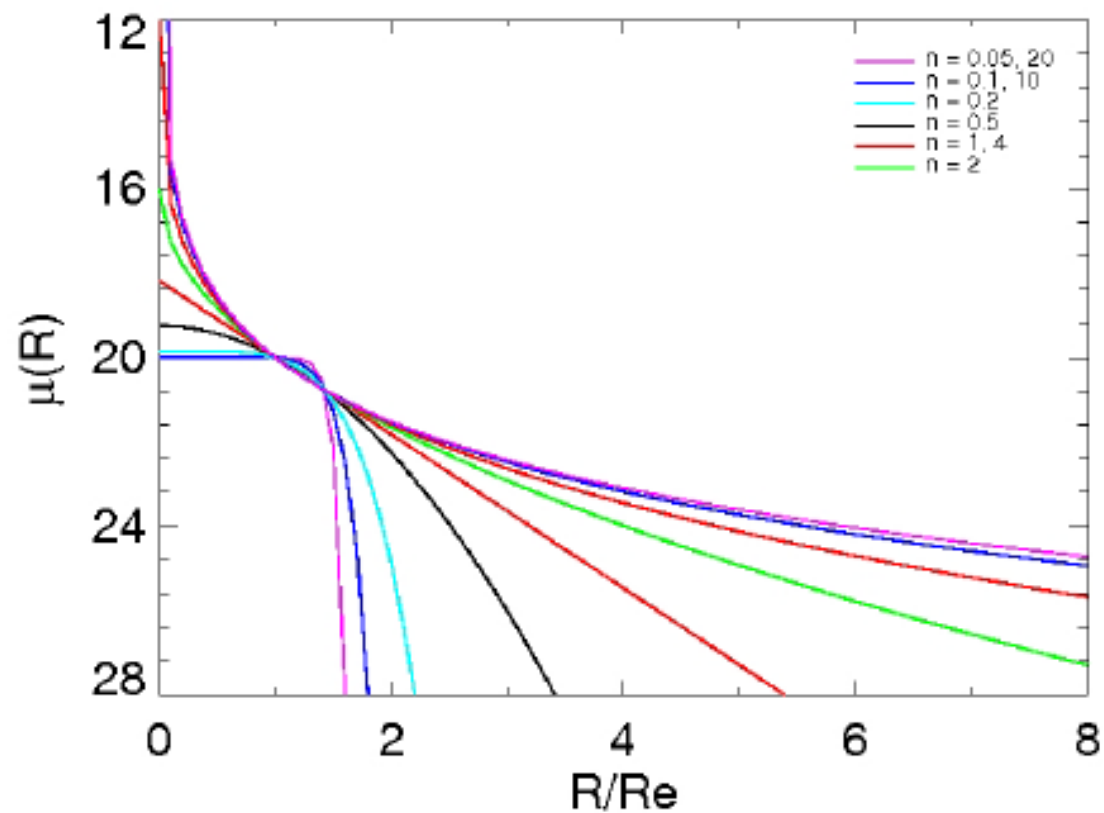


$$R = \left(x^{c+2} + \left(\frac{y}{q} \right)^{c+2} \right)^{\frac{1}{c+2}}$$



Surface Brightness Models

Strategy: Fit Sérsic (SS) and Sérsic + Exponential (BD) surface brightness to cluster galaxies

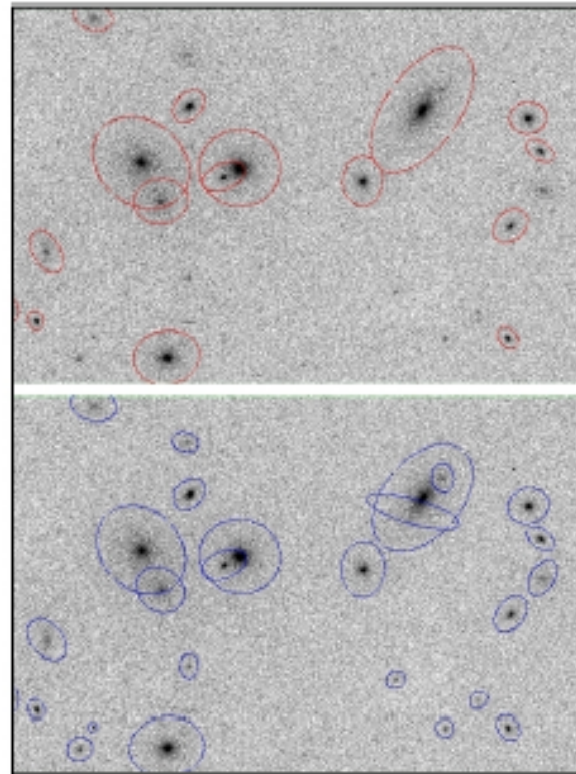


DGCG: Algorithm

- ① Initial parameters (SExtractor; Bertin & Arnouts 1996)
- ② Masking & Sky
- ③ Sorting
- ④ Fitting - Single, simultaneous
- ⑤ Bumpiness

DGCG: Algorithm

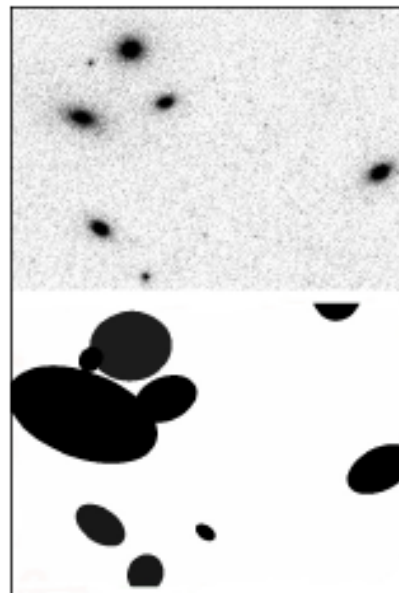
1.- Reads SExtractor catalog (Bertin & Arnouts 1996)



Star/Galaxy Classification (PPP)

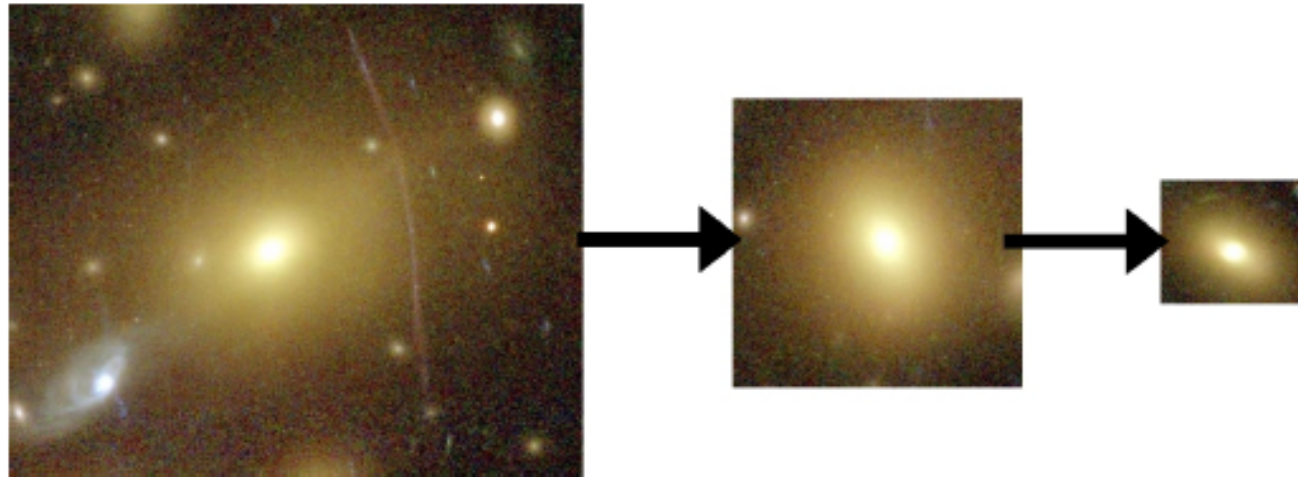
DGCG: Algorithm

2.- Creates Mask and computes Sky for every object on the catalog



DGCG: Algorithm

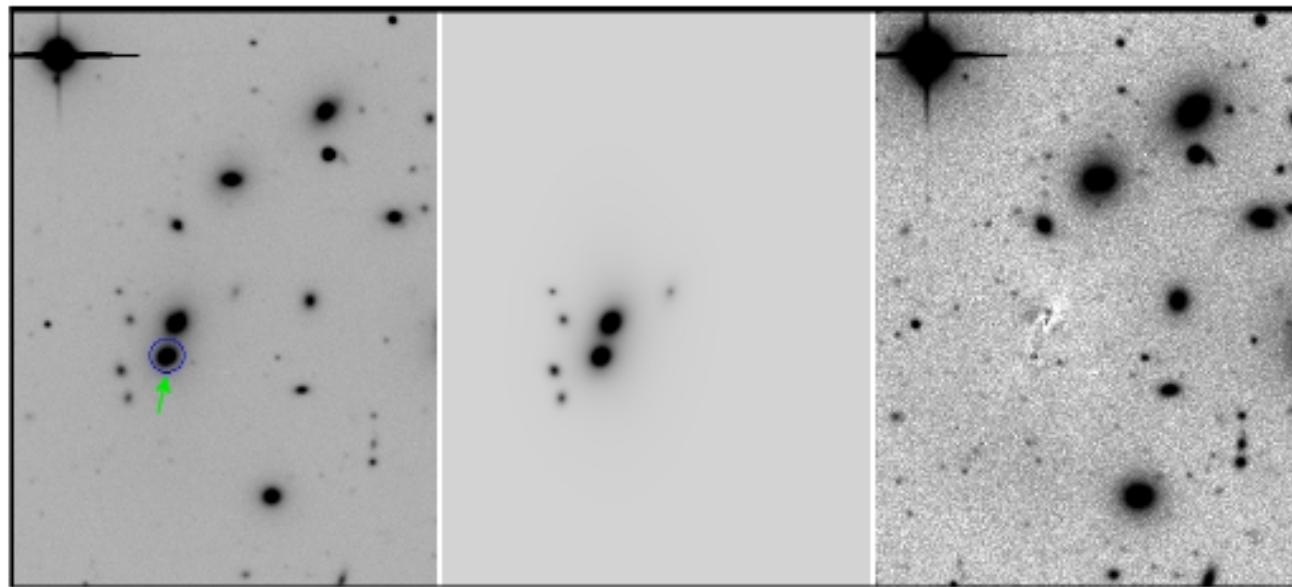
3.- Fit order: Brightest to faintest galaxy. Selects PSF. We used the closest star to the galaxy of interest.



DGCG: Algorithm

4.- Galaxy Fitting:

Simultaneous fitting for Neighbors galaxies. Use Mask for the rest.



DGCG: Algorithm

5.- Computes Bumpiness, Bulge to Total ratio, SNR, local χ^2_ν within Kron Radius

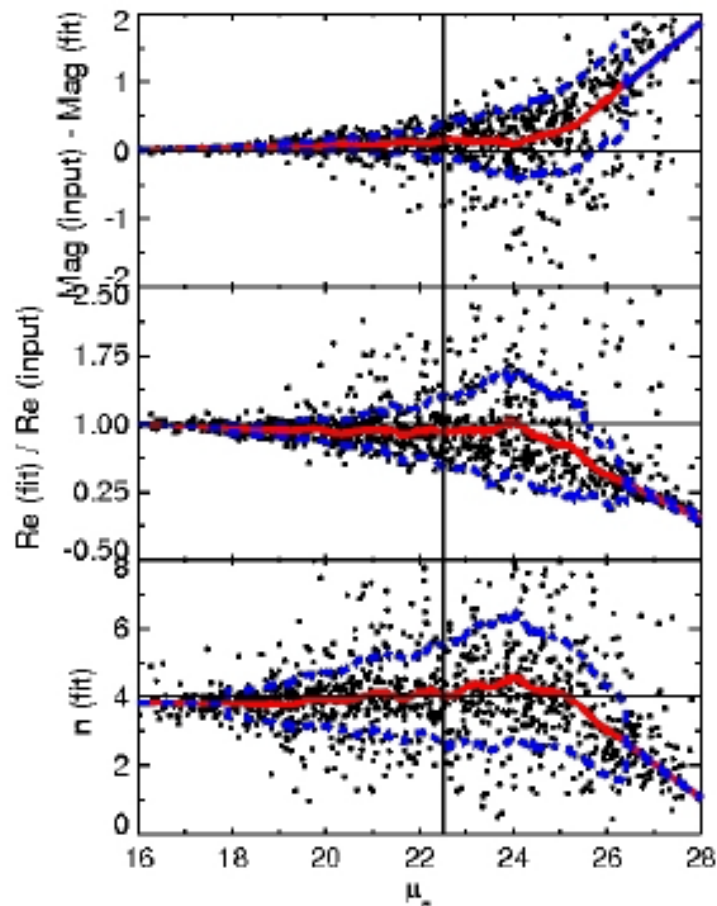
Bumpiness: Blakeslee et al. (2006)

$$BPN = 10 \frac{\sqrt{\langle [I - S(Re, n)]^2 \rangle - \langle \sigma_s^2 \rangle}}{\langle S(Re, n) \rangle} \quad (3)$$

Bars

Surface Brightness Photometry on Artificial Galaxies (Validation)

Fits on GEMS artificial galaxies (SS)

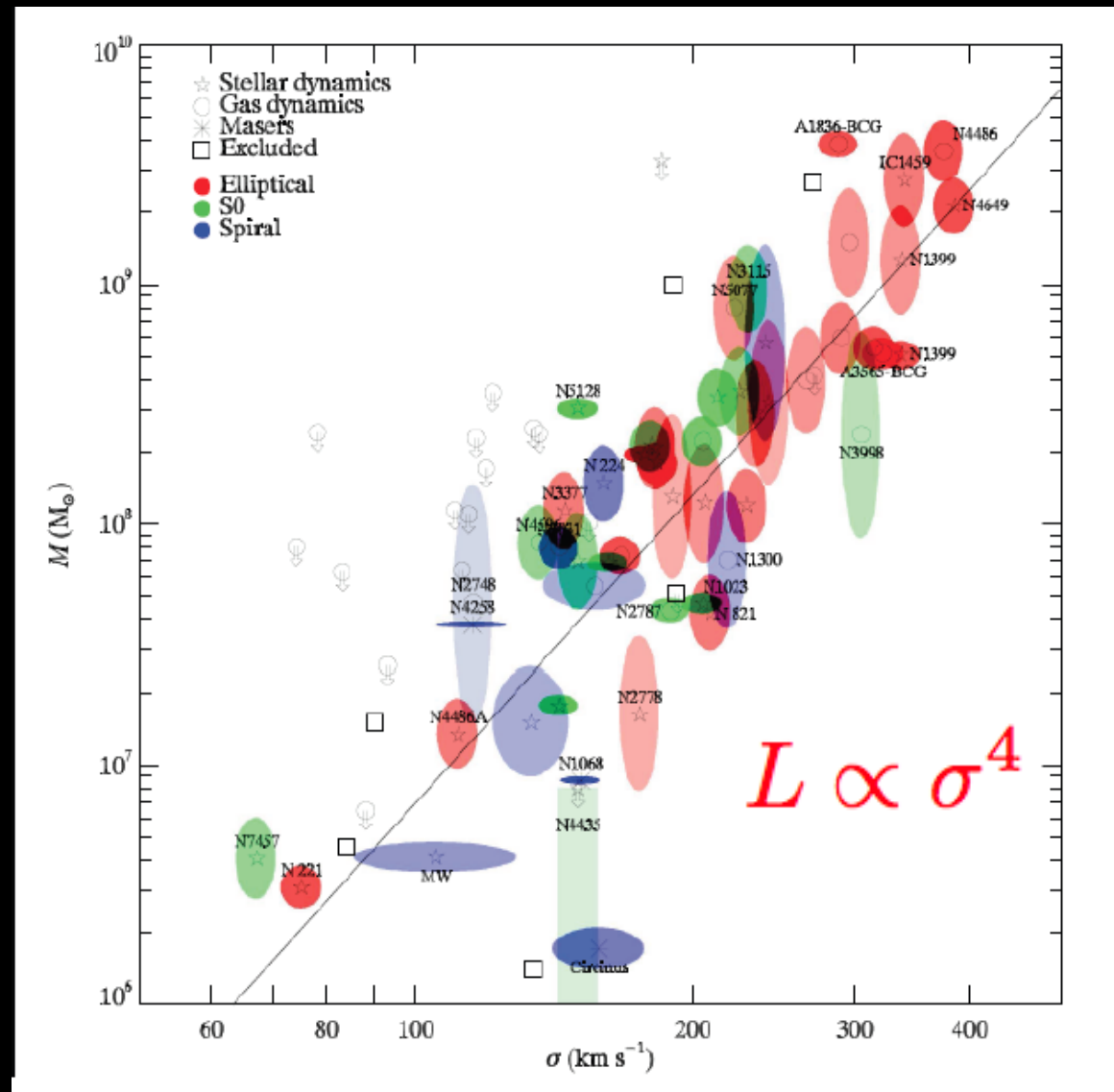


The Sérsic index
is hard to recover!

Results

- 1453 from 21 low- z ($0.02 \leq z \leq 0.07$) Abell clusters
- $R = 18$ mag completeness limit
- 1453 galaxies with both SS and BD acceptable fits
- The final sample contains 304 E, 557 S0 and 548 S galaxies
- 297 galaxies whose BD fits resulted in Sérsic index less than 0.2 ($n < 0.2$)
- It took 258.89 hrs (~ 10.79 days) for SS models and 337.23 hrs ~ 15.72 days for the BD models, respectively, to process the whole sample.

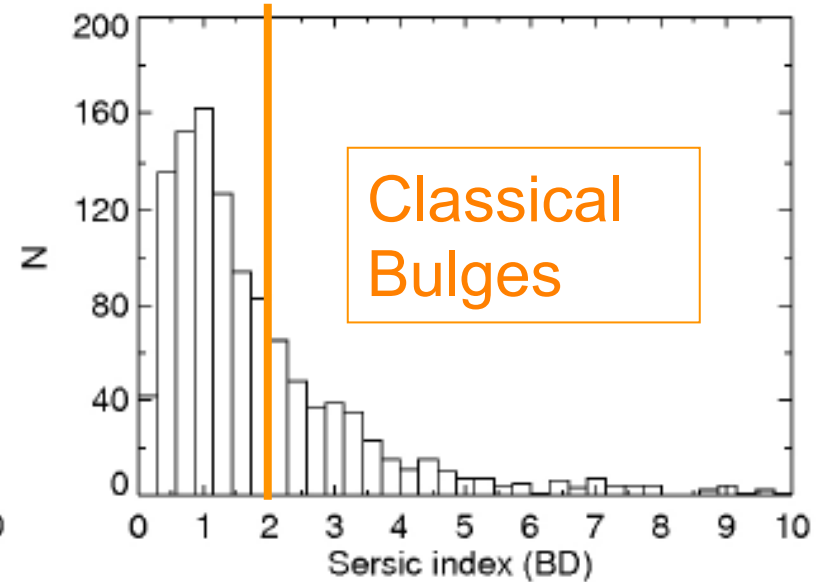
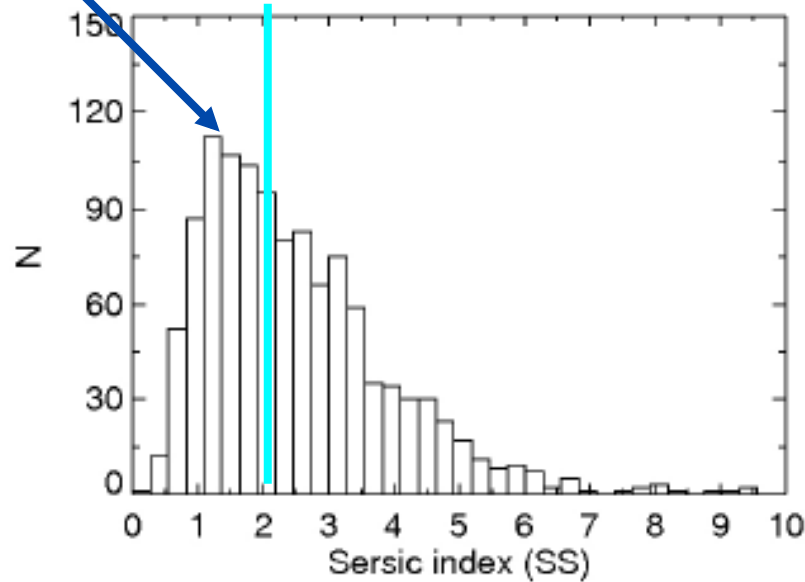
Black hole mass correlates with the velocity dispersion of their host spheroidal component.



Ultramassive BH : $M_{\bullet} \geq 10^{10} M_{\odot}$

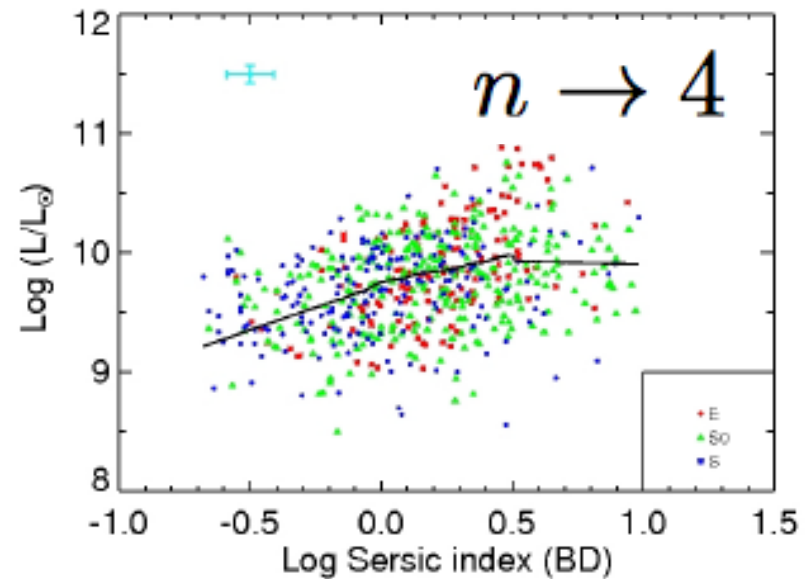
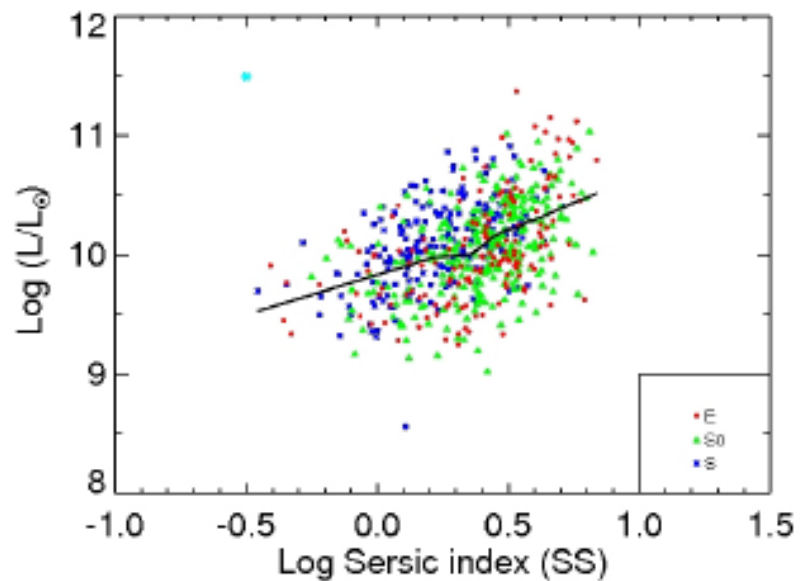
Distribution of Sérsic Index

Pseudo Bulges



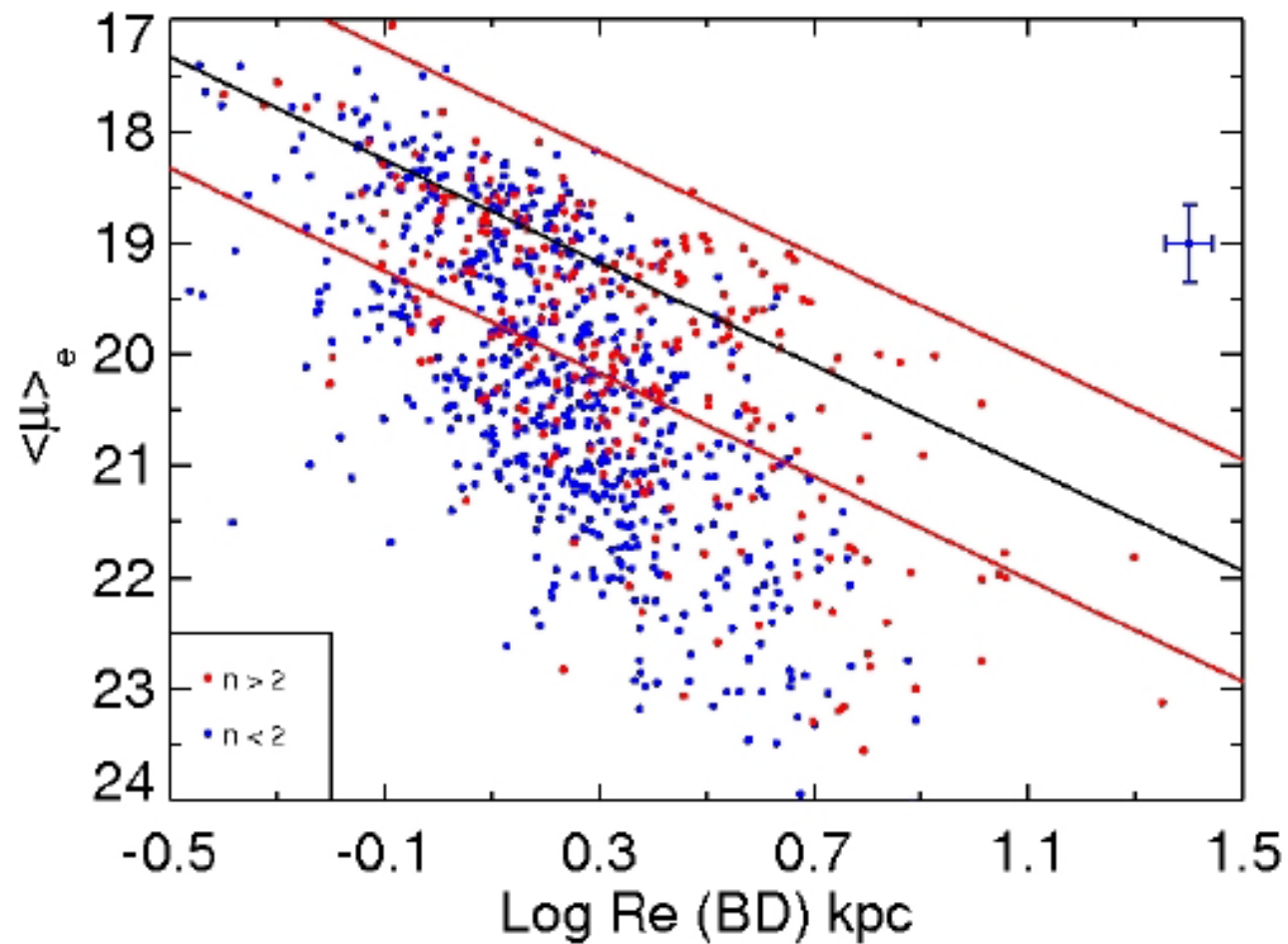
Blanton et al. (2003); Blanton & Moustakas (2009); Fisher & Drory (2010).

Luminosity vs. Sérsic index



More luminous galaxies have larger Sérsic Indexes ($n \rightarrow 4$)

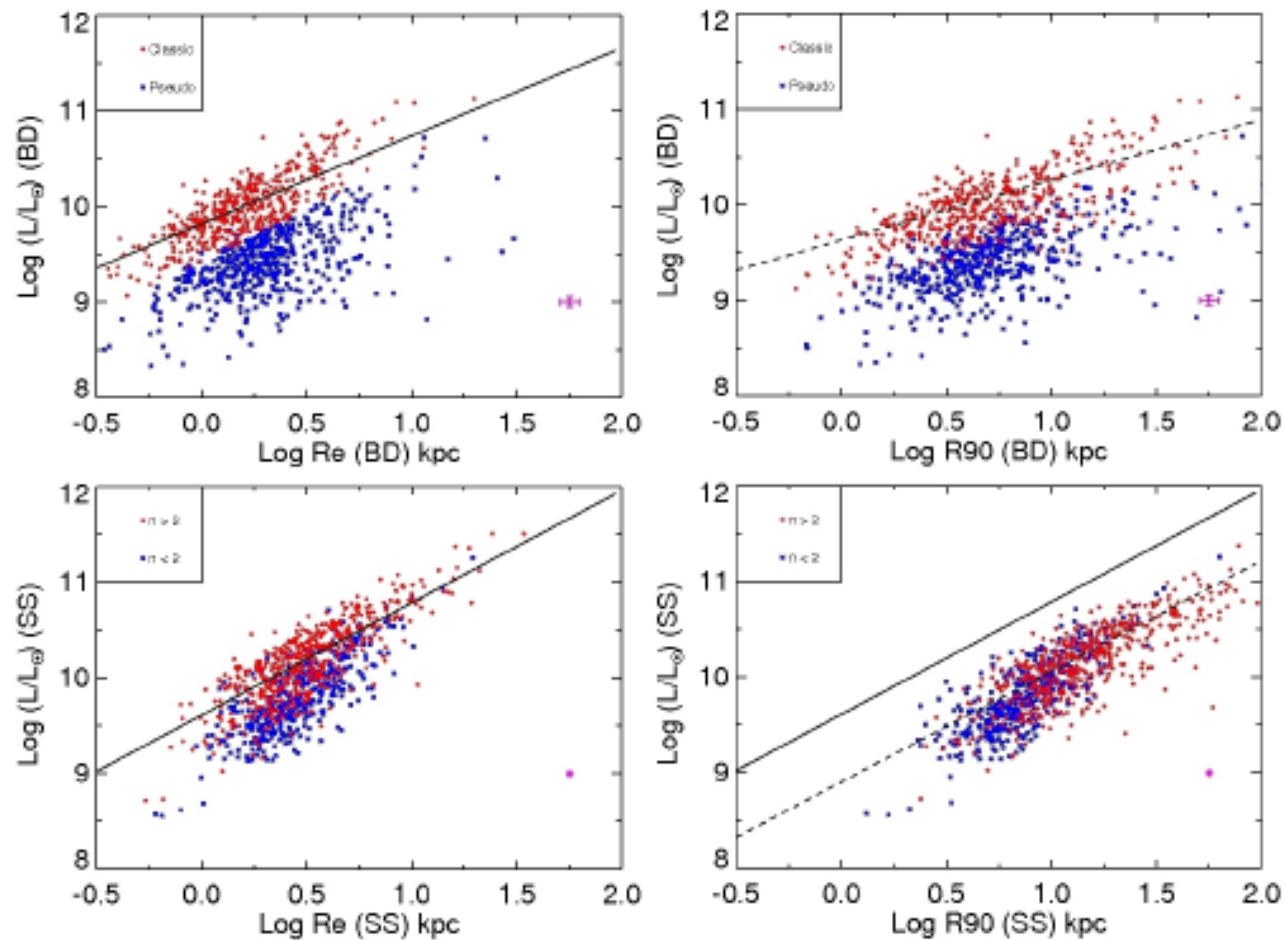
Kormendy Relation



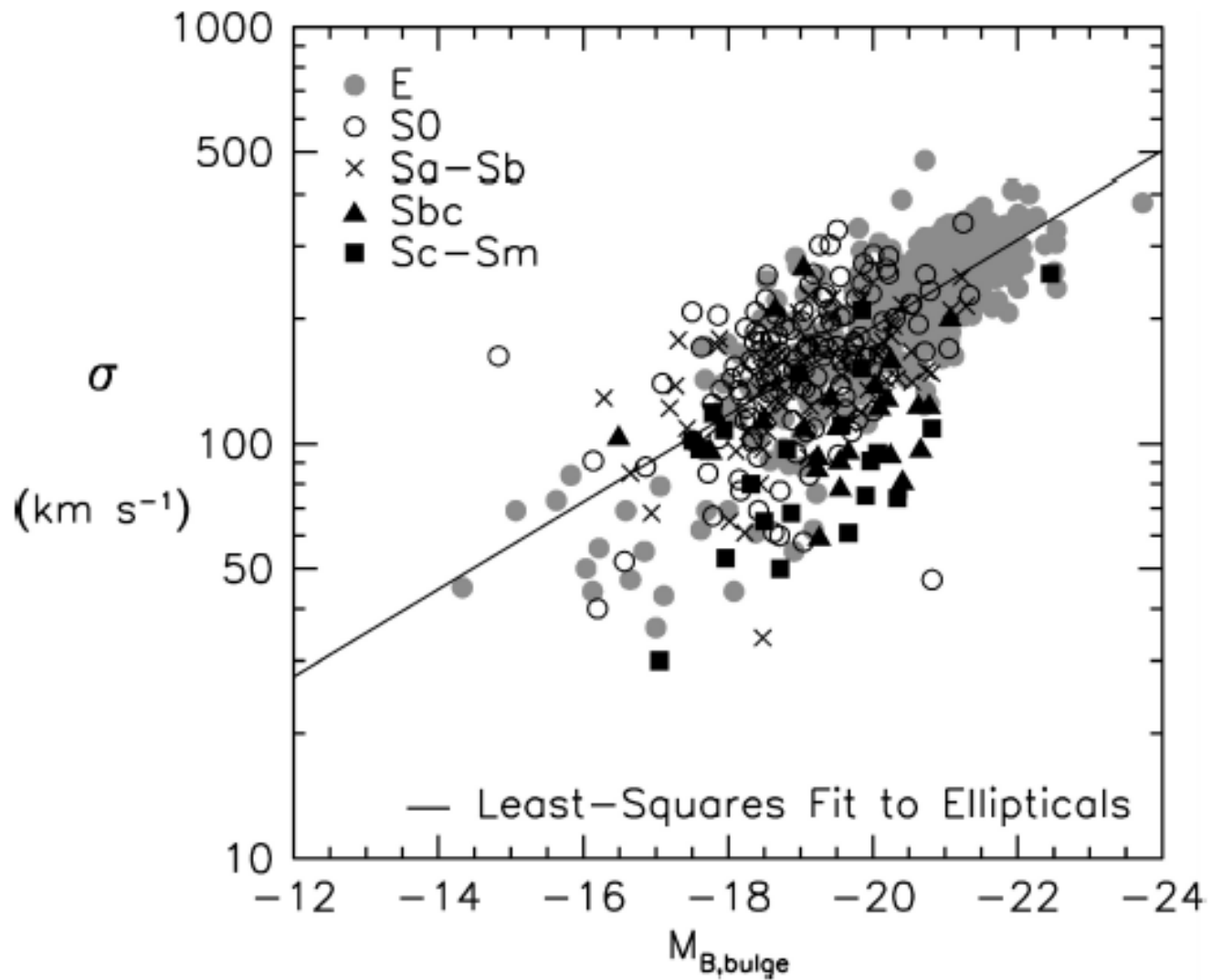
Gadotti (2009)

Pseudobulges: 65% $n < 2$
Diagnostic Diagrams

Luminosity-Size Relation

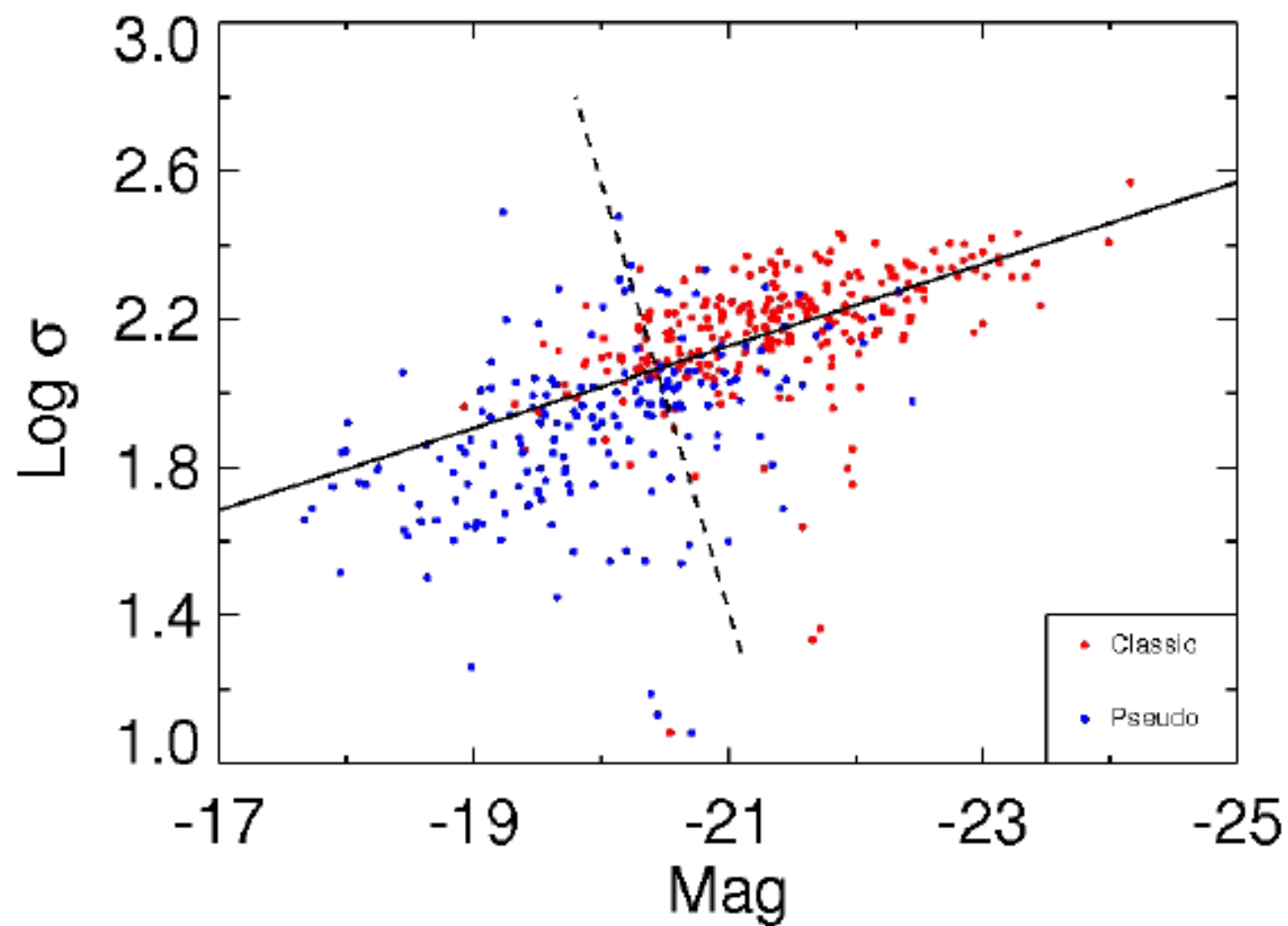


Top: Bulges, Bottom: Galaxy



Kormendy & Kennicutt (2004)

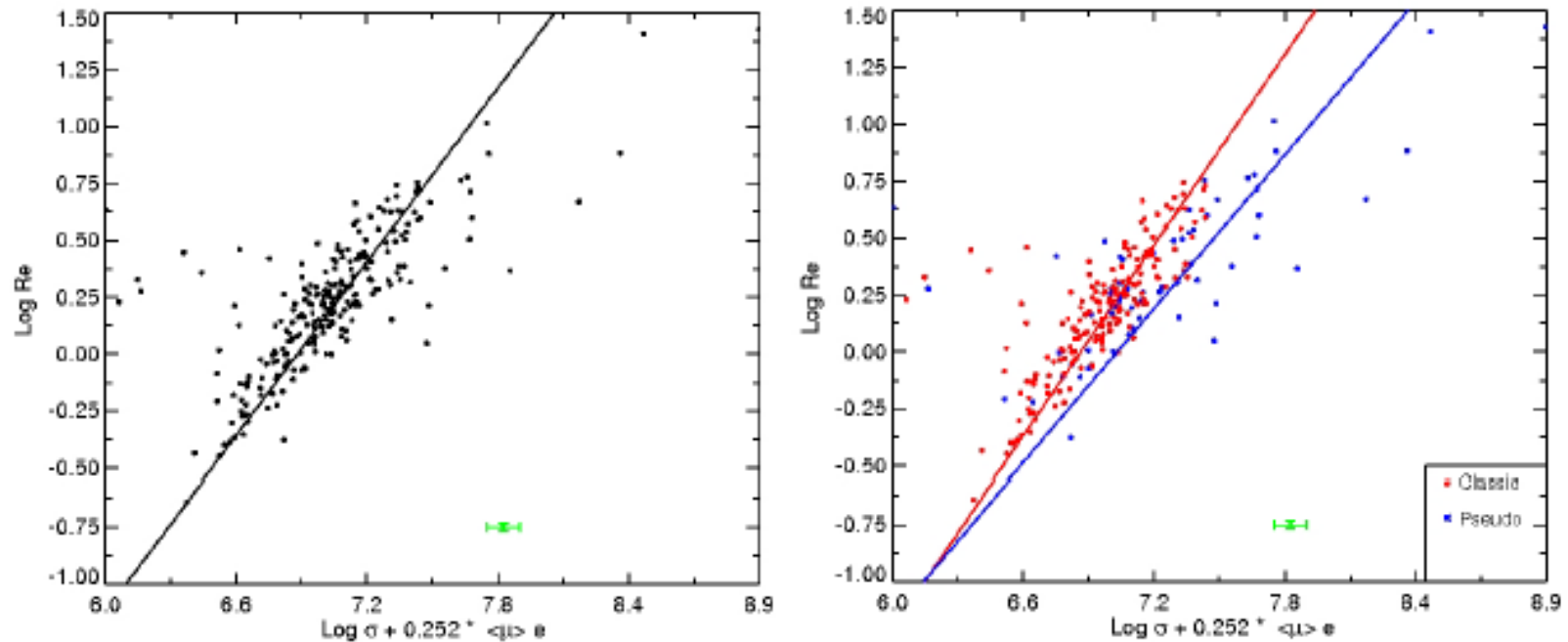
Faber-Jackson Relation



$$\log \sigma = -0.11 \text{ Mag} - 0.23,$$

(8)

Fundamental Plane



Normal bulges $\log Re = 1.43 \pm 0.023(\log \sigma + 0.252) \langle \mu \rangle_e + 9.8 \pm 0.16$,

Pseudobulges $\log Re = 1.16 \pm 0.039(\log \sigma + 0.252) \langle \mu \rangle_e + 8.1 \pm 0.077$

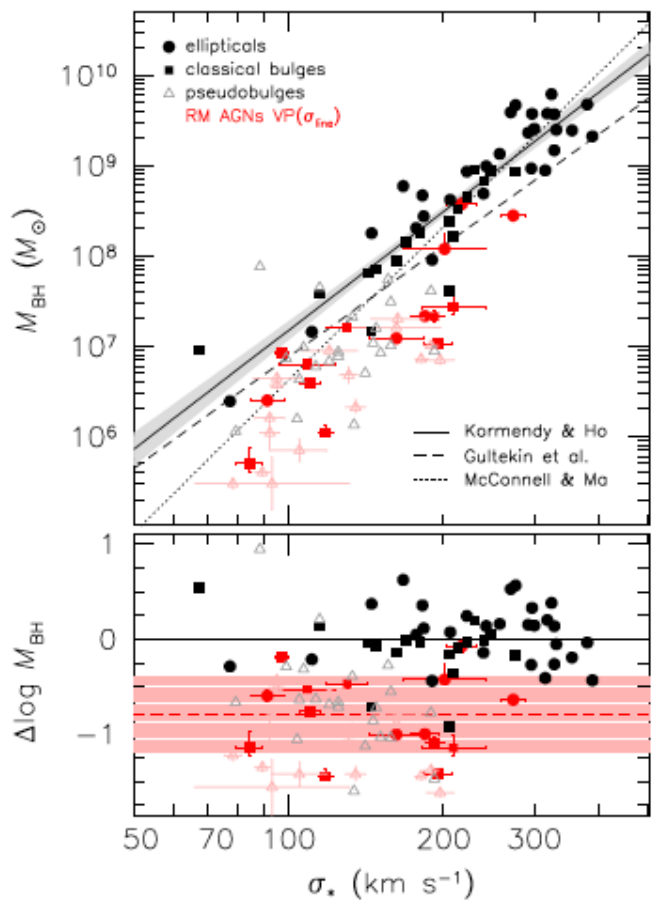


FIG. 2.— (Top) The $M_{\text{BH}} - \sigma_*$ relation for inactive galaxies (black points) and RM AGNs (red points). The masses for the RM AGNs represent the virial product, $\text{VP} = c\tau \Delta V^2 / G \equiv M_{\text{BH}} / f$, with $\Delta V = \sigma_{\text{line}}(\text{H}\beta)$ measured from rms spectra. Classical bulges and ellipticals are highlighted as filled symbols, and pseudobulges are plotted as open symbols. Error bars are suppressed for the inactive galaxies to reduce crowding. The best-fit relation of Kormendy & Ho for classical bulges and ellipticals (Equation 2) is given by the solid line; the gray shading represents its 1σ scatter. The fits of Gültekin et al. (2009) and McConnell & Ma (2013) are shown as dashed and dotted lines, respectively. (Bottom) Residuals of the data points with respect to the Kormendy & Ho fit for ellipticals and classical bulges. The red dashed line and associated shaded pink band mark the average offset and standard deviation for the 15 RM AGNs hosted by ellipticals and classical bulges: $\langle \Delta \log M_{\text{BH}} \rangle = -0.79 \pm 0.42$.

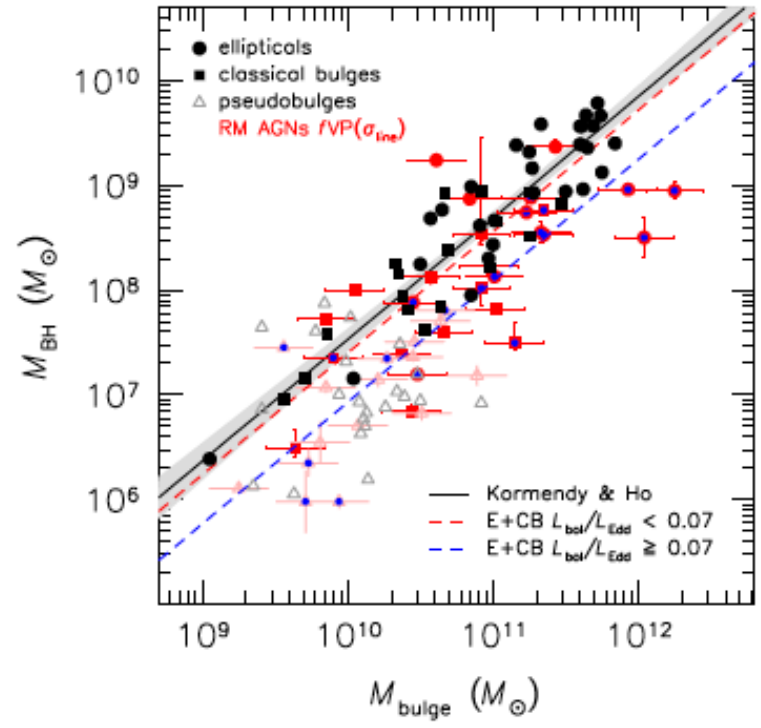


FIG. 5.— $M_{\text{BH}} - M_{\text{bulge}}$ relation for inactive galaxies (black points) and RM AGNs (red points). Error bars are suppressed for the inactive galaxies to reduce crowding. The virial products, derived from $\sigma_{\text{line}}(\text{H}\beta)$ measured from rms spectra, have been scaled by $f = 6.3$ for classical bulges and ellipticals and by $f = 3.2$ for pseudobulges. The best-fit relation of Kormendy & Ho for classical bulges and ellipticals (Equation 3) is given by the solid line; the gray shading denotes its 1σ scatter. AGNs hosted by classical bulges and ellipticals with $L_{\text{bol}}/L_{\text{Edd}} < 0.07$ are denoted by filled red symbols and the red dashed line; those with $L_{\text{bol}}/L_{\text{Edd}} \geq 0.07$ are highlighted with a blue center and the blue dotted line.

Ho & Kim (2014) bulge mass best correlates BH mass.

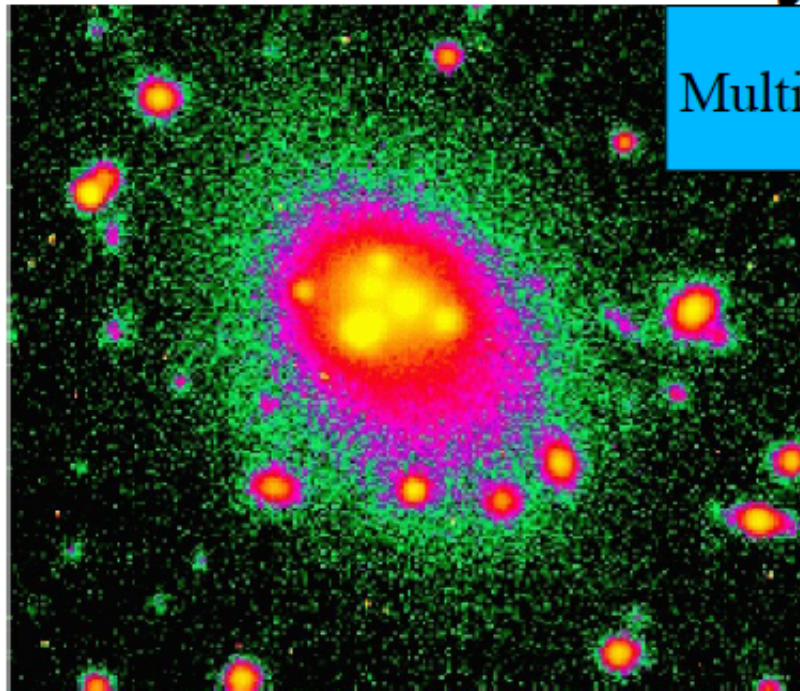
What are cD Galaxies?

- cD are supergiant galaxies up to 4 mags. Brighter than M^* . They can concentrate up to half the total cluster light.

$$L_{cD} = 5 \times 10^{12} L_{\odot}$$

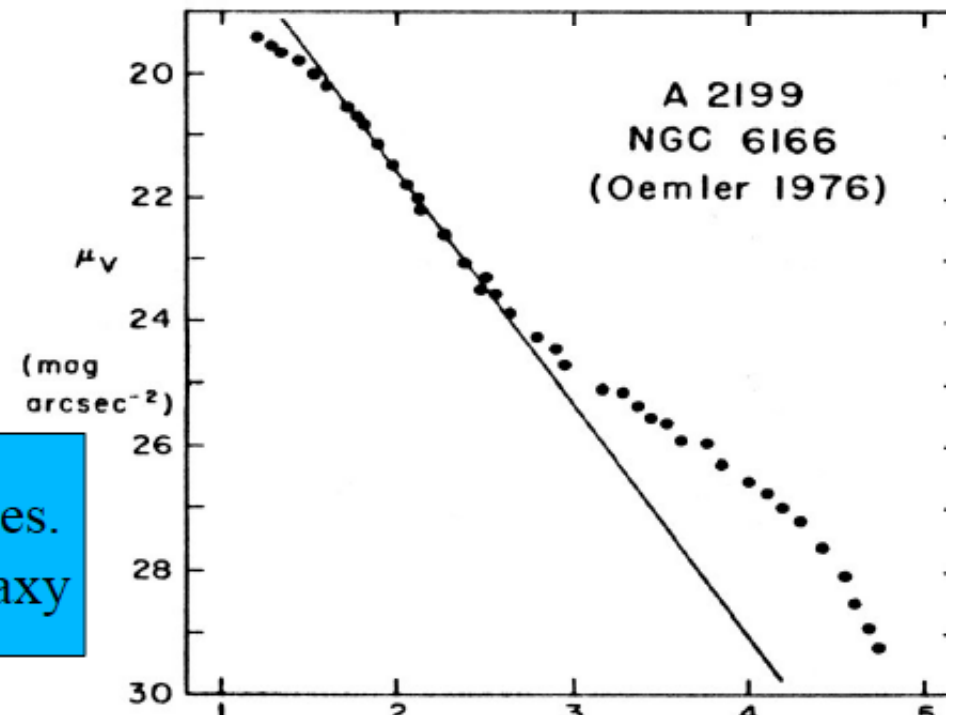
- They are usually found at the center of the galaxy distribution or in local density maxima (Beers & Tonry 1986).
- In some cases they have blue cores (McNamara & O'Connell 1989), or multiple nuclei (Morgan & Lesh 1965)
- 1/4 of the most luminous radio galaxies (WAT) are cD. The term was introduced in a study of optical counterparts of luminous radio galaxies (Matthews, Morgan, & Schmidt, 1964)
- Faint extended envelopes (Oemler 1974, Schombert 1988)

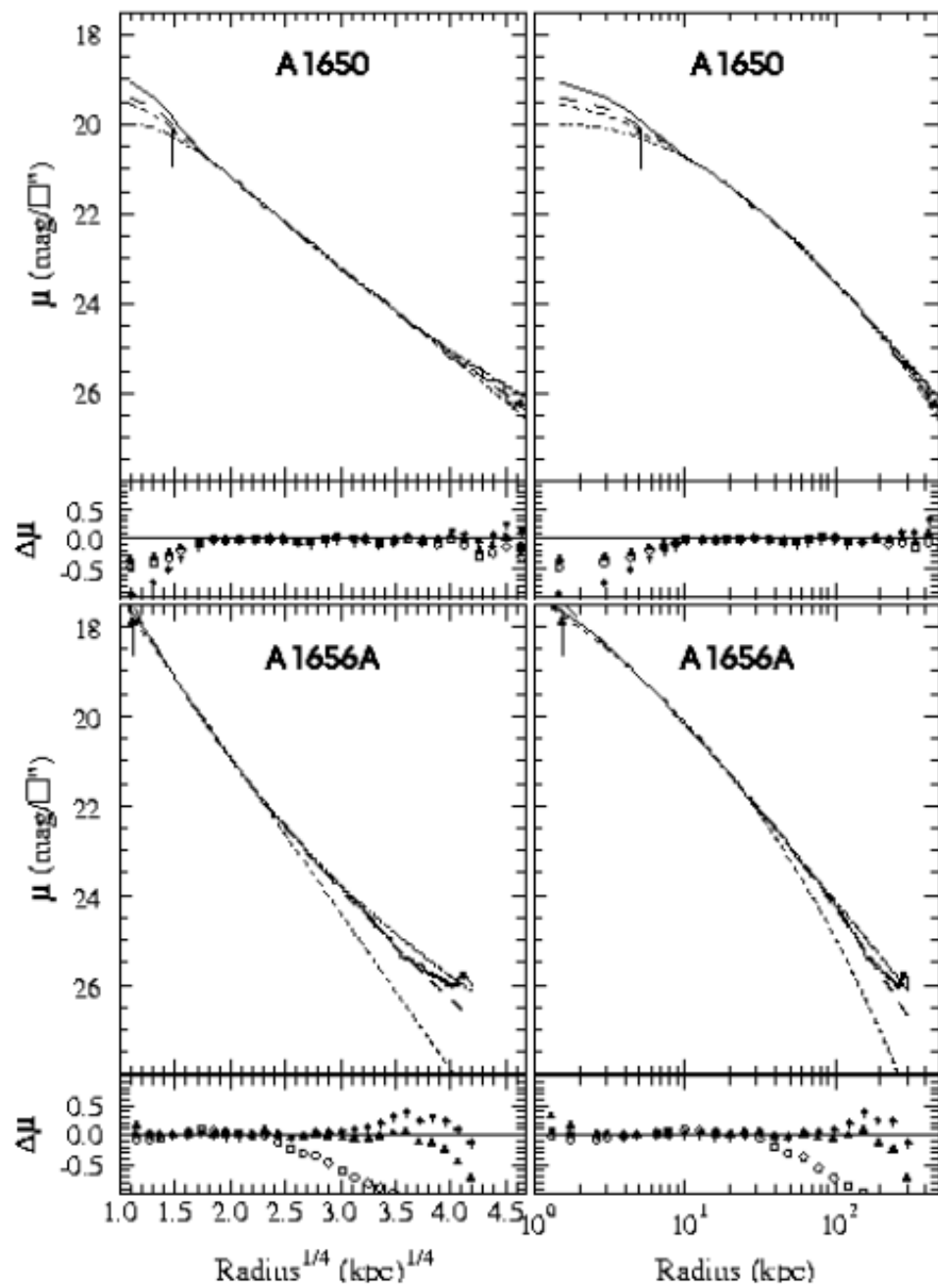
Commonly thought as...



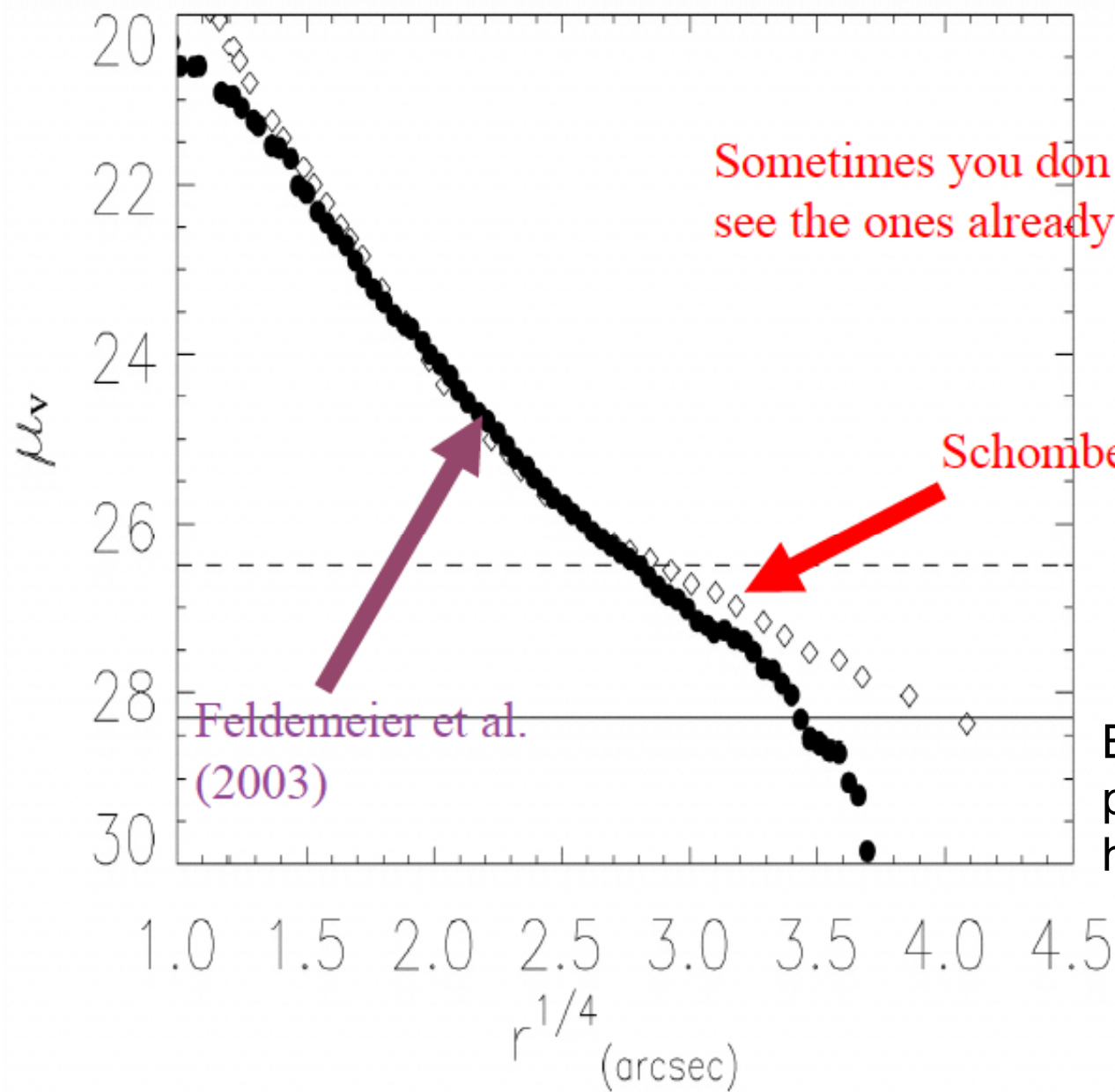
Multiple Nuclei: many of them are just projections

cD galaxies have faint-extended envelopes.
Schombert's (1988) definition of cD galaxy





Sometimes
you see the envelopes!!!

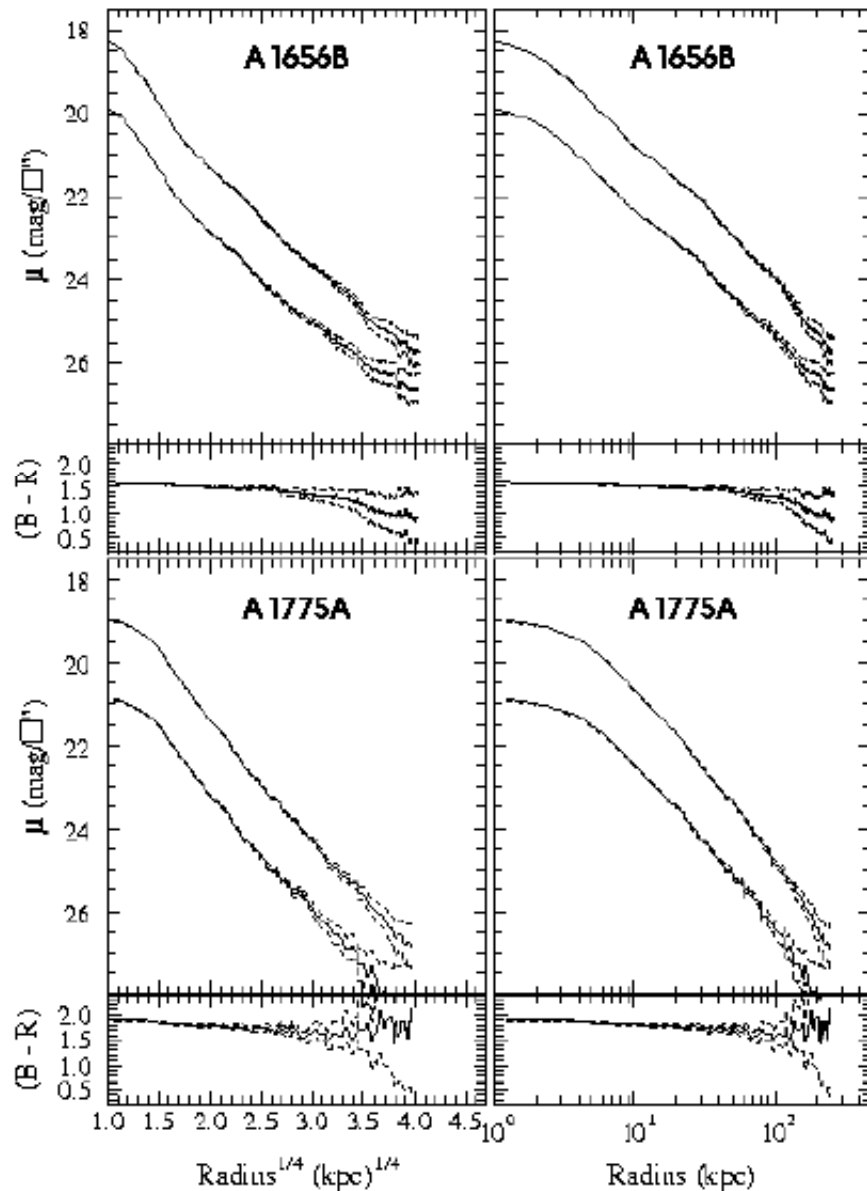


Sometimes you don't see the ones already discovered!!

Schombert (1988)

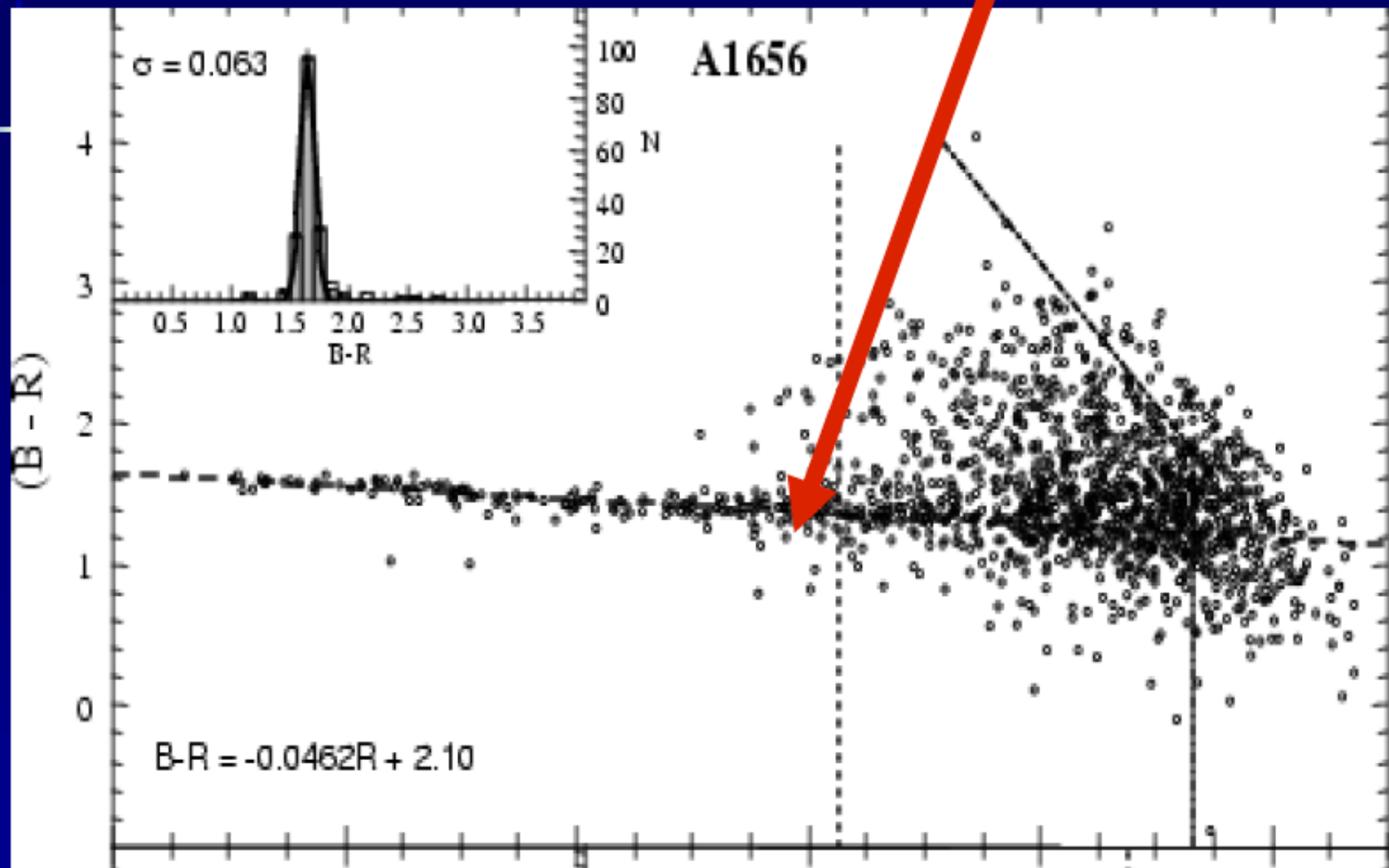
Feldemeier et al. (2003)

Because of these problems, the term BCG has been adopted.



cD galaxies have color gradients, the outskirts are bluer than the inner parts. Beware there are large error... B-R~1.3!!!

The colors of cD envelopes are similar to those of dwarf galaxies



$R=14$

$R=23$

How do cD galaxies form?

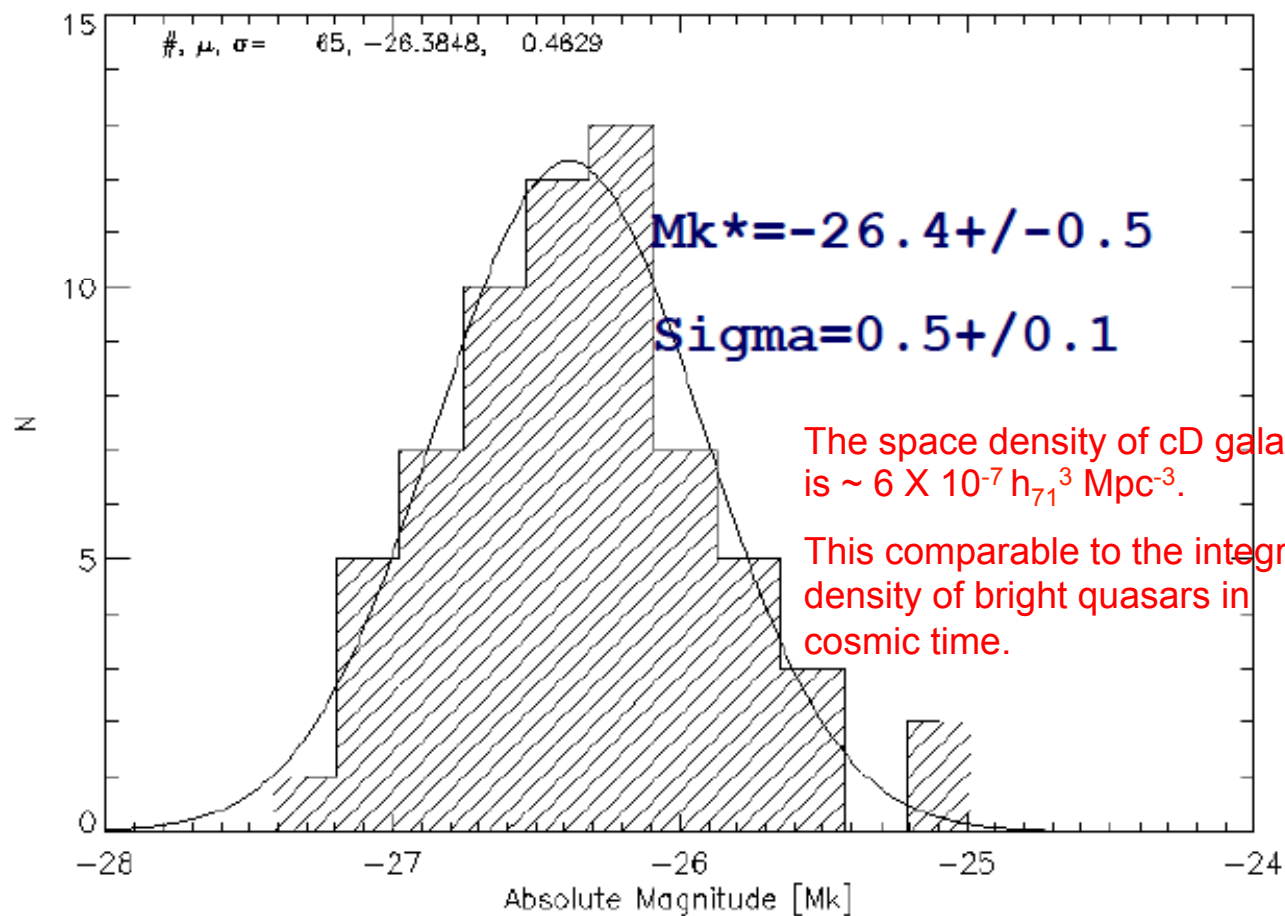
- Merger of giant galaxies. Not very efficient once the cluster have virialized.
- Dynamical effects: tidal truncation, tidal stripping, ♣ dynamical friction, etc.
- The fashionable harassment...
- Disruption of Galaxies. We see some evidence at low redshift.

♣ Tidal stripping (Aguilar & White 1985, 1986)

cD Galaxy Working Definition

- Take Abell statistical sample (**1682 clusters**). Richness 1 (50 galaxies, $m_3 - m_3 + 2$ within 1 Abell radius $3 h^{-1}_{50}$ Mpc) estimated redshift $0.02 < z < 0.2$; solid angle approx. 2/3 of the sky (14 438.2 sq. deg.).
- **cD are BGCs in BM I - BM I-II clusters**. There are **137 cD galaxies in Abell's sample**. cD galaxies overwhelm satellite galaxies. Classification from Leir & van den Bergh (1977). Leir's M.Sc. Thesis at U of T.

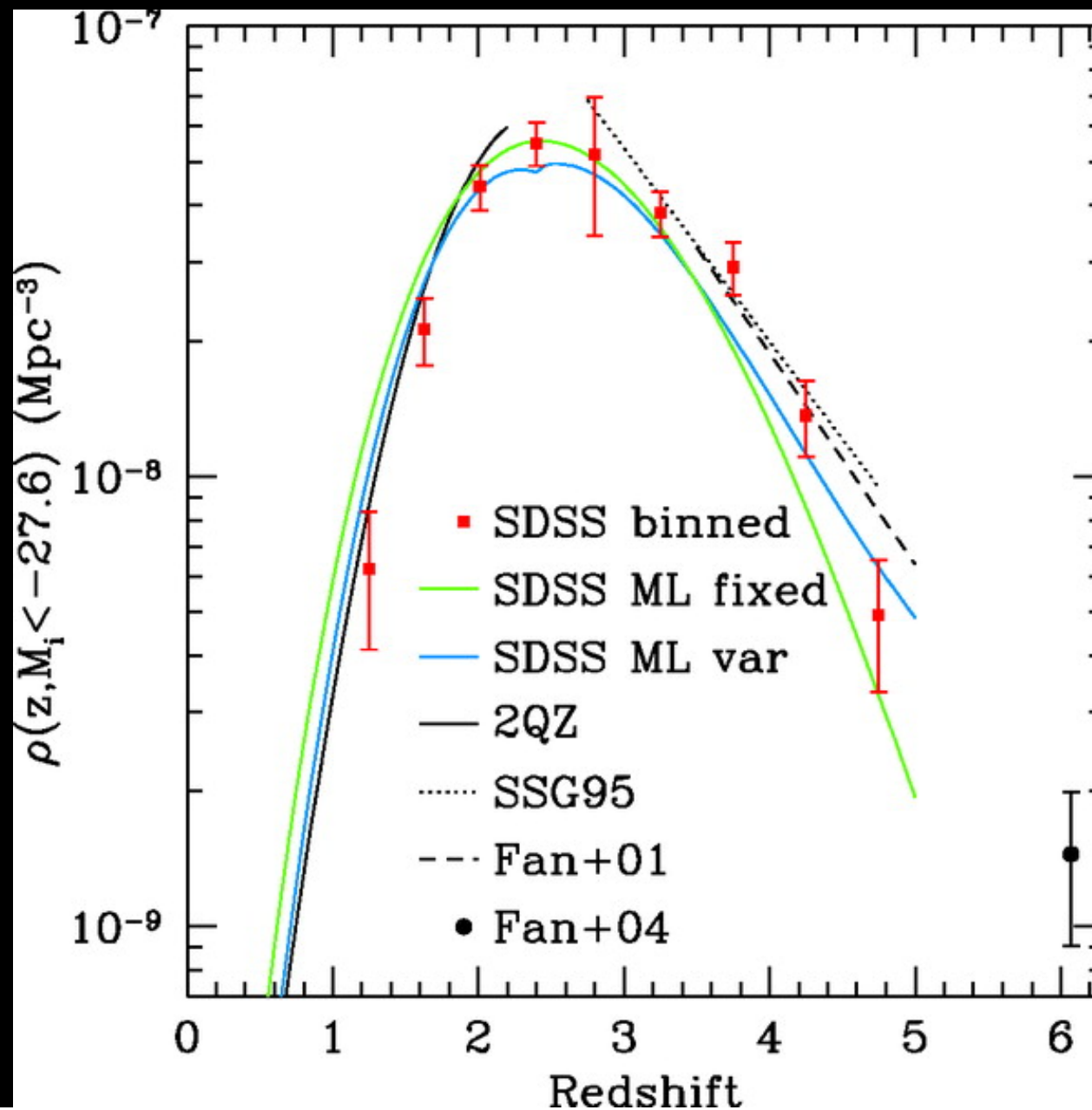
The LF of cD galaxies

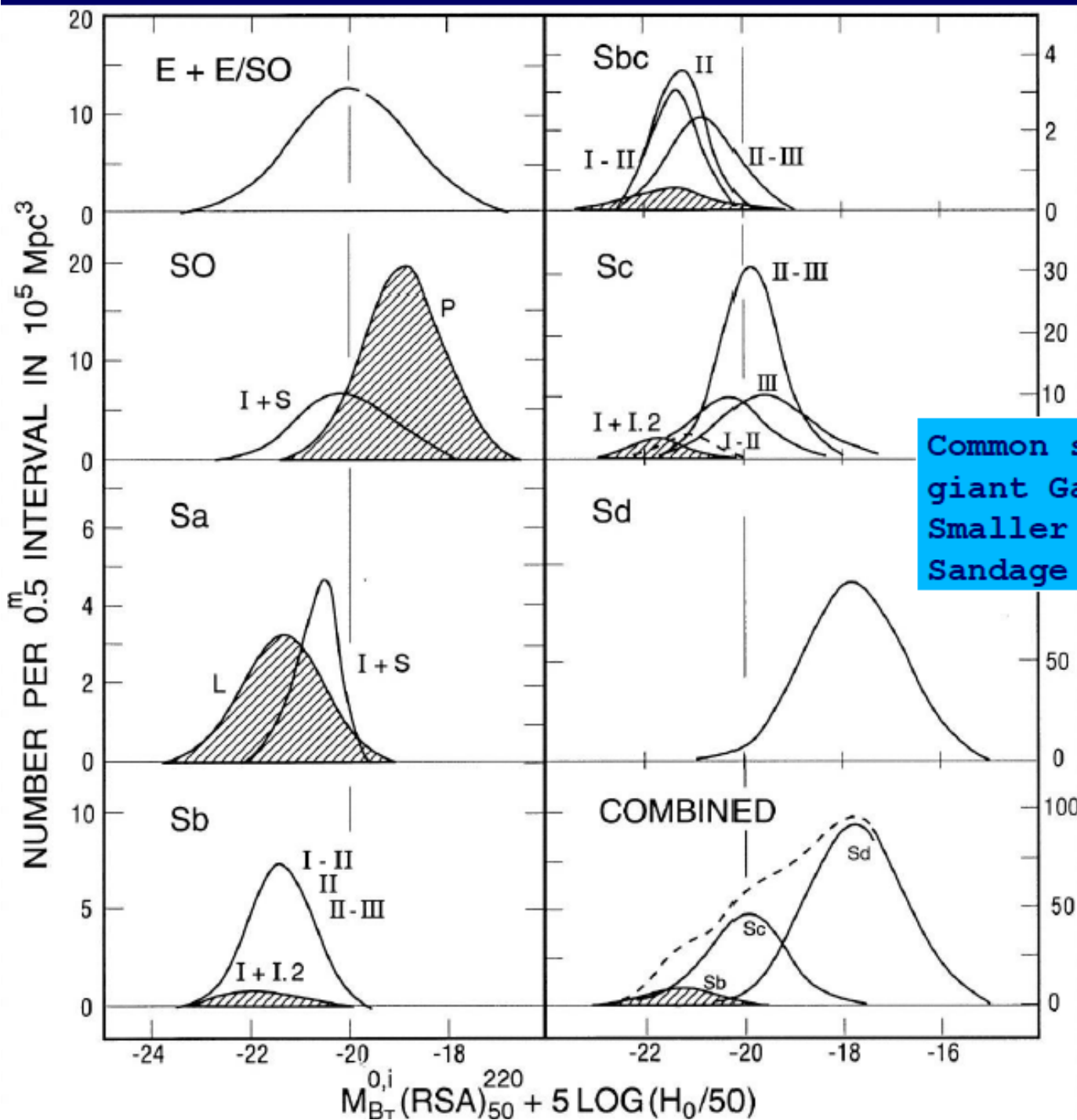


The space density of cD galaxies is $\sim 6 \times 10^{-7} h_{71}^3 \text{ Mpc}^{-3}$.

This comparable to the integrated density of bright quasars in cosmic time.


Bright QSO from SDDS DR3, Richards et al. 2006





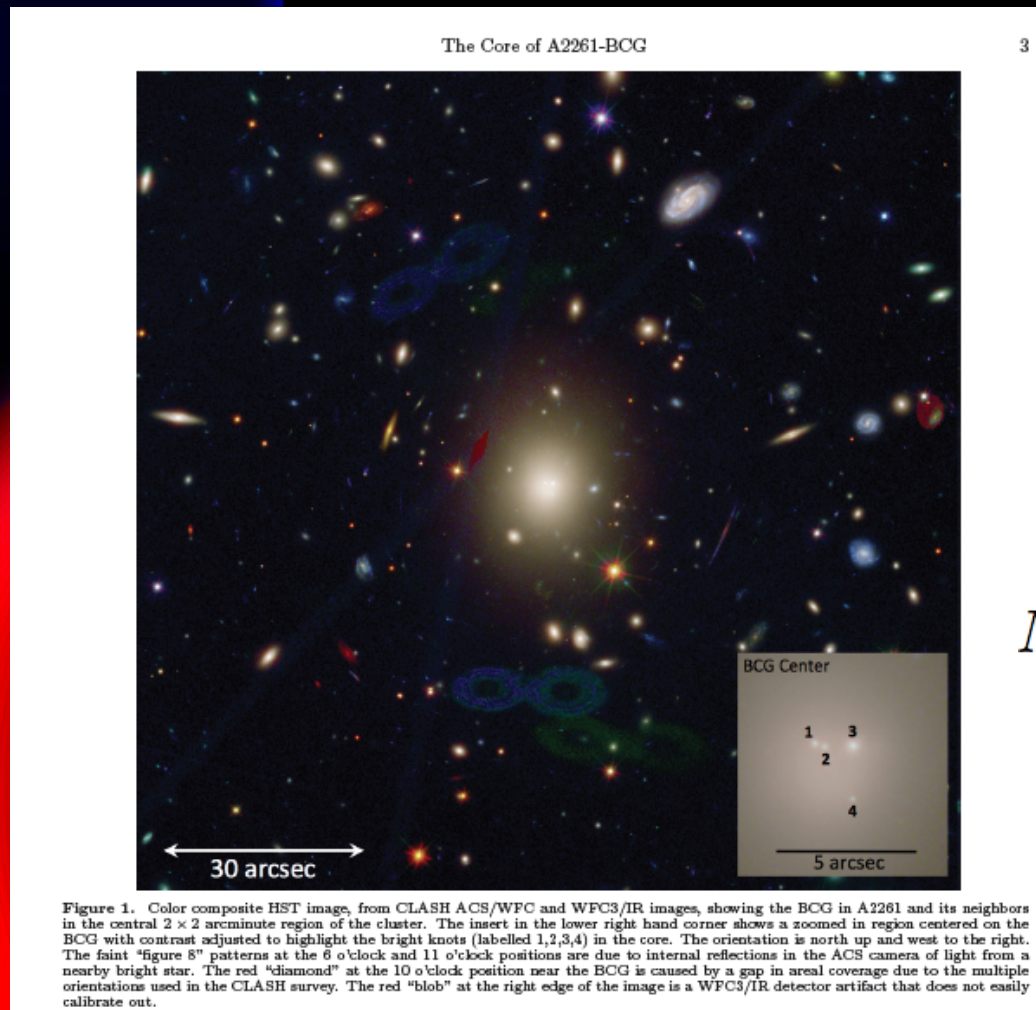
Connection

Common shape of F for
 giant Galaxies, but
 Smaller dispersion. (Sd, 0.8)
 Sandage 2005, de Lapparent 03



cD Galaxies & Other High Surface Brightness Galaxies follow Gaussian LF, but cDs have the narrowest distribution.

...And we bumped into Postman et al. 2012



A2261-BGC

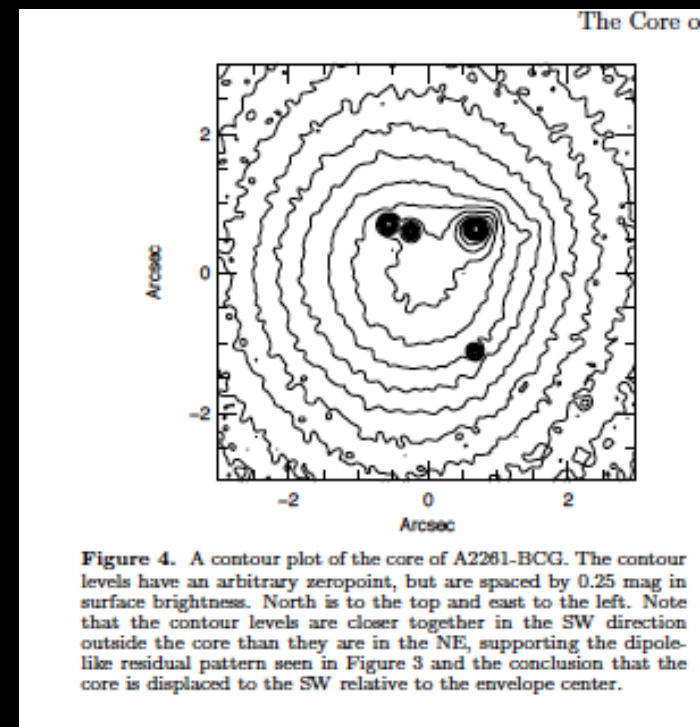
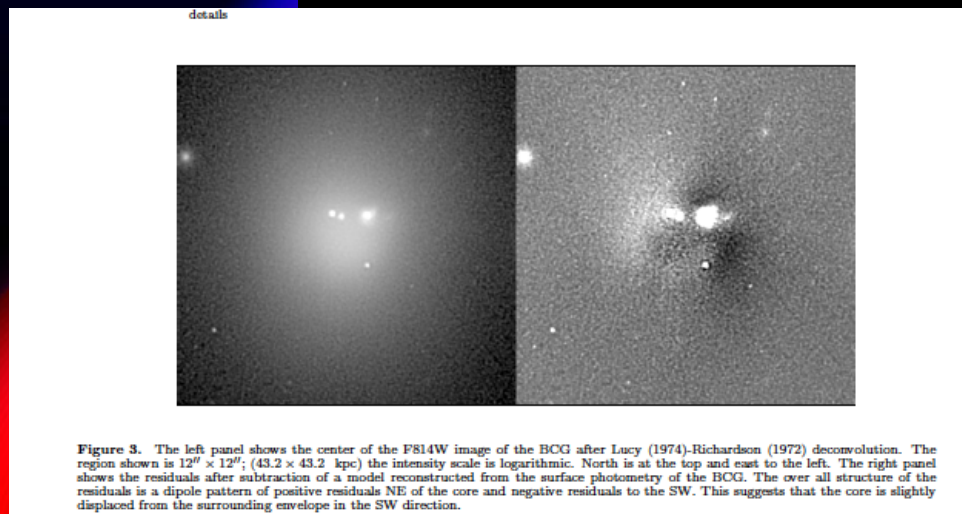
$z=0.2233$

$\alpha_{2000}=17:22:27.18$

$\delta_{2000}=+32:07:57.1$

$$M_{2500} = (2.9 \pm 0.5) \times 10^{14} M_{\odot}$$

Postman et al. 2012 were introducing the largest core yet detected in any galaxy...



The Nuker Law (Lauer et al. 1995)

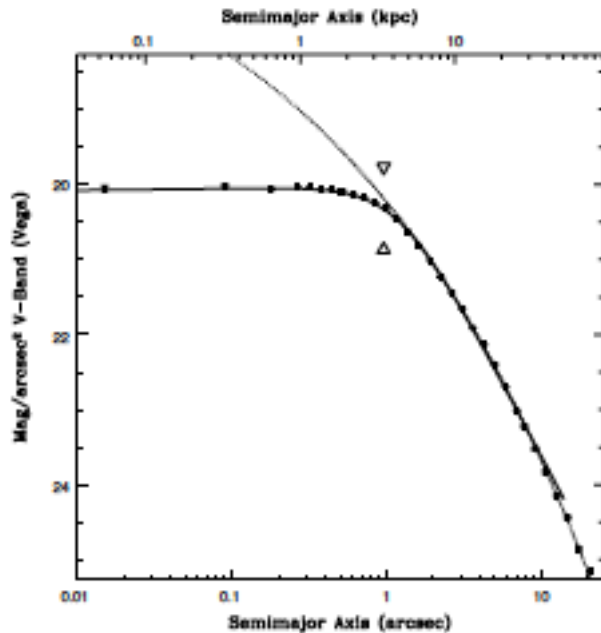


Figure 5. The central surface brightness profile of A2261-BCG as measured (solid points) is shown with two "Nuker-law" profile fits (Lauer et al. 1995). The error bars are smaller than the points, but for the central few measurements. For comparison to previous studies the profile is normalized to $z = 0$ V-band (Vega). The solid line is the best-fitting Nuker-law and features a slightly depressed ($\gamma = -0.01$) cusp as $r \rightarrow 0$. The dotted line is an $r^{1/4}$ -law (an $n = 4$ Sérsic-law) fitted to the envelope. The triangles indicate the cusp-radius.

$$I(r) = 2^{(\beta-\gamma)/\alpha} I_b \left(\frac{r_b}{r}\right)^\gamma \left[1 + \frac{r}{r_b}\right]^{(\gamma-\beta)/\alpha}$$

$$r_\gamma \equiv r_b \left(\frac{1/2 - \gamma}{\beta - 1/2}\right)^{1/\alpha}$$

$$\gamma = -0.01, \beta = 1.56, \alpha = 2.41 \pm 0.18, r_b = 1.2''$$

$$r_\gamma = 0''.89 \pm 0''.02; 3.2 \pm 0.1 \text{ kpc}$$

Luminosity vs. Cusp Radius

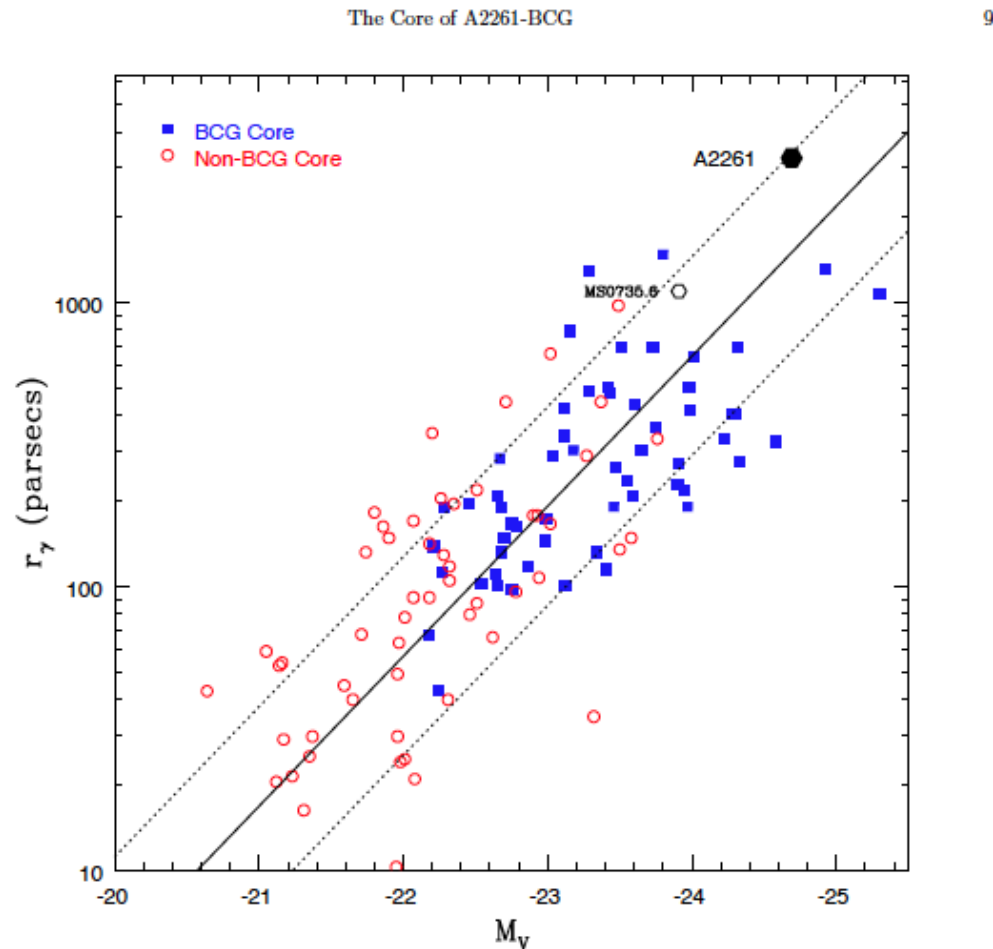
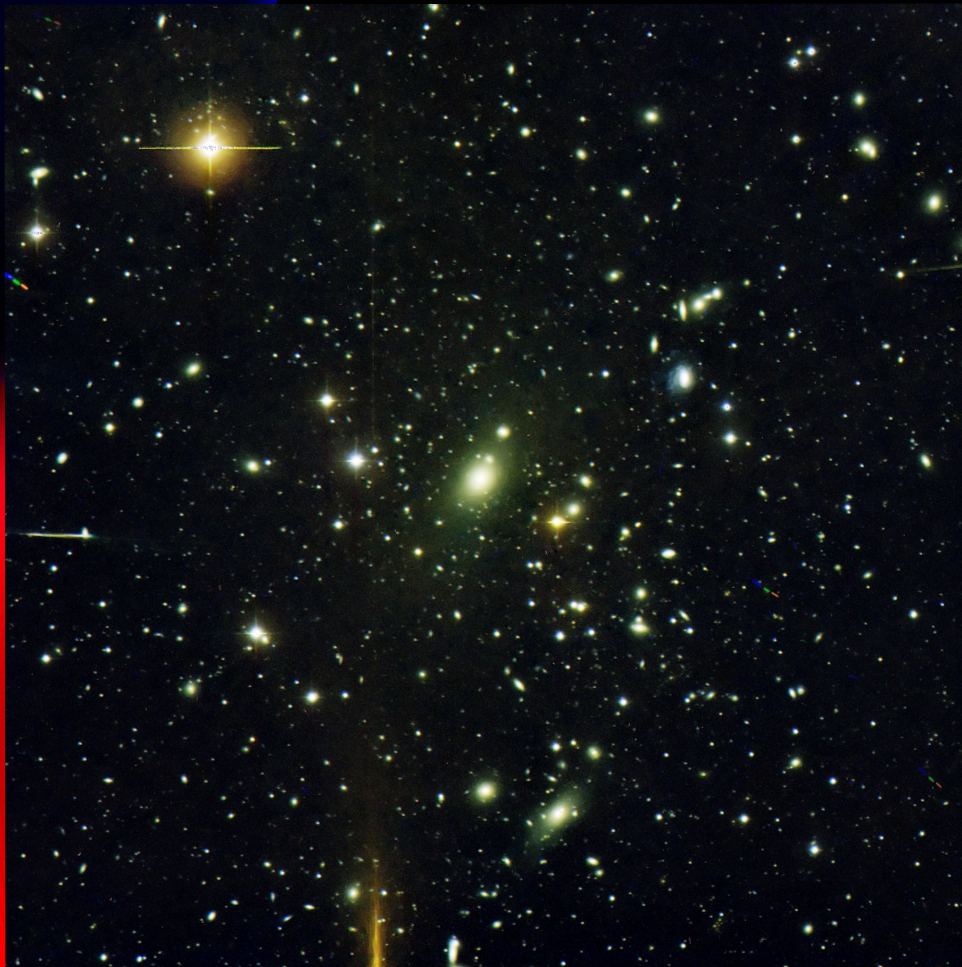


Figure 7. The relationship between cusp-radius and total galaxy luminosity (V-band Vega-based). The galaxy sample plotted was assembled in Lauer et al. (2007a) from a variety of sources (the figure is adopted from Figure 5 in that paper). The BCGs in particular come from the Laine et al. (2002) sample. The Lauer et al. (2007a) $r_\gamma - L$ relationship (also given in equation 4) is plotted; the dotted lines indicate $\pm 1\sigma$ scatter about the mean relationship. A2261-BCG is plotted at the top, clearly has a cusp-radius larger than all other galaxies in the sample. The large core in the MS0735.6+7421 BCG discovered by McNamara et al. (2009) is also plotted for comparison.

But, I recalled Abell 85. James P. Brown and I have worked on it back in 1995 at DA&A, UofT



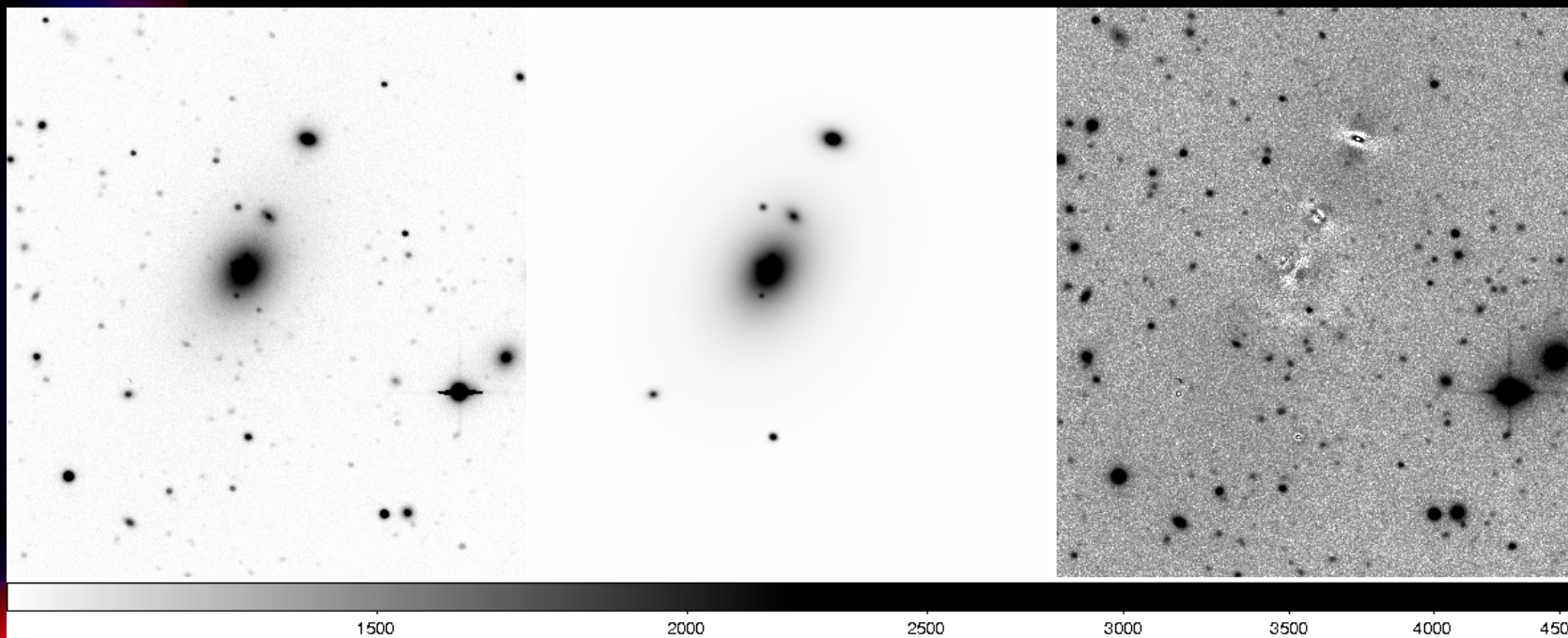
LOCOS, *BR* images,

Taken with the KPNO 0.9m
Telescope using the T2KA CCD.
Seeing ~ 1.6 arcsec

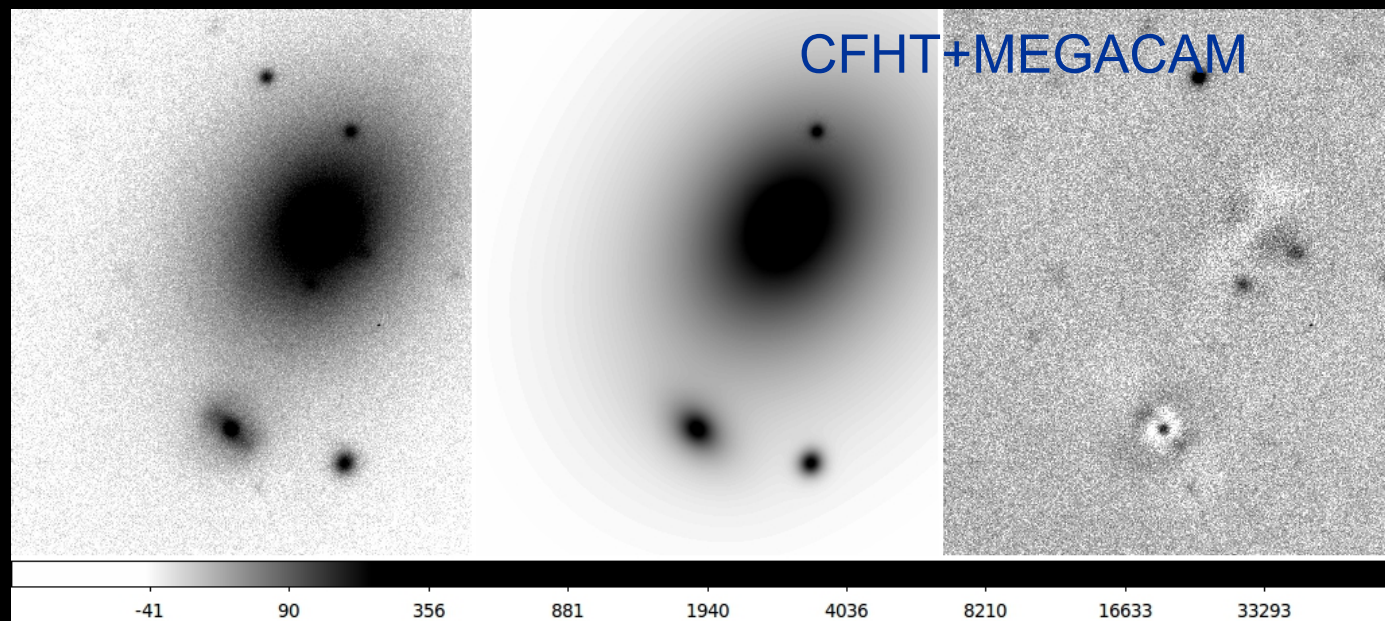
Scale 0.68 arcsec/pixel
 $z=0.05529 \pm 2.4 \times 10^{-4}$

Hoessel and collaborators during
the early 80's have already singled
out that that HOLM 015A had a
very large core (modified Hubble
Profile)

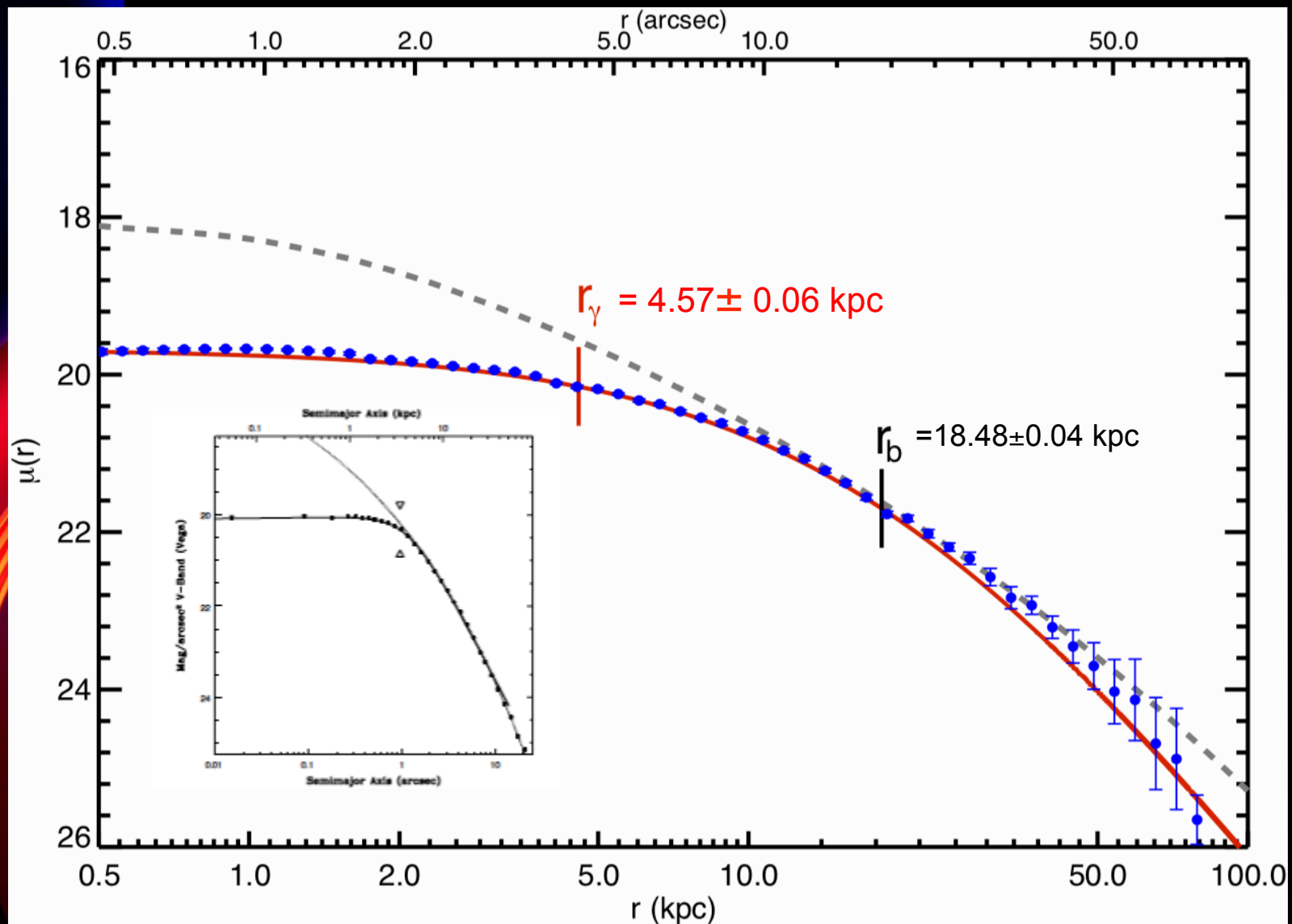
KPNO 0.9m + T2KA



We
tried
DGCG
on A85
using
the
Nuker
Law.



The surface brightness profile of Holm 15A



A2266-BGC

$$r_\gamma = 0''.89 \pm 0''.02; 3.2 \pm 0.1 \text{ kpc}$$

Table 1. Holm 15A: Nuker law fits

$\mu_b \left[\frac{\text{mag.}}{\text{arcsec}^2} \right]$	$r_b \text{ [arcsec]}$	$r_b \text{ [kpc]}$	α	β	γ	$r_\gamma \text{ [kpc]}$	e	P. A.	Data
21.78 ± 0.01	17.21 ± 0.04	18.48 ± 0.04	1.24 ± 0.01	3.33 ± 0.02	0.0 ± 0.0	4.57 ± 0.06	0.26	-33.33	★
22.32	19.09	20.50	1.22	3.62	0.0	4.57	0.24	-34.07	•

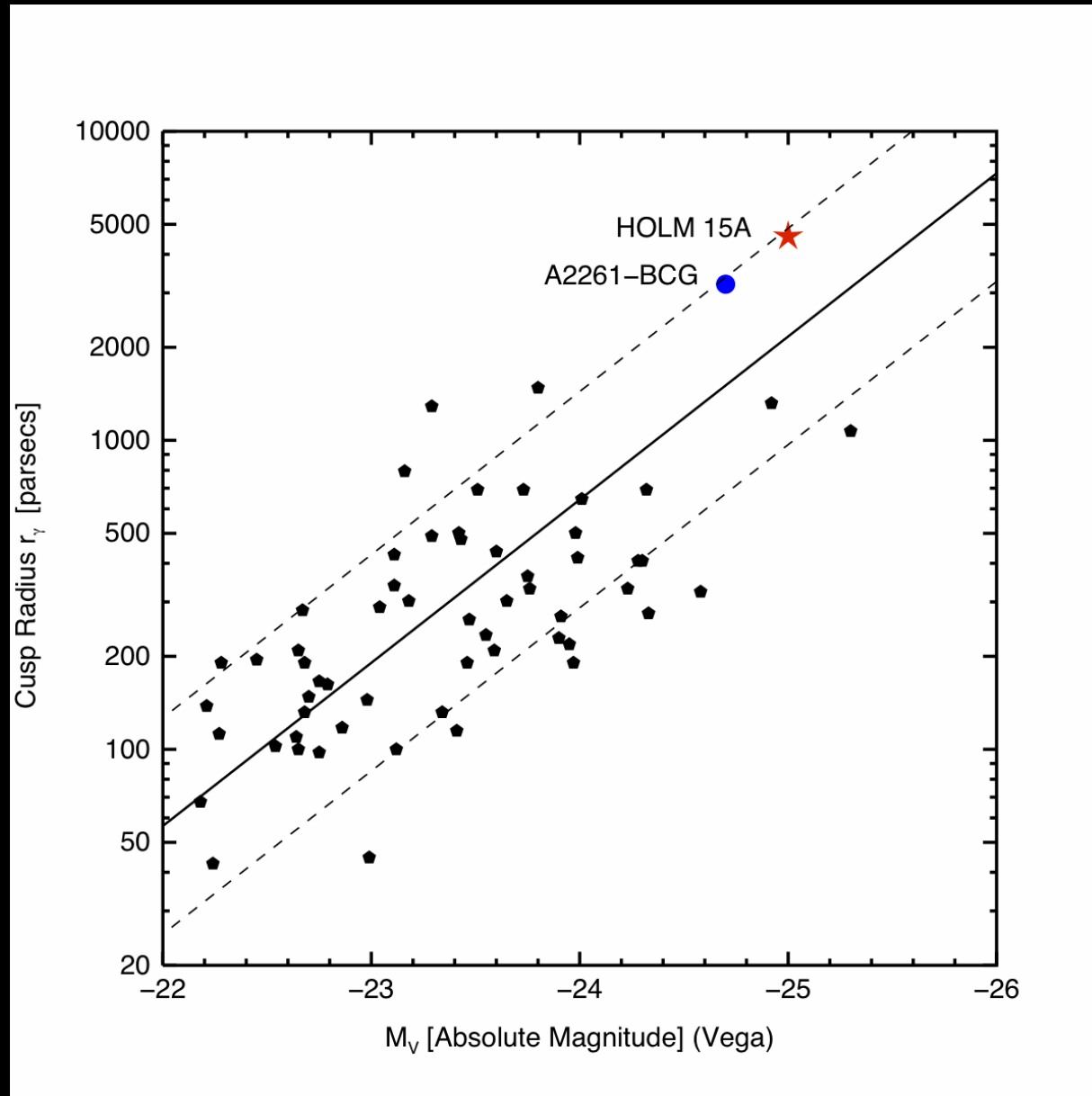
COLUMNS– 1: surface brightness, 2: break radius in arcsec, 3: break radius in kpc, 4: outer power index, 5: power index at r_b , 6: outer power index, 7: cusp radius in kpc, 8: ellipticity, 9: position angle in degrees, 10: Data Source.

★ LOCOS (López-Cruz 1997), Telescope: KPNO 0.9m, CCD: T2KA; pixel scale: $0''.68/\text{pixel}$. filter: R (Kron-Cousins), exposure time: 900 s, seeing: $1''.6$ FWHM, FOV: $23'.2 \times 23'.2$.

• MENEaCS (Sand et al. 2011), Telescope: CFHT 3.5m, CCD: MegaCam; pixel scale: $0''.187/\text{pixel}$, filter: r' (SDSS), exposure time: 120 s, seeing: $0''.74$ FWHM, FOV: $0'.96 \times 0'.94$.

Note. — centroid: $\alpha_{2000} = 00^{\text{h}}41^{\text{m}}50^{\text{s}}.467$, $\delta_{2000} = -09^{\circ}18'11''.57$

The Largest Core known so far!!!



Bonfini et al. (2015) has challenged this result.

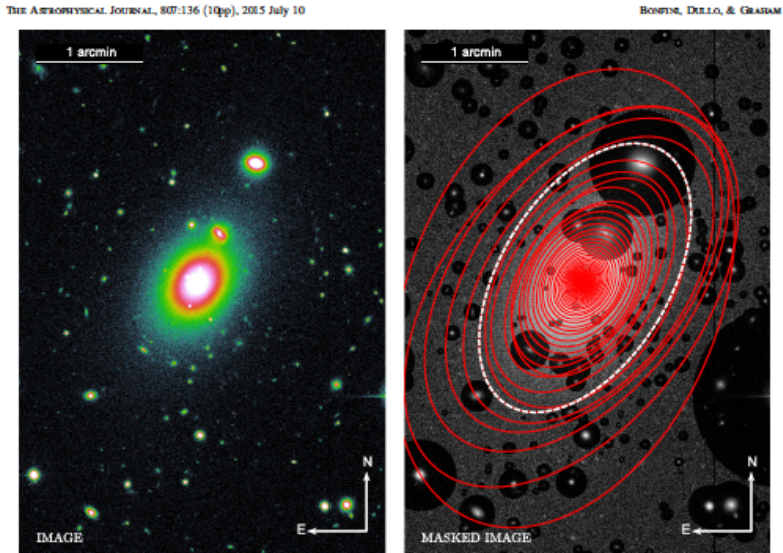


Figure 1. CFHT-MegaPrime r -band image (left) and relevant mask (right) for Holm 15A. The masked areas (see Section 2) have been arbitrarily decreased in intensity so as to still show the contaminating objects. We overlap the elliptical isophotes from IRAF *allpix* (the thickness of the ellipses is not considered in this representation). The dashed ellipse corresponds to the limiting surface brightness at which we truncated our 1D analysis ($\mu \sim 25.5 \text{ mag arcsec}^{-2}$; see Section 2). This ellipse also corresponds to the physical extent of the 2D fit (everything outside the dashed curve was masked for the 2D fit; see Section 2).

Table 1
CFHT-MegaPrime Image Characteristics

Target (1)	R.A. (J2000) (hh:mm:ss) (2)	Decl. (J2000) (dd:mm:ss) (3)	D (Mpc) (4)	$m - M$ (mag) (5)	Camera/Filter (6)	Exposure (s) (7)	Scale ($''$ /pixel) (8)
Holm 15A	00 ^h 41 ^m 30 ^s .5	-09 ^o 18'11"	253	37.02	MegaPrime/ r	120	0.386

Notes. Details of the CFHT-MegaPrime image used for the current work. (1) Target name. (2, 3) Target coordinates from NED. (4) Luminosity distance from NED, corresponding to a redshift $z \sim 0.057$ (see footnote 2). (5) Distance modulus. (6) CFHT camera and filter. (7) Total exposure time. (8) Image pixel scale.

Core-Sérsic (Trujillo et al. (2004))

$$I(R) = I' \left[1 + \left(\frac{R_b}{R} \right)^\alpha \right]^{-\frac{\gamma}{\alpha}} \exp \left(-b \left[\frac{R^\alpha + R_b^\alpha}{R_e^\alpha} \right] \right)^{\frac{1}{(\alpha n)}}$$

$$I' = I_b 2^{-\frac{\gamma}{\alpha}} \exp b \left(\frac{2^{-\frac{\gamma}{\alpha}} R_b}{R_e} \right)^{\frac{1}{n}}$$

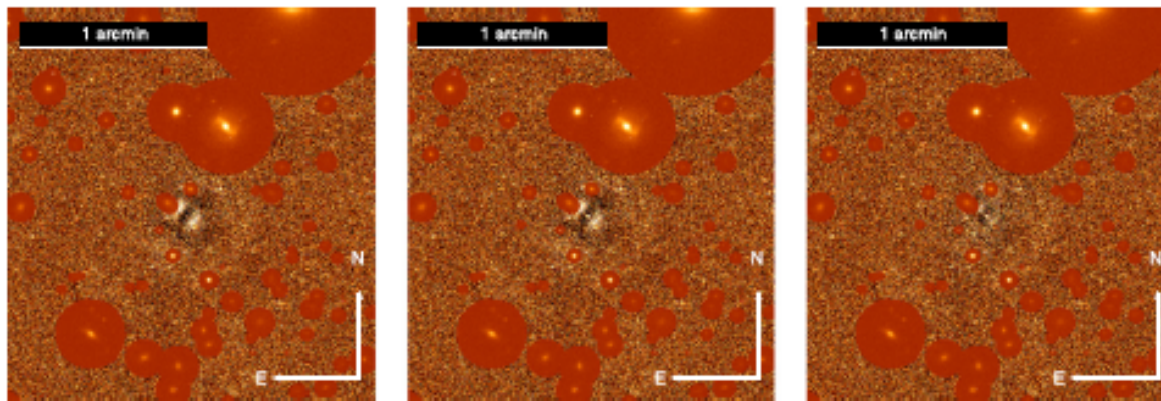
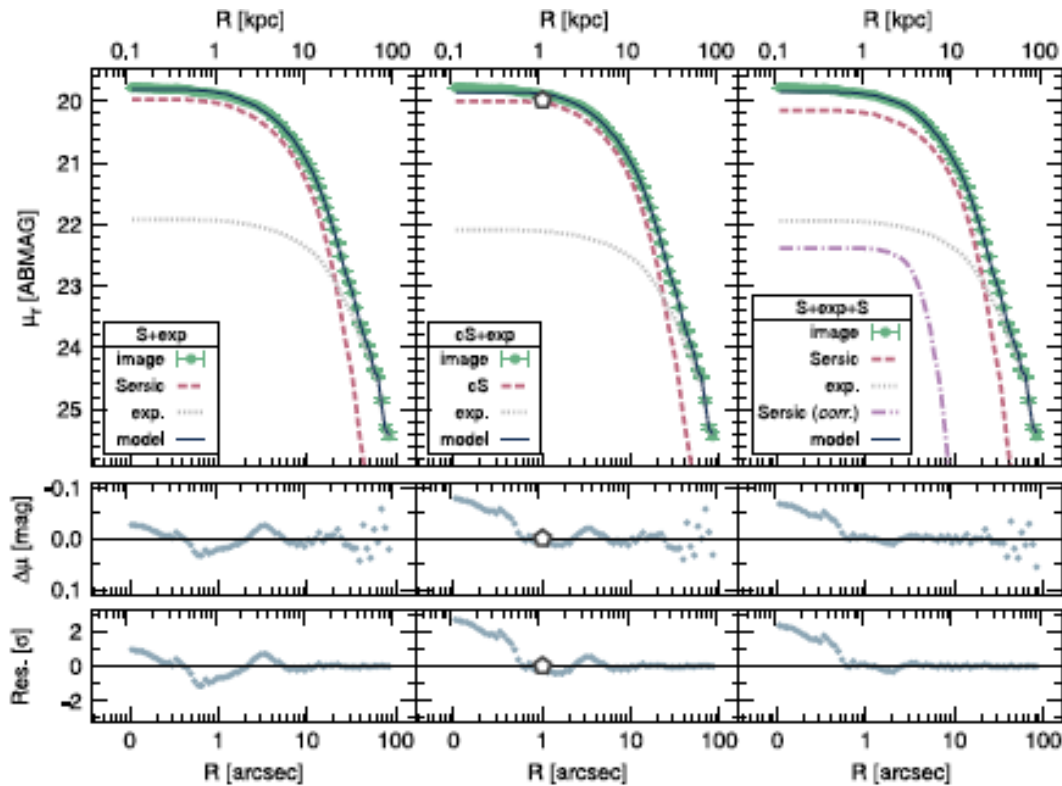


Figure 4. 2D image analysis: results of the fits to the 2D image of Holm 15A performed assuming: a Sérsic+exponential model (left); a core-Sérsic+exponential model (middle); and a Sérsic+exponential model plus an additional inner Sérsic component intended to compensate for the single ellipticity of the 2D model components (right). Top. The green data points represent the major-axis surface brightness profile measured over the isophotes defined using IRAF *ellipse* (i.e., the same measurement presented in Figure 2). The curves represent the surface brightness profile of the model images measured over exactly the same isophotes. The continuous curves show the global models, while the dashed curves represent their sub-components. We stress that these are *not* fits to the 1D profile, but rather surface brightness measurements (projections) of the 2D models. The pentagon indicates the location of the core-Sérsic model's break radius. The panels underneath the profiles represent the data residuals about the fitted models, first expressed in terms of the difference in surface brightness, and then in terms of residuals (in units of counts) divided by the standard deviation as measured on the "sigma" image. Bottom. The actual residual images that were minimized by GALFIT-CORSAIR. Masked objects are highlighted as in Figure 1.

Bonfini et al. (2015) used GALFIT-CORSAIR and showed that Core-Sérsic profile doesn't work for Holm 15A, but recovered López-Cruz et al. (2014) Nuker fit.

The Mexican Contribution to the field (ca. 1978)!!

THE TEMPERATURE AND DYNAMICS OF THE IONIZED GAS IN THE NUCLEUS OF OUR GALAXY

LUIS F. RODRIGUEZ* AND ERIC J. CHAISSON†
Harvard-Smithsonian Center for Astrophysics
Received 1978 August 18; accepted 1978 September 26

ABSTRACT

Observations of the H65 α (23.4 GHz), H84 α (10.9 GHz), and H94 α (7.8 GHz) radio recombination lines from Sgr A West are presented. We suggest that a core-halo model can satisfactorily account for the reported radio and infrared observations of this source. Due to instrumental limitations, the observed infrared lines are dominated by the core, while the observed radio radiation arises mostly in the halo. Although more than a factor of 10 brighter than the halo, the core is responsible for only about one-fourth of the integrated thermal continuum from Sgr A West. Our model implies that the neon abundance determination from infrared observations, previously considered consistent with the solar value, should be revised upward by a factor of 4. This suggested enrichment of neon relates strongly to our derivation of an unusually low electron temperature, $T_e = 5000 \pm 1000$ K, since nebular cooling is expected to be enhanced by an overabundant heavy-element content. The dynamical structure of Sgr A West can be explained in terms of Keplerian rotation due to the gravitational field of the normal stellar population plus a central mass point of $5 \times 10^6 M_\odot$.

... a central mass point of $5 \times 10^6 M_\odot$

No. 3, 1979

IONIZED GAS IN GALAXY

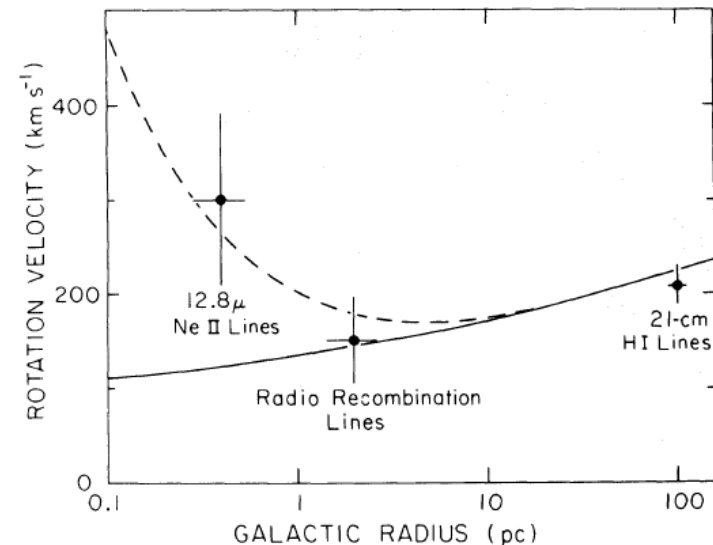
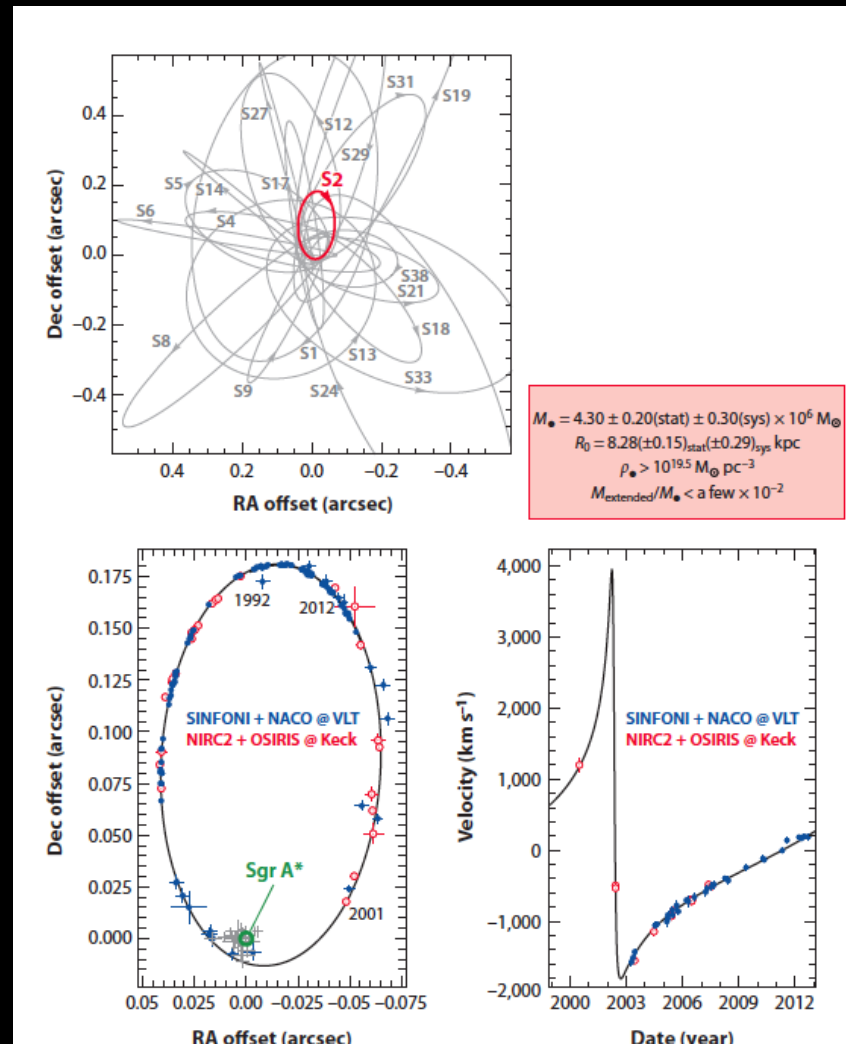


FIG. 4.—A tentative rotation curve for the galactic nucleus. *Solid curve*, model having normal stellar population; *dashed curve*, model having the same normal stellar population plus a mass point of $5 \times 10^6 M_\odot$. Error bars are estimated to be twice the standard deviation.

SMBH Masses: Directly



Measuring the mass of the SMBH in our Galaxy through stellar motions.

SMBH Masses: Directly

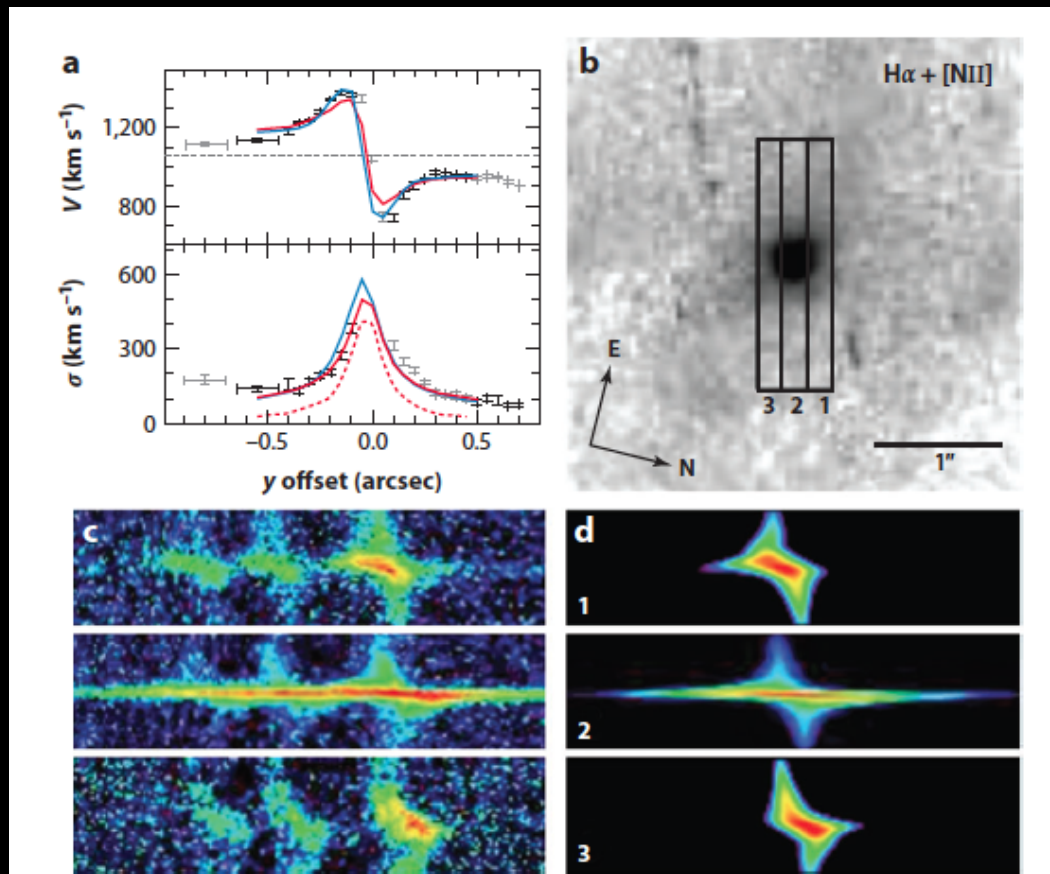


Figure 6

Space Telescope Imaging Spectrograph (STIS) observations of NGC 4374 (Bower et al. 1998) as analyzed by Walsh, Barth & Sarzi (2010). (b) WFPC2 continuum-subtracted H α + [NII] image showing the nuclear disk of ionized gas. The footprints of the three slits are overlaid. (a) The radial profiles of [NII] λ 6583 mean velocity and velocity dispersion (in kilometers per second) along the central slit position. Superposed are predictions of the best black hole model with (blue curve) and without (red solid curve) correction for asymmetric drift. The red dotted curve shows the contribution from rotational line broadening; because these velocity dispersions are smaller than the ones observed, the intrinsic velocity dispersion must be significant. The bottom panels show (c) the continuum-subtracted, two-dimensional STIS spectra of the H α + [NII] region and (d) the synthetic spectra of [NII] λ 6583. The vertical axis is the spatial direction, and the horizontal axis shows wavelength increasing toward the right. The three slit positions correspond to the locations labeled in panel b.

SMBH Masses: Scaling Laws

$$\frac{M_{\bullet}}{10^9 M_{\odot}} = (0.544^{+0.067}_{-0.059}) \left(\frac{L_{Ks, \text{bulge}}}{10^{11} L_{Ks\odot}} \right)$$

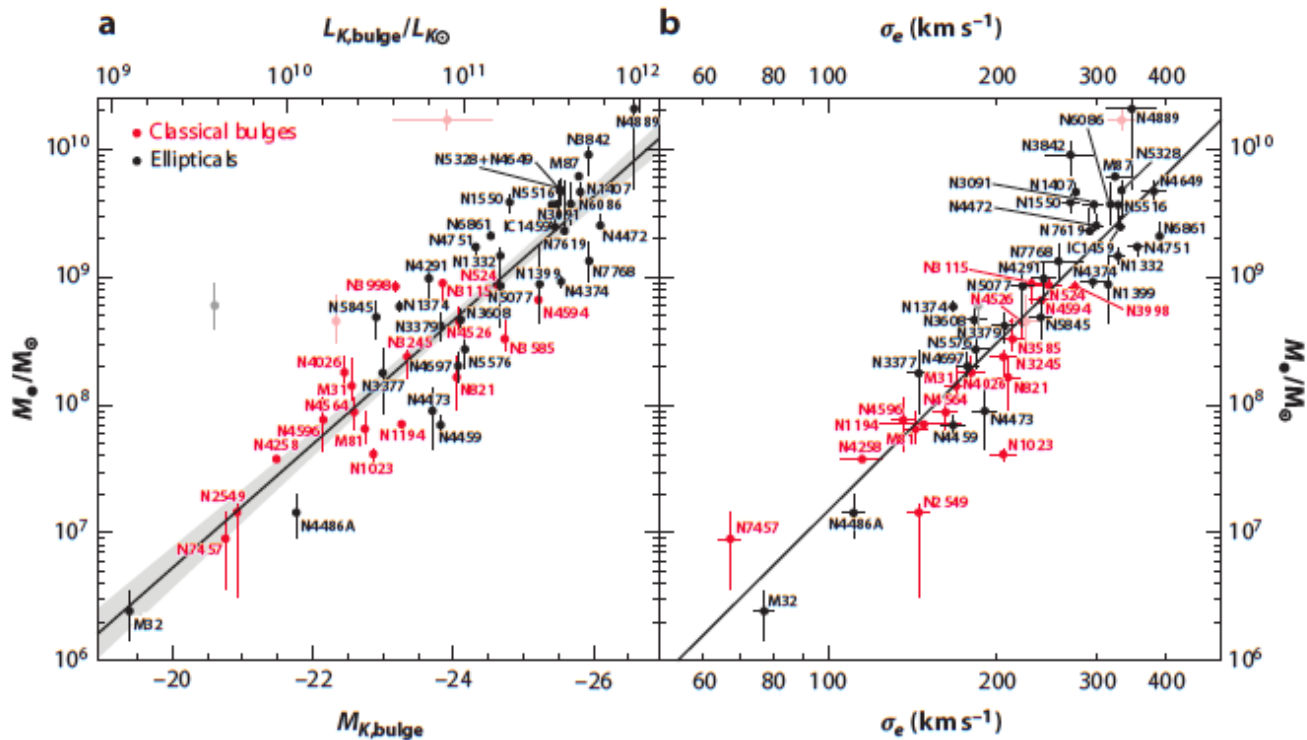
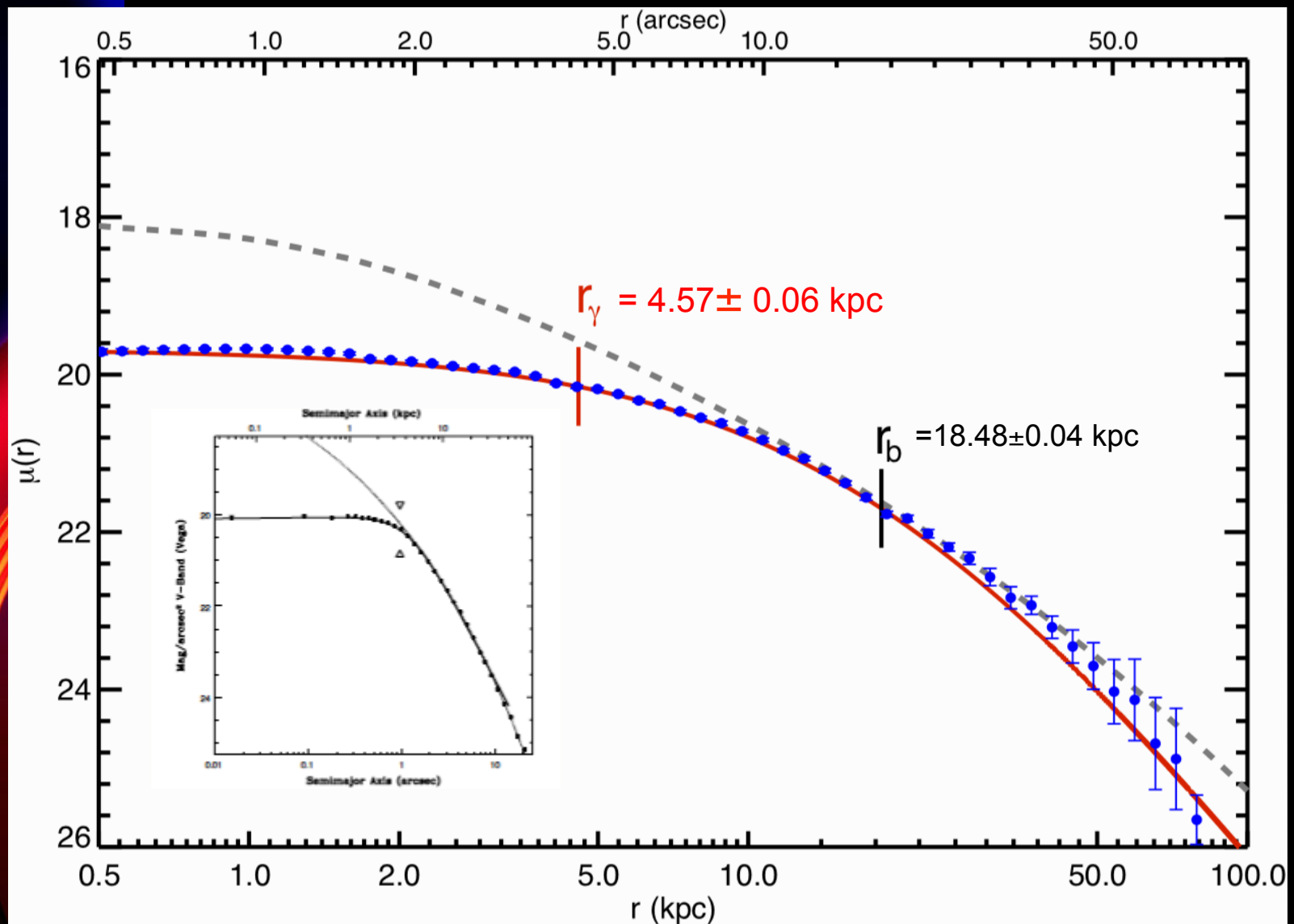


Figure 16

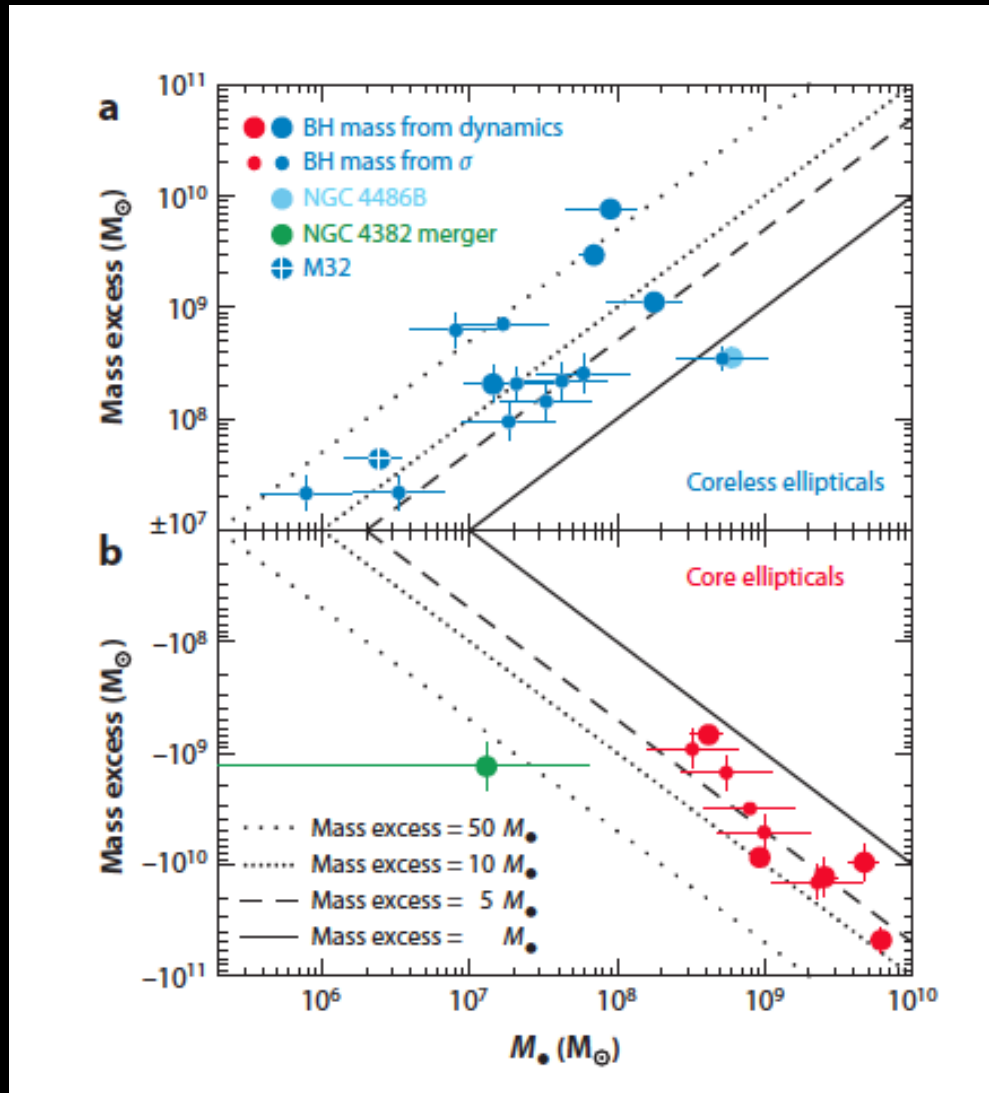
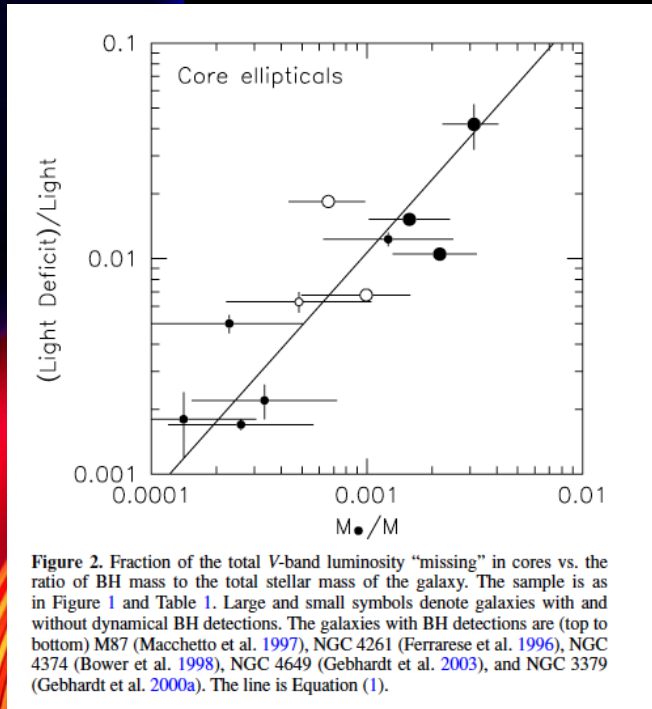
Correlation of dynamically measured black hole mass M_{\bullet} with (a) K -band absolute magnitude $M_{K, \text{bulge}}$ and luminosity $L_{K, \text{bulge}}$ and (b) velocity dispersion σ_e for (red) classical bulges and (black) elliptical galaxies. The lines are symmetric least-squares fits to all the points except the monsters (points in light colors), NGC 3842, and NGC 4889. Figure 17 shows this fit with $1\text{-}\sigma$ error bars.

$$\frac{M_{\bullet}}{10^9 M_{\odot}} = (0.310^{+0.037}_{-0.33}) \left(\frac{\sigma}{200 \text{ km s}^{-1}} \right)^{4.38 \pm 0.29}$$

The surface brightness profile of Holm 15A



Missing light correlates with BH mass (Kormendy & Bender 2009; Kormendy & Ho 2013)



$$\frac{M_{BH}}{10^8 M_{\odot}} = (0.81^{+0.18}_{-0.15}) \left(\frac{L_{V, def}}{L_{V\odot}} \right)^{0.95^{+0.18}_{-0.13}}$$

extra" in coreless galaxies or (b) "missing" in cores versus black hole (BH) mass. Large galaxies with dynamical M_{BH} measurements; small ones use M_{BH} given by the $M_{BH}-\sigma$ relation. Mass excesses tend to be larger than mass deficits and scatter with M_{BH} . It is important to note that the merger in progress NGC 4382 has a high luminosity but deviates to small M_{BH} . If cores are excavated by BH binaries, this suggests they do not lack a big BH (or BH binary) but rather that this BH (or BH binary) is not resident at the center. Updated from Kormendy & Bender (2009).

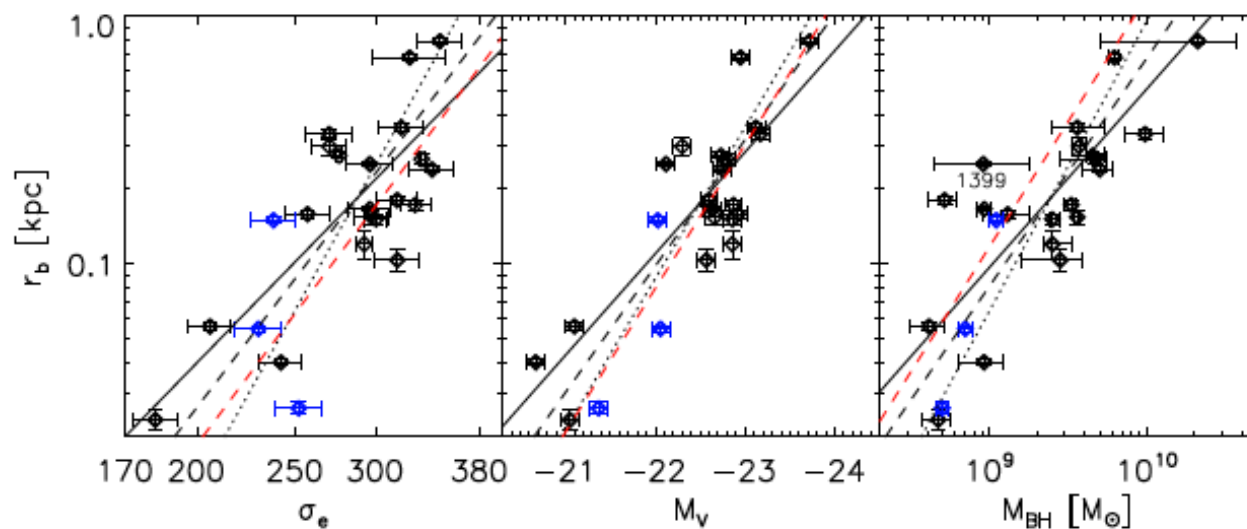


Fig. 7.— From left to right: the break radius is plotted as a function of velocity dispersion within the effective radius σ_e , luminosity of the host galaxy and black hole mass. The 20 galaxies with reliable M_{BH} are represented by the black diamonds. The additional three galaxies without reliable M_{BH} (NGC 4552, NGC 5813, NGC 5846) are shown in blue. The black lines show our fits to all black datapoints. The solid lines show the fits when x -axis variables are treated as independent variables/predictors, the dotted lines are fits when these parameters are the response, and the dashed lines represent the symmetrical bisector regression; see Table 5. The red dashed lines are the bisector fits from Dullo & Graham (2012), i.e. their equations 5, 6, and 12, respectively.

$$\log \left(\frac{M_{\bullet}}{3 \times 10^9 M_{\odot}} \right) = (0.59 \pm 16) + (0.92 \pm 0.20) \log \left(\frac{r_b}{\text{kpc}} \right)$$

Lauer et al. 2007

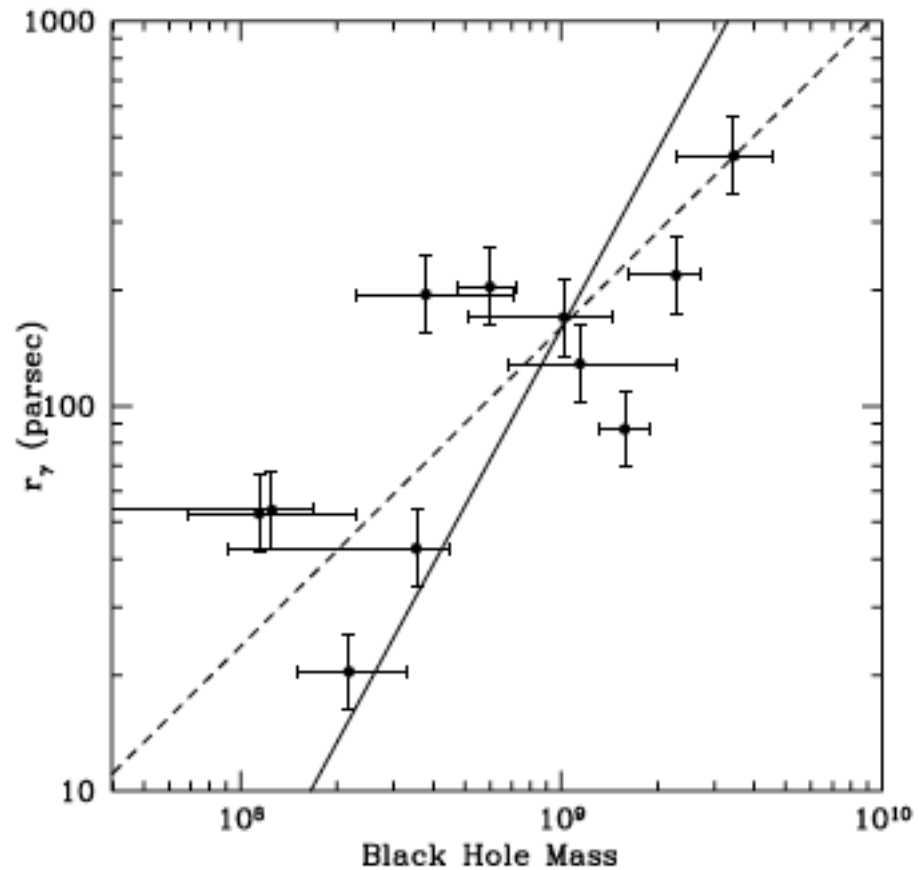


FIG. 7.—BH mass vs. core size, r_γ , for the 11 core galaxies that have M_\bullet measurements. The dashed line is the symmetric fit between M_\bullet and r_γ provided by eq. (24), while the solid line gives the fit presented in eq. (25), which assumes that r_γ is the independent variable. [See the electronic edition of the *Journal* for a color version of this figure.]

$$\log \left(\frac{r_\gamma}{\text{kpc}} \right) = (0.59 \pm 0.18) \log \left(\frac{M_\bullet}{10^9 M_\odot} \right) + (2.19 \pm 0.10)$$

How Big the SMBH in Holm 15A?

Table 2. Holm 15A: Black Hole Mass Estimates

Relation	M_{\bullet} [M_{\odot}]	Reference
$M_{\bullet} - \sigma$	$\sim 2.1 \times 10^9$	Kormendy & Ho (2013, Eqs. 6)
$M_{\bullet} - L_{K,bulge}^*$	$\sim 9.2 \times 10^9$	Kormendy & Ho (2013, Eqs. 7)
$M_{\bullet} - L_{V,def}$	$\sim 2.6 \times 10^{11}$	Kormendy & Bender (2009, Eq. 3)
$M_{\bullet} - r_b$	$\sim 1.7 \times 10^{11}$	Rusli et al. (2013, Eq. 13)
$M_{\bullet} - r_{\gamma}$	$\sim 3.1 \times 10^{11}$	Lauer et al. (2007, Eq. 26)

*Taking the entire galaxy as a classical bulge, and correcting the value of H_0

How Big the SMBH in Abell 85 Dark Matter Halo?

$$\frac{M_{\bullet}}{10^8 M_{\odot}} = 0.168 \left(\frac{V_{circ}}{200 \text{ km s}^{-1}} \right)^{5.45}$$

$$\sigma_{cl} = 752 \pm 34 \text{ km s}^{-1}, \quad V_{circ} = \sqrt{2} \sigma_{cl};$$

hence, the mass of the SMBH for Abell 85 DM halo is :

$$M_{\bullet} \sim 1.5 \times 10^{11} M_{\odot}$$

SMBH Dynamical Effects

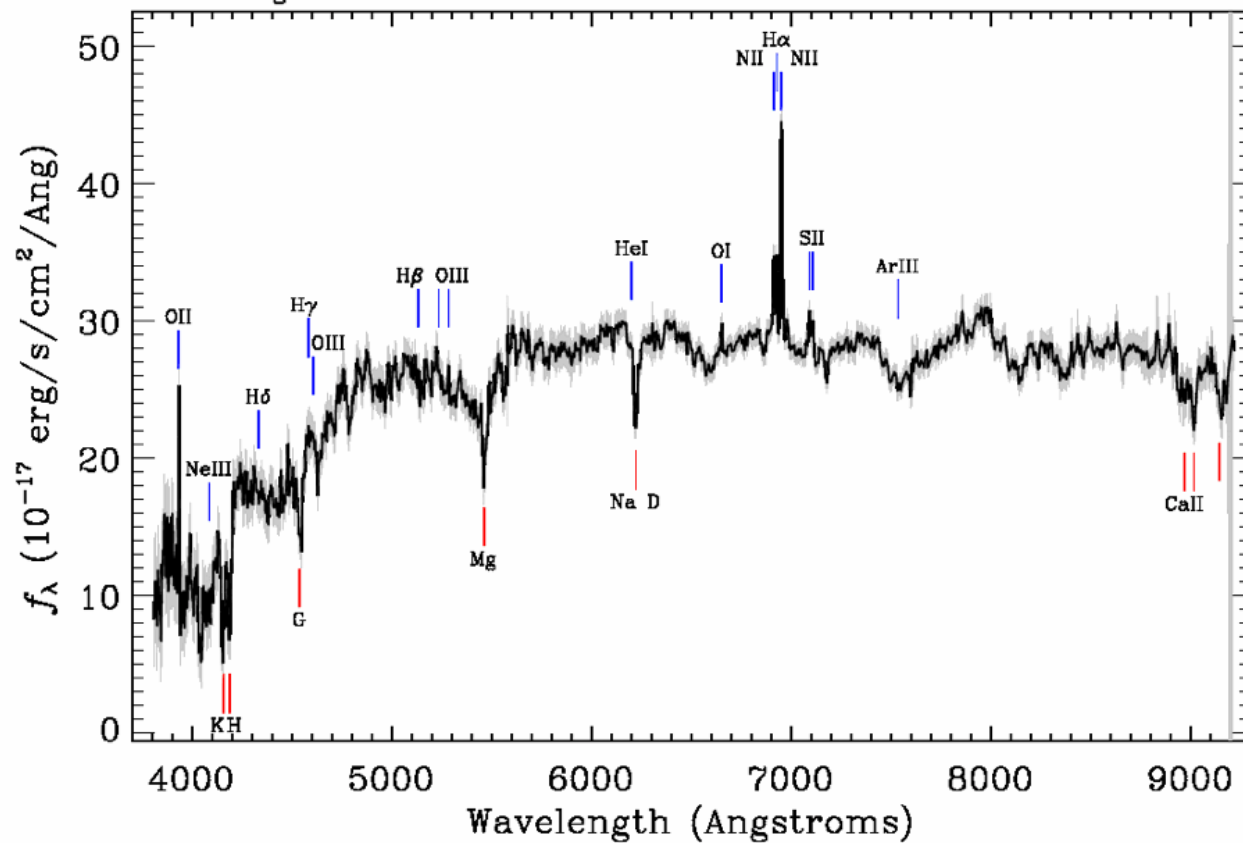
$$r_f = \frac{GM_\bullet}{\sigma^2} \text{ (Influence radius);}$$

$$M_\bullet = 10^{10} M_\odot \text{ and } \sigma \approx 310 \text{ km s}^{-1} :$$

$$r_f \sim 450 \text{ pc } (0.42'')$$

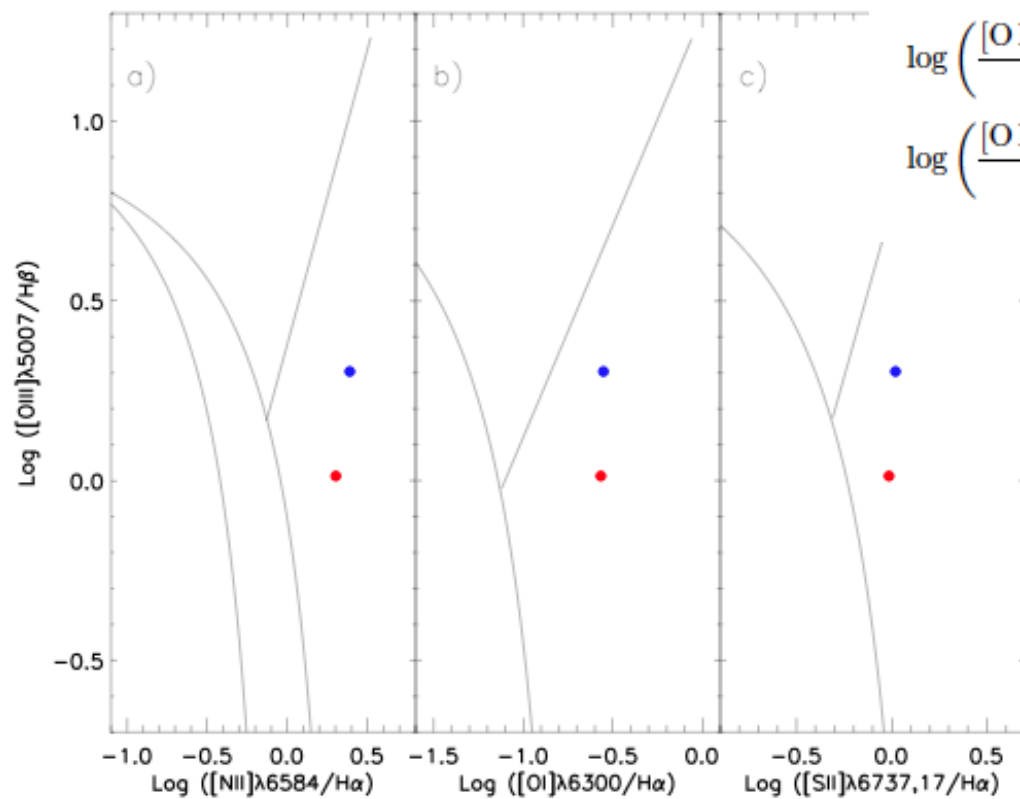
Holm 15A is a LINER

Survey: *sdss* Program: *legacy* Target: *GALAXY_RED GALAXY ROSAT_A ROSAT_D*
RA=10.46027, Dec=-9.30315, Plate=656, Fiber=390, MJD=52148
 $z=0.05536 \pm 0.00002$ Class=GALAXY AGN
No warnings.



Cores are explained by the scouring action of Supermassive Binary Black Holes (SBBH).

Is there a SBBH in HOLM 015A? Maybe...



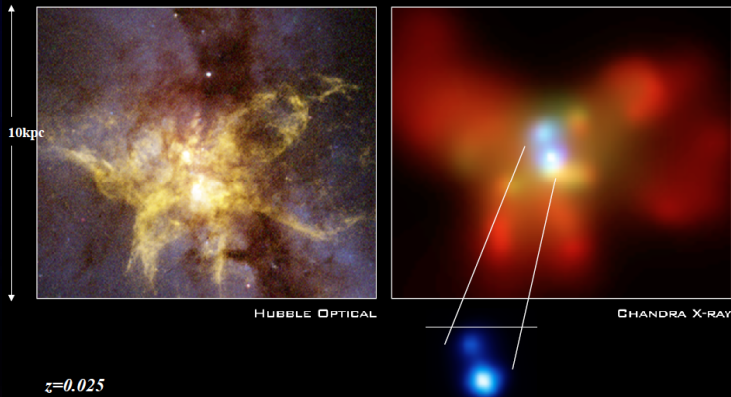
$$\log\left(\frac{[\text{O III}]\lambda 5007}{\text{H}\beta}\right) = 0.013, \log\left(\frac{[\text{N II}]\lambda 6584}{\text{H}\alpha}\right) = 0.302;$$

$$\log\left(\frac{[\text{O I}]\lambda 6300}{\text{H}\alpha}\right) = -0.567, \log\left(\frac{[\text{S II}]\lambda 6737,17}{\text{H}\alpha}\right) = -0.014.$$

A2261-BGC doesn't show emission lines.

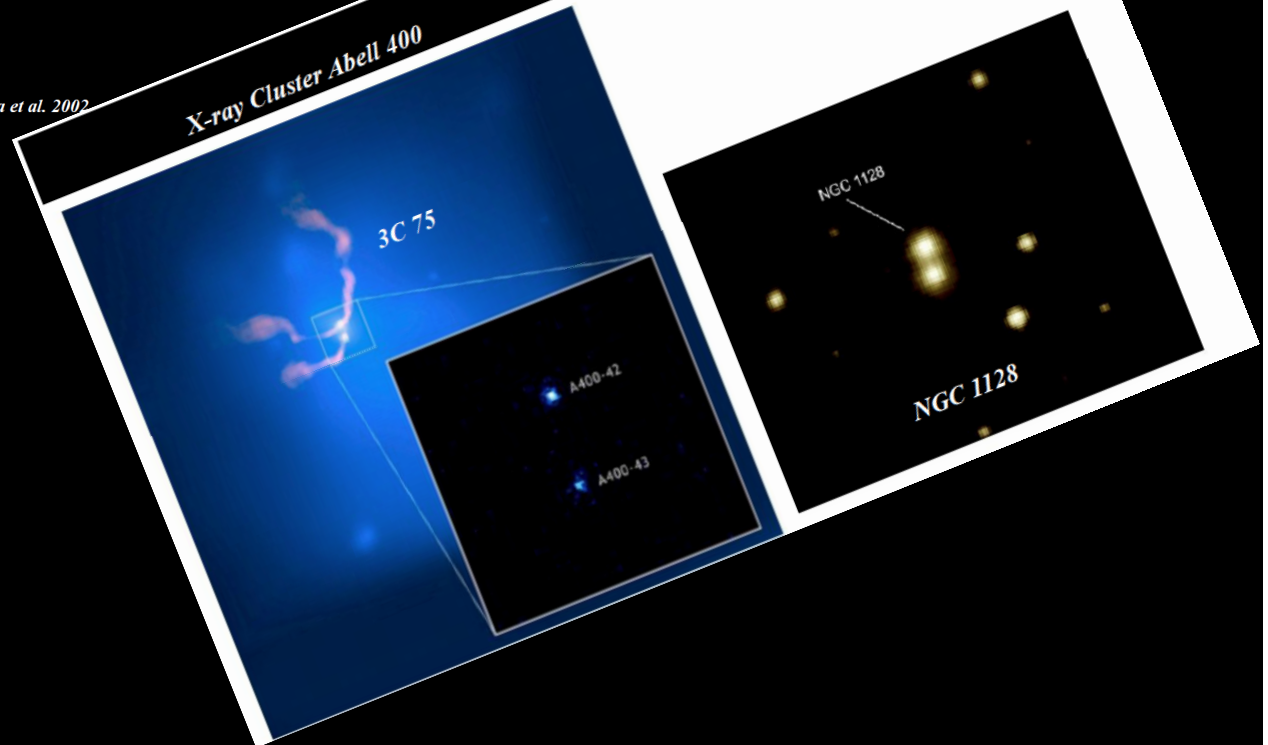
Are there SMBH binaries?

X-ray Image of a binary black hole system in NGC 6240



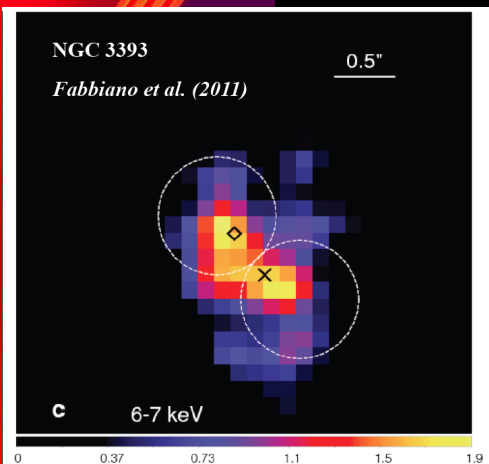
Komossa et al. 2002

X-ray Cluster Abell 400

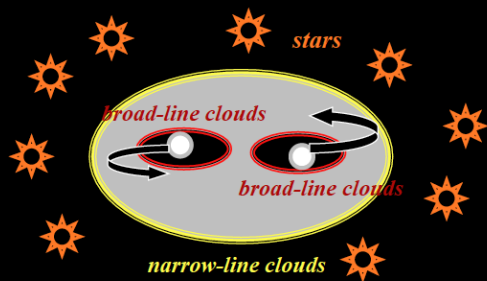


NGC 1128

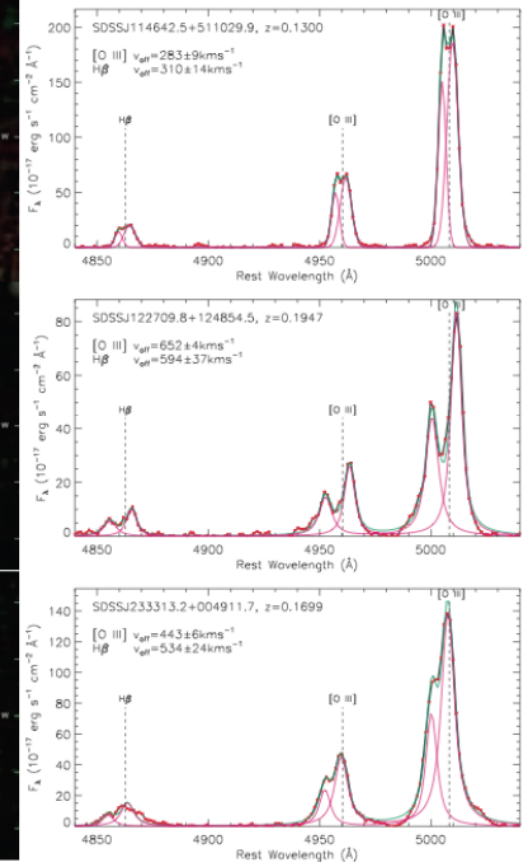
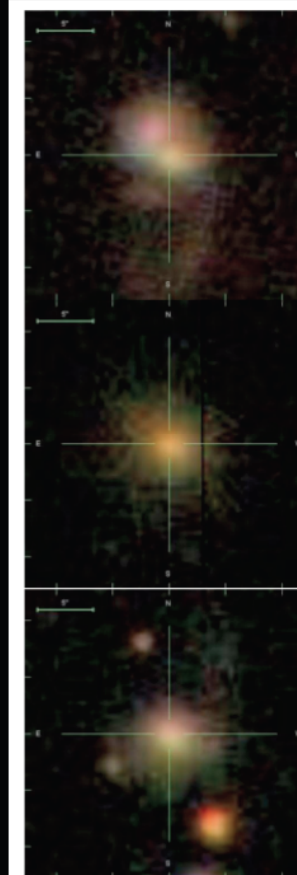
NGC 1128



SMBM binary



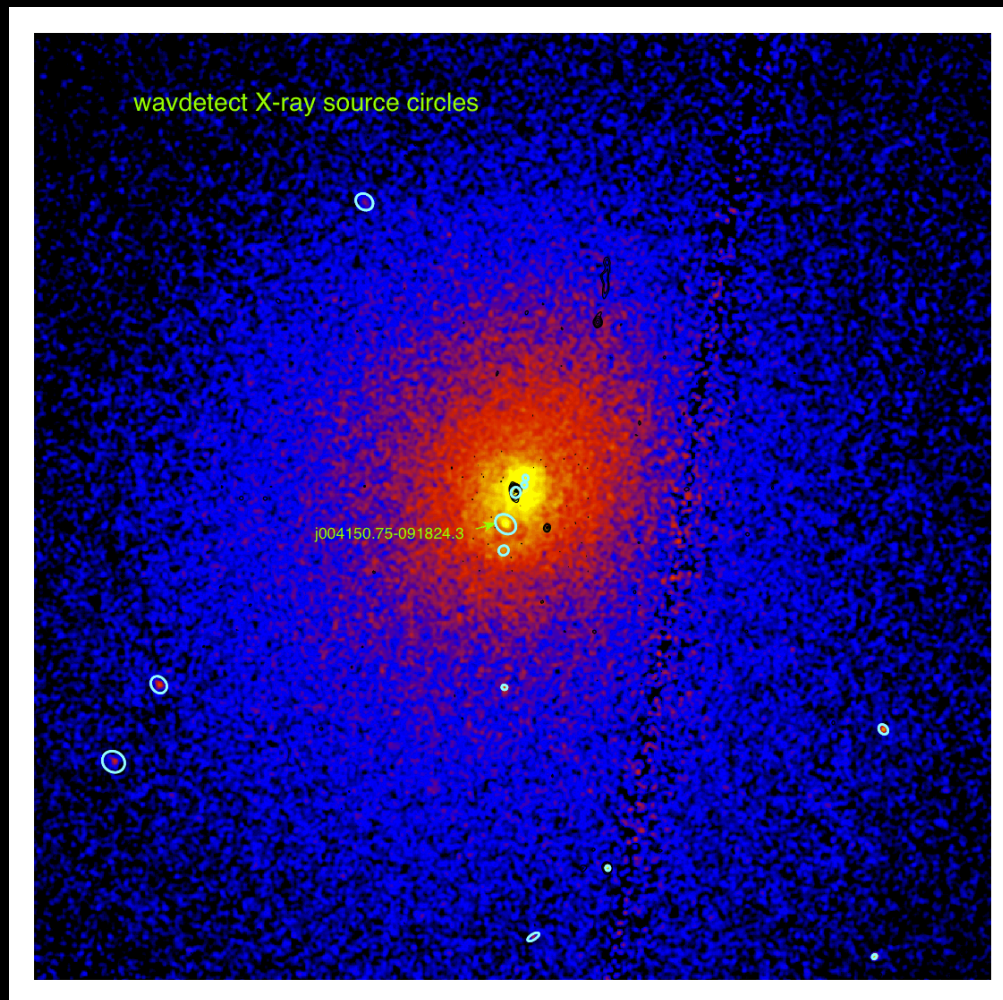
SDSS images & spectra of selected double narrow -line quasars



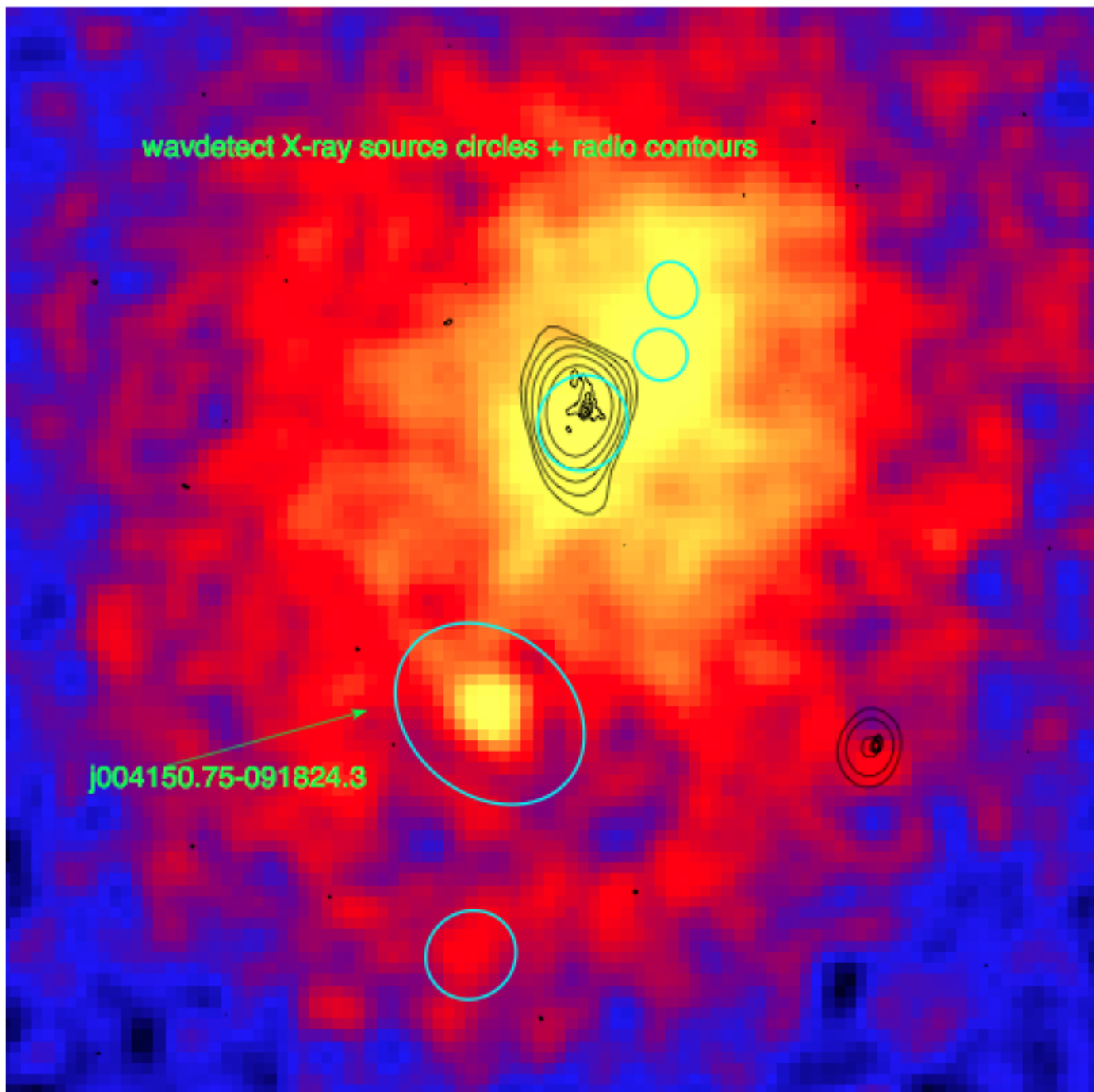
But Holm 15A has no double narrow-lines

(Liu et al. 2009)

Is there a SMBH binary?



Holm 15A



Holm 15A

Maybe: A candidate supermassive binary black hole system in the brightest cluster galaxy of RBS 797 ($z \sim 0.35$) reported by Gitti and collaborators 2014

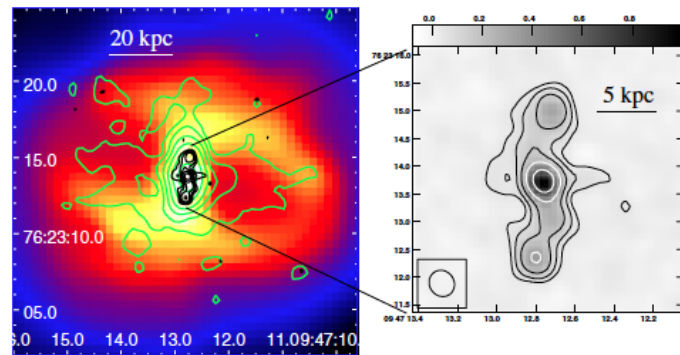


Fig. 2. The 4.8 GHz VLA contours obtained from the combined A- and B-array archival observations of RBS 797, imaged at different resolutions, are overlaid onto the *Chandra* image of the central region of the cluster (*left panel*). Green contours: 4.8 GHz VLA map at $1''.38 \times 1''.33$ resolution, obtained by setting ROBUST=+5, UVTAPER=250; the rms noise is $0.01 \text{ mJy beam}^{-1}$ and the contour levels are 0.03, 0.06, 0.12, 0.24, 0.48, and $0.96 \text{ mJy beam}^{-1}$; the total flux density is $\sim 4 \text{ mJy}$, with a peak flux density of $1.5 \text{ mJy beam}^{-1}$. Black contours (best visible in the zoom in the *right panel*): 4.8 GHz VLA map at $0''.49 \times 0''.44$ resolution, obtained by setting ROBUST=0, UVTAPER=0; the rms noise is $0.01 \text{ mJy beam}^{-1}$ and the contour levels are 0.04, 0.08, 0.16, 0.32, and $0.64 \text{ mJy beam}^{-1}$; the total flux density is $\sim 2.8 \text{ mJy}$, with a peak flux density of $1.0 \text{ mJy beam}^{-1}$.

EVN observations. 56 minutes on source

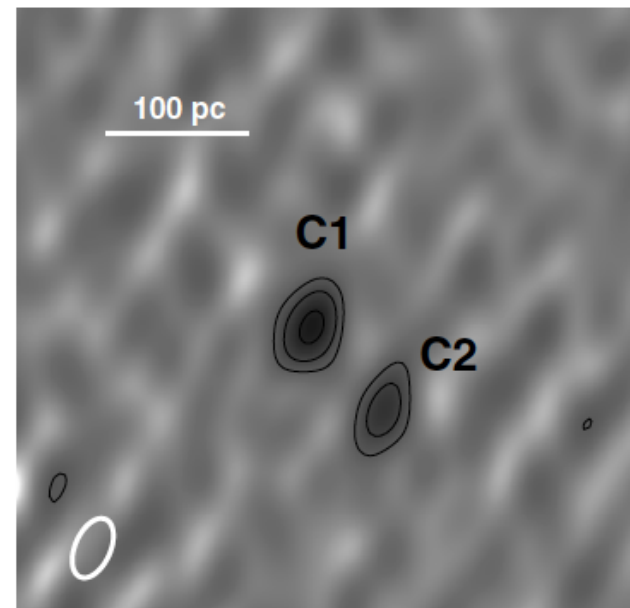


Fig. 1. 5 GHz EVN map of the BCG in RBS 797 at a resolution of $9.4 \times 5.3 \text{ mas}^2$ in PA -24° (the beam is shown in the lower-left corner). The rms noise is $36 \mu\text{Jy beam}^{-1}$ and the peak flux density is $0.53 \text{ mJy beam}^{-1}$. The contours levels start at 3σ and increase by a factor of 2. Two components, separated by 16 mas ($\sim 77 \text{ pc}$), are clearly detected.

Gitti et al. 2014

The Eddington Limit

Black Hole



*Outward
force*



$$\frac{L}{4\pi r^2 c} \sigma_T$$

*Inward
force*



$$\frac{GMm_p}{r^2}$$

$$L_E = \frac{4\pi GMm_p c}{\sigma_T} = 1.3 \times 10^{38} \frac{\text{erg}}{\text{s}} \left(\frac{M}{M_\odot} \right).$$

Black Hole Growth

$$L = \epsilon \dot{M} c^2 = \eta L_E \propto M$$

$$M = M_0 \exp\{t/t_E\}$$

$$t_E = 4 \times 10^7 \left(\frac{\epsilon/\eta}{10\%} \right) \text{ years}$$

Starting from a stellar mass, there is barely enough time to grow the observed quasar black hole during the age of the Universe at $z=7.1$ for $e=10\%$... but the radiative efficiency may be small due to trapping of radiation:

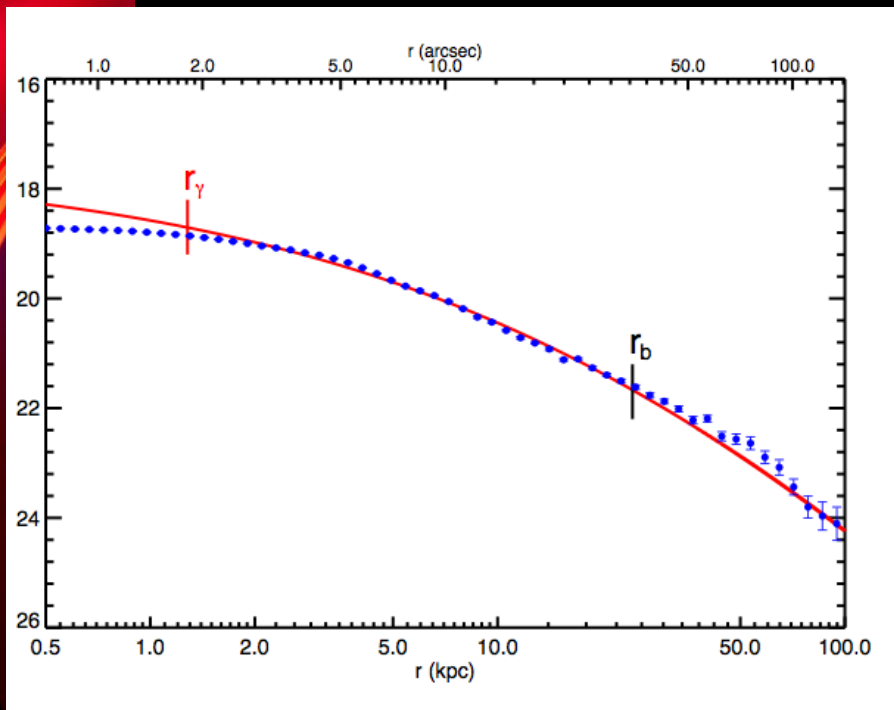
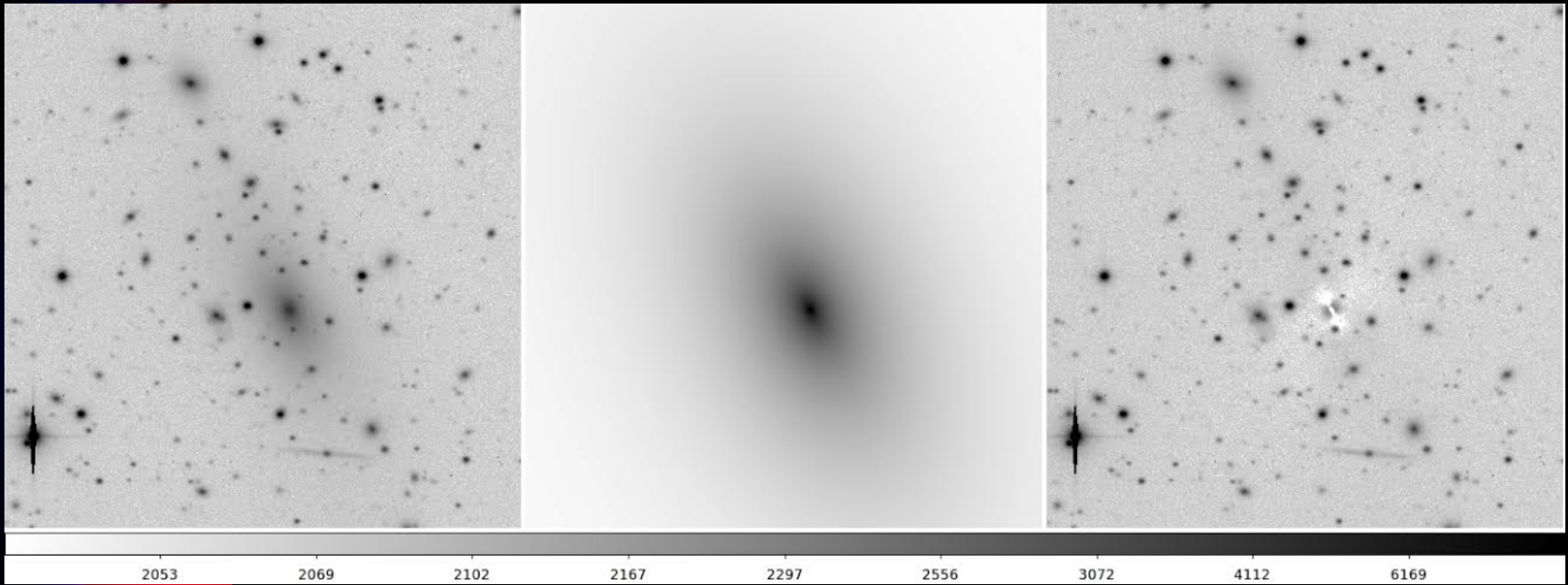
$$v_{\text{diff}} \sim (c/\tau) \ll v_{\text{infall}}$$



How about the brightest cD in the local universe?



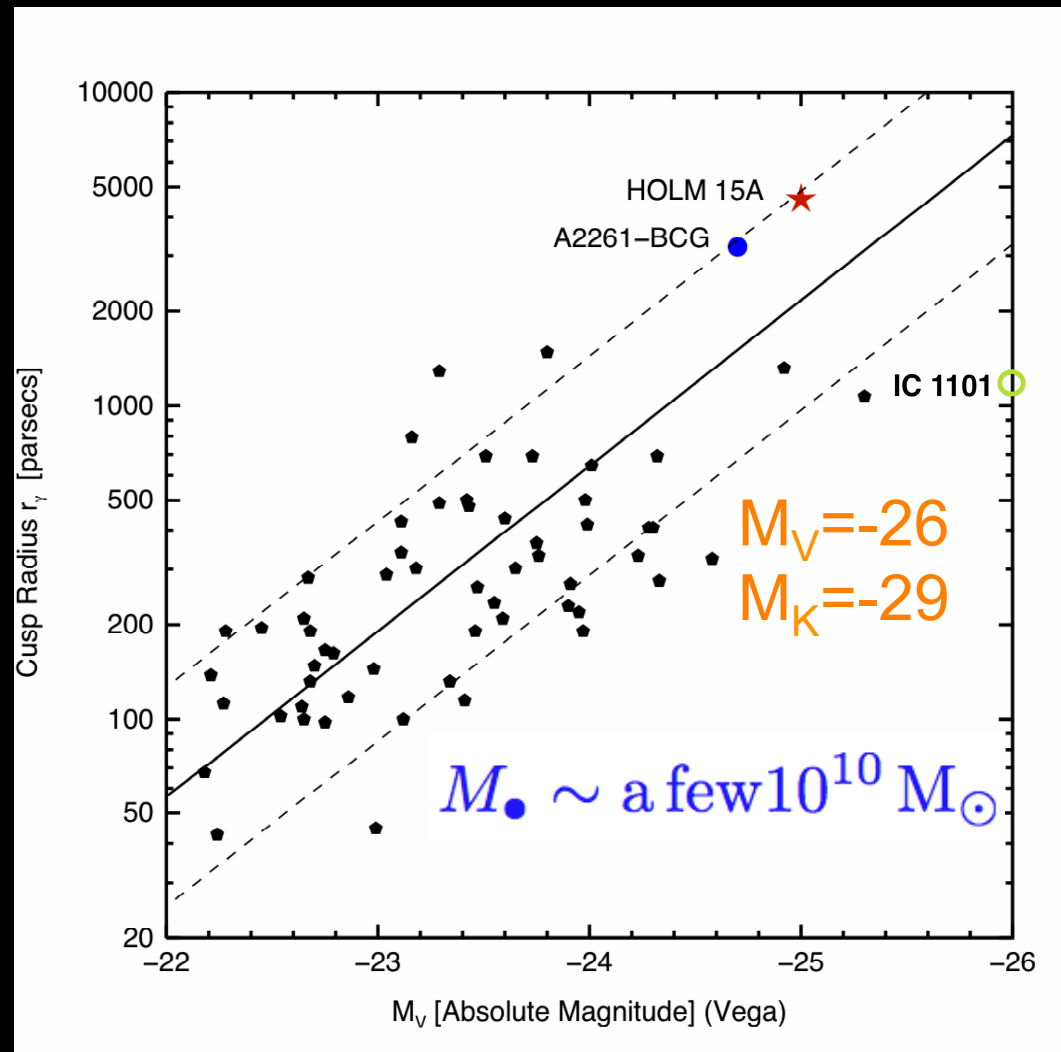
IC 1101 in Abell 2029,
 $z=0.077947$



IC 1101

$$r_\gamma \sim 1.3 \text{ kpc}$$

IC 1101 breaks away




Conclusions

We have found that Holm 15A has the largest core known so far. A central AGN supports the presence of a central BH, which could be ultramassive $\sim 10^{10} M_{\text{sun}}$. This seems to be a limit for maximum BH growth. Very large cores seem to be rare and may represent a relatively brief phase in the evolution of BCG, as the merging times for SMBH binaries should be relatively short (Kahn et al. 2014). After SMBH binary merger cusp regeneration would induce core shrinkage (Merritt 2006). If SMBH growth is regulated by galaxy mergers (e.g., Merritt 2006, Booth & Schaye 2011), their final masses were set, perhaps, by initial conditions (e.g., Treister et al 2013).

Holm 15A presents the best conditions to explore the effects of SMBH. We are proposing follow up observations (i.e., detection of BH gravitational lens using the background of SMG). I did not mention gravity waves nor dark mater annihilation ...

López-Cruz, O., Añorve, C., Birkinshaw, M., et al. ApJ , 795, L31



Desde el 15 de octubre de
2014, estamos en las
noticias en el mundo.

México (Excélsior), Guatemala, Perú, Colombia, Chile,
Argentina, Brasil, España (El País).

También en radio y televisión.

Recientemente (6 de noviembre), New Scientist, una
de las revistas de divulgación de la ciencia más
importante en el mundo.

15 de noviembre, 2015, Conozca Más

ASTRONOMÍA



Este es el núcleo de Abell 85, uno de los espectros visibles

HISTORIA DE UN DESCUBRIMIENTO ENORME

Curiosamente, las circunstancias de este «literalmente» gigantesco hallazgo son peculiares: no solo requirió varios años, sino también un descuido por parte de astrónomos de EUA. En 1995, Omar López realizó sus estudios de doctorado en la Universidad de Toronto, donde tuvo su primer contacto con la galaxia Holm 15A. En una parte de sus observaciones astronómicas de 1993, detectó que esta galaxia era peculiar, ya que tenía un bajo brillo superficial, pero no logró más debido a la falta de factores de comparación que le permitieron llegar a datos precisos, además de que las herramientas de análisis eran aún inadecuadas.

Para el 2012, los astrónomos estadounidenses Marc Postman y Tod Lauer, junto con un equipo de 15 investigadores dedicados a crear estallidos de supernovas en galaxias, y usando como herramienta al telescopio espacial Hubble, detectaron en la galaxia más brillante en el cúmulo Abell 2261-A 2261-BGG, un agujero negro masivo, el más grande descubierto hasta entonces. Este anuncio hizo retomar la

investigación de López Cruz debido a que las características descritas por ellos se asemejaban a las que él había estudiado en sus años en Toronto. Y las circunstancias, a poco más de una década después, permitieron al investigador del INAOE tener las herramientas suficientes para avanzar en el estudio de Holm 15A.

A cada integrante del equipo que con él (mencionado en la entrevista con él) le asignó una tarea en específico, ya que la investigación de la galaxia requirió de utilizar datos de su propia tesis doctoral e imágenes de diferentes telescopios ubicados en el espacio y en tierra, y procesarlas en supercomputadoras que corren programas escritos especialmente para atender este estudio.

El esfuerzo colaborativo de los científicos y la tenacidad del Dr. Omar López mostró un resultado impresionante: en el centro de la galaxia Holm 15A vive un agujero negro con una escala de tamaño, mucho más grande que el de la A2261-BGG. Y así estos investigadores encontraron el mayor hoyo negro jamás hallado.

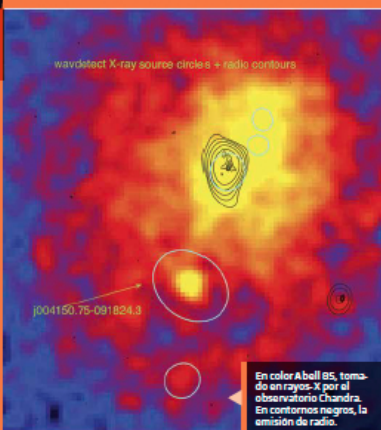


Fotografía de Abell 85, desde el Kitt Peak National Observatory en Arizona

Foto: Omar López Cruz

En el artículo "En verdad existen los hoyos negros" (de la edición anterior), el astrónomo y divulgador británico Marcus Chown expone qué es un agujero negro, cómo se han clasificado, además de las ideas que hoy en torno de estos cadáveres de estrellas supergigantes. Te invitamos a leerlo para conocer más del tema.

El equipo estadounidense había estudiado a Holm 15A y no le dio importancia. **¡Se les fue el descubrimiento!**



En color Abell 85, tomado en rayos-X por el observatorio Chandra. En contornos negros, la emisión de radio.



DR. OMAR LÓPEZ CRUZ, DEL INAOE

De primera mano, el investigador en jefe del proyecto nos explica todo lo relacionado a este importante descubrimiento astronómico.

¿CÓMO SE HIZO ESTE DESCUBRIMIENTO?

Dr. Omar López Cruz: Es un proceso bastante largo. Desde 1995 nos dimos cuenta de que la galaxia central, la más brillante de Abell 85 [Abell es un cúmulo de galaxias], a la que llamamos Holm 15A, tenía la peculiaridad de ser de muy bajo brillo superficial. Es decir, que en lugar de estar concentrada hacia el centro, era difusa. ¡Incluso llegamos a decir que era «esponjosa»! Aunque ya en 1980 un colega había reportado que ésta tenía propiedades especiales, no sabíamos cuáles eran. No teníamos referente. Entonces, gracias a un programa nuevo dirigido por mis colegas Postman y Lauer, usando el telescopio espacial Hubble, equipado ahora con una cámara especial para abarcar un área grande del Universo y así buscar supernovas en galaxias - particularmente querían establecer la tasa de supernovas en galaxias elípticas -, se descubrió que en la galaxia Abell 2261-BGG existía un aplanchamiento en el centro, y la declararon como la de mayor núcleo («core», en inglés) que jamás se haya medido. Este «aplanchamiento» se creó que se debe a la interacción de dos agujeros negros binarios.

Por las leyes de escalamiento que aplicamos, inferimos que entre mayor sea dicho core, mayor será el agujero negro en la galaxia. Pero no se detectó ninguna firma de un agujero negro binario, ni actividad de AGN [ver recuadro] por el espectro, ni gas en A2261-BGG. Y entonces nos acordamos de Abell 85.

Yo no la quise medir. Le di los datos a mi estudiante Christopher Ahn. Su trabajo es muy importante, ya que él llevó a cabo un programa de ajuste de brillo superficial. Así, ya teníamos la herramienta hecha. Él me dio los resultados para interpretarlos y vimos que Holm 15A era mucho mayor que Abell 2261-BGG.

¿ES EL AUTOR PRINCIPAL DE LA INVESTIGACIÓN, QUIÉNES SON LOS COLABORADORES QUE MENCIONAS?

Los astrónomos colaboramos con mucha gente. Christopher Ahn, de la Universidad Autónoma de Sinaloa, hizo el ajuste de brillo superficial. M. Birkinshaw analizó los datos en [el espectro de] radio, y D.M. Worrall, en el de rayos X. Ambos provienen de la Universidad de Bristol. Héctor Ibarra Medel, mi estudiante del INAOE, hizo la dinámica del cúmulo. Wayne A. Barkhouse, de la Universidad de Dakota del Norte, desarrolló el análisis de una dimensión. Juan Pablo Torres Papaquín, de la Universidad de Guanajuato, realizó el análisis del espectro con el que determinamos que es un tipo de AGN de baja luminosidad. Verónica Motta, de la Universidad de Valparaíso, nos ayudó con la dinámica.

¿PLATICAMOS SOBRE EL GRUPO ESTADOUNIDENSE AL QUE SE LE ESCABULLÓ ESTE DESCUBRIMIENTO?

Marc Postman y Tod Lauer son grandes investigadores que han conducido programas muy importantes en cuanto al estudio de agujeros negros y galaxias elípticas... Ellos y su grupo son el referente en el estudio de agujeros negros. Cuando encontramos y reportaron la galaxia Abell 2261-BGG, se abrió la oportunidad de hacerle justicia a Holm 15A. Eso fue lo que me hizo tener que pensar: «Veamos si es más grande». Como te dije antes, le pedí a Christopher que la midiera para no sesgar dicha medición, porque en el fondo quería que fuera más grande. Así que Christopher la midió y resultó que es mayor por un kiloparsec [ver recuadro «poniéndolo en perspectiva»], es decir, ¡descubrimos un monstruo! Y con eso rompimos el récord establecido en 2012.

PONIÉNDOLO EN PERSPECTIVA

¿Qué es eso de «kiloparsec» (o kpc)? Simple: es una unidad astronómica para medir la distancia hacia objetos fuera de nuestro sistema solar. Un parsec equivale a aproximadamente 3.26 años luz, por lo que mil parsecs o un kpc son unos 3,260 años luz.

¿QUÉ APORTA ESTE DESCUBRIMIENTO A LA EVOLUCIÓN DE GALAXIAS?

A los agujeros negros de 10⁹ veces la masa del sol los llamamos supermasivos. Con esto hablamos de agujeros ultramasivos de 10⁹. Esto nos dice que quizá los agujeros negros no crezcan más. ¿Por qué? Aún no lo sabemos.

¿QUÉ IMPLICA ESTE TRABAJO PARA LA ASTROFÍSICA NACIONAL?

Que este es el nuevo paradigma en las bases de datos. Para hacer este trabajo sólo usamos datos de mi tesis y todo lo demás, son datos públicos del SLOAN, del telescopio espacial de rayos X Chandra, del observatorio orbital XMM-Newton, también de rayos X, entre otros. El resultado nos dice que aún no hemos explotado las bases de datos públicas y que todavía hay mucho por descubrir.

¿CÓMO TE SIENTES CON TU DESCUBRIMIENTO?

Imagina que tienes un problema que empezó en 1995, que no te lo has quitado de encima en 19 años y de repente lo acomodas. ¡Se siente maravilloso! Es un gran alivio y regocijo, y lo que no anticipamos es que esta noticia trascienda de esta manera.

¿CÓMO ES EL SIGUIENTE PASO?

Solicitar el telescopio Hubble y el Gran Telescopio de Canarias para hacer un estudio dinámico. ☺

NewScientist

WEEKLY November 9-14, 2014



Flight time:
3907 days

Distance travelled:
6 billion kilometres

Touchdown:
12 November 2014, 15:38 UTC

COMET CHASER

Can Rosetta pull off the most daring space encounter ever?

Science and technology news
www.newscientist.com
Focus on diversity
No2994 US\$5.95 CAN\$3.95



MY FAMILY AND OTHER STRANGERS

When you can't recognize your loved ones

ACID TEST
Can you really dissolve a body?

GOODBYE KITTY
Schrödinger's cat is finally dead

CHIP FOR BRAINS
The man who is remaking the computer in your head

IN BRIEF



On tamedness and tail-dropping in island lizards

LIZARDS are famed for a rather extreme escape tactic: they shed their tails to avoid predators. But for Erismia's wall lizard, found across Greece's Cyclades islands, its tendency to jettison the tail – and indeed its tolerance for disturbances – depends on which island it lives on.

The Cyclades became isolated from each other more than 11,000 years ago. This stranded such island lizards, meaning they diverged into different predators. As a result, the lizards, *Palafoxia erismiae*, have been walking different evolutionary paths. Kelsey Drick of the University of Michigan, Ann Arbor, tested the lizards' fear

factor by walking towards 813 male lizards from 27 islands plus the mainland. The average lizard ran when Drick got within 1.8 metres, but some lizards got to 10 centimetres.

She found a link between the distance zone in which a lizard flees and the island it hailed from. Lizards from smaller islands with fewer types of predators, or that had been disconnected from mainland predators for longer, let her come far closer compared with their more fearful relatives (*Evolution*, doi.org/10.1111/evo.12477).

In lab tests, Drick also found that lizards from safer islands were less likely to ditch their tails when put under physical pressure through prodding. "If you're living in a predator-free environment, it would be evolutionarily disadvantageous to spend your time running away while other lizards are foraging and mating," she says.

Printable transistor can tell what ails you

GUT a respiratory bug? Printable plastic transistors that can directly detect pathogens in blood or saliva could one day tell you if you're in a flask.

Transistors easily switch on and off in response to an applied voltage, which allows electric current to flow between two terminals. But they can respond to other things as well, says Mariana Medina-Sánchez, now

at the Leibniz Institute for Solid State and Materials Research in Germany. She and her colleagues have engineered an inkjet-printable transistor that can recognise the protein biomarkers of common diseases, switching on only when it has detected them.

The team printed a transistor using a special ink embedded with a common antibody called human immunoglobulin G, which binds

to antigens from a number of common viruses, bacteria and fungi. When a disease protein binds to an antibody, it changes the transistor's electrical properties, altering the voltage level at which it turns on (*Advanced Functional Materials*, doi.org/10.1002/adfm.201401016).

Eventually, the team says, doctors could print out a sheet of the device – each equipped with a different disease antibody – and diagnose people in a snap.

Plants make their own sunscreen

THEY bask in the sun for hours, but just like us, plants need to protect themselves from damaging ultraviolet rays. Now we know how they do it.

Many plants use a group of chemicals called strigolactones to defend against the sun while they absorb light for photosynthesis. These aromatic compounds sit in the upper cell layers of these plants' leaves and one type – strigoyl malate – provides the bulk of this UV protection.

A team led by Timothy Zwier of Purdue University in West Lafayette, Indiana, has probed how strigoyl malate works. They found that it filters out the entire spectrum of UV-B radiation, which is known to damage plant and human DNA.

"It can absorb all wavelengths of UV-B radiation, with no gaps in coverage," says Zwier, who says his finding could be useful for developing more resistant plants.

Giant galactic core in black hole battle

THE largest galactic core ever seen may be the nemesis of a battle for black hole supremacy.

Cesar Lopez, Chief of Mexico's National Institute of Astrophysics, Optics and Electronics and his team measured the core of galaxy Holm 15A, 650 million light years from Earth. They found it was a record-breaking 15,000 light years across – about one sixth the diameter of the entire Milky Way.

The core's size suggests the black hole it hosts could weigh 100 billion times the mass of our sun – nearly as much as the Milky Way (*Astrophysical Journal Letters*, doi.org/10.1088/2031-5122/aa000000). If so, it probably formed as two or three separate black holes jostled for position before merging into one, puffing up Holm 15A's core in the process.